
Mach number dependence of the turbulent dynamo: solenoidal vs. compressive flows

...or „The forcing matters“

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17.11.2011



Turbulence forcing

Purely hydrodynamical, supersonic turbulence

MHD turbulent dynamo (subsonic versus supersonic)

Polaris Flare: Bensch, Stutzki, Ossenkopf 2001

distance ~ 150 pc

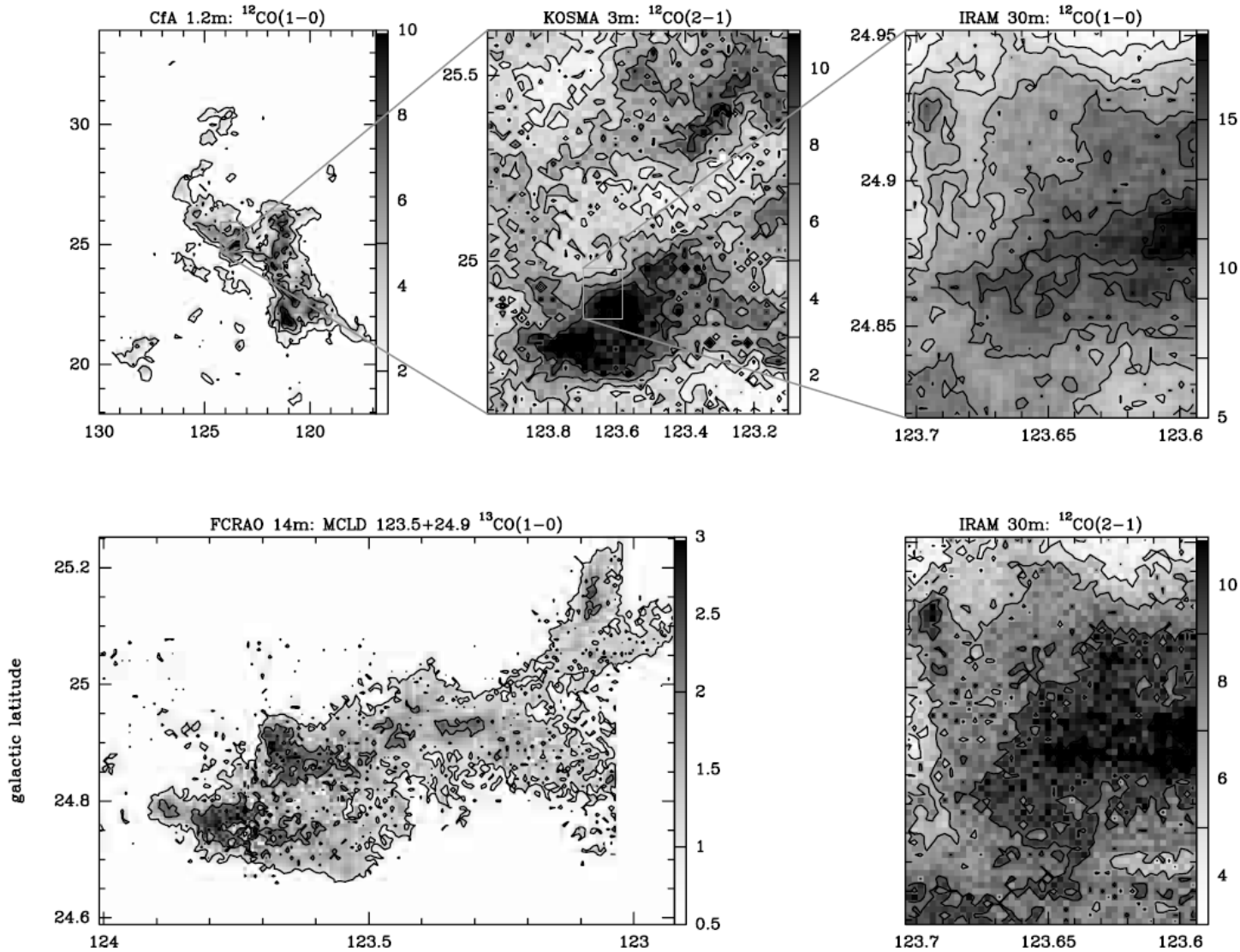
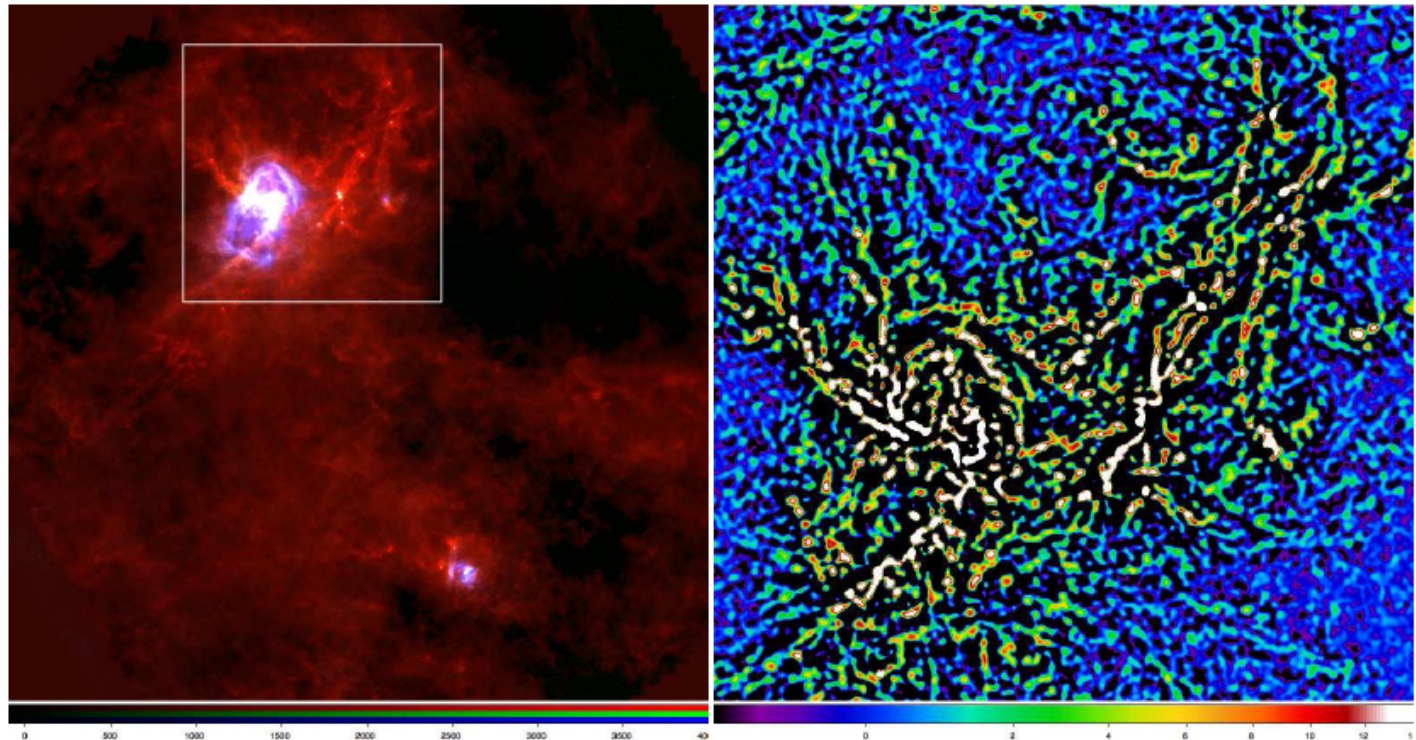


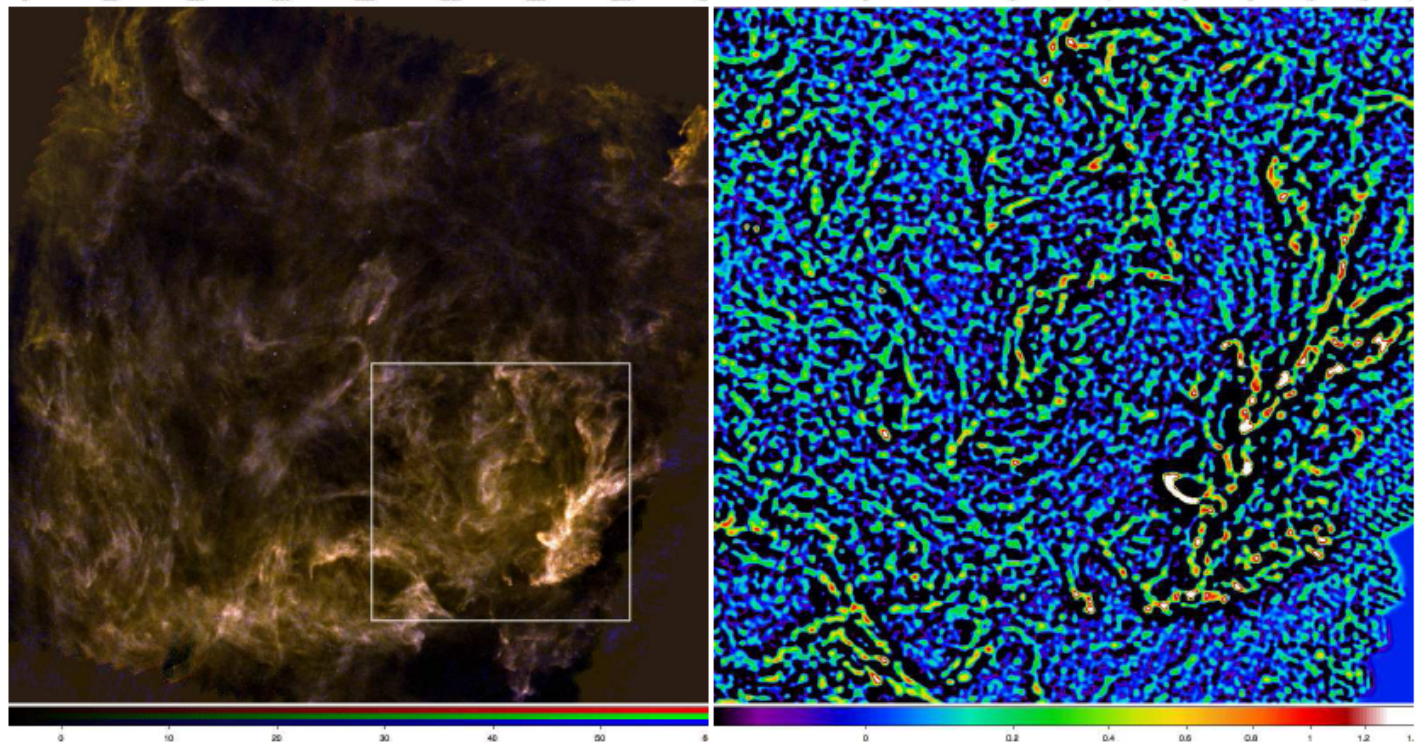
Fig. 8. Velocity integrated spectral line maps of the rotational transition $^{12}\text{CO } J = 1 \rightarrow 0$, $^{12}\text{CO } J = 2 \rightarrow 1$ and $^{13}\text{CO } J = 1 \rightarrow 0$, observed towards the Polaris Flare, and one of its cores, MCLD 123.5+24.9. The transition and the telescope are indicated at the top of each panel. The line intensity is given in main beam brightness temperature, T_{mb} . Iso-intensity levels are shown from 2 to 8 in steps of 2 (CfA map), 1 to 11 by 2 (KOSMA), 1 to 4 by 1 (FCRAO), 5 to 17 by 2 (IRAM, $^{12}\text{CO } J = 1 \rightarrow 0$), 3 to 11 by 2 (IRAM, $^{12}\text{CO } J = 2 \rightarrow 1$), in units of K km s^{-1}

Aquila and Polaris:
Men'shchikov et. al.
2010

Aquila
FOV 15pc
(star-forming)



Polaris Flare
FOV 9pc
(quiescent)

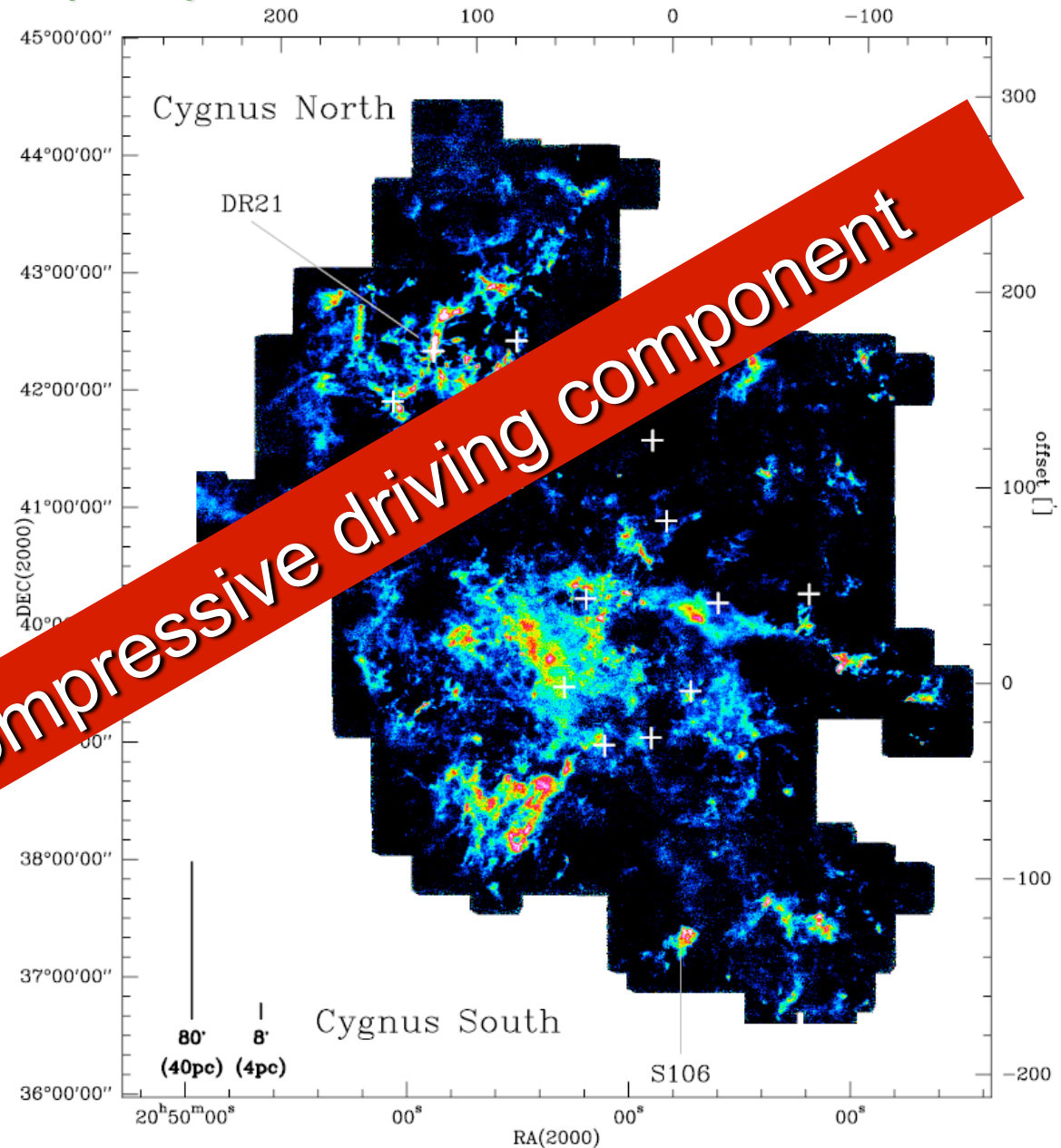


Cygnus X: Schneider et al. (2011)

Giant molecular cloud complex

Turbulence driven by
Supernova explosions,
Ionization fronts,
Protostellar jets/winds,
Gravitational collapse,
Galactic rotation, MRI

Significant compressive driving component



Turbulence driving experiments

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Typical setup for forced turbulence simulations:

e.g., Vazquez 1994, Padoan+ 1997, Passot+ 1998, Stone+ 1998, Mac Low 1999, Klessen+ 2000, Heitsch+ 2001, Cho+ 2002, Boldyrev+ 2002, Li+ 2003, Padoan+2004, Jappsen+ 2005, Ballesteros+ 2006, Kritsuk+ 2007, Dib+ 2008, Offner+ 2008, Kowal+ 2008, Schmidt+ 2009, Cho+ 2009

- Periodic boundary conditions **“Turbulence in a box”**
- Isothermal EOS: $P = c_s^2 \rho$
- Neglect self-gravity
- Driven to supersonic speeds (Mach 5-10)
- Large-scale **Forcing term f**

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) = 0$$

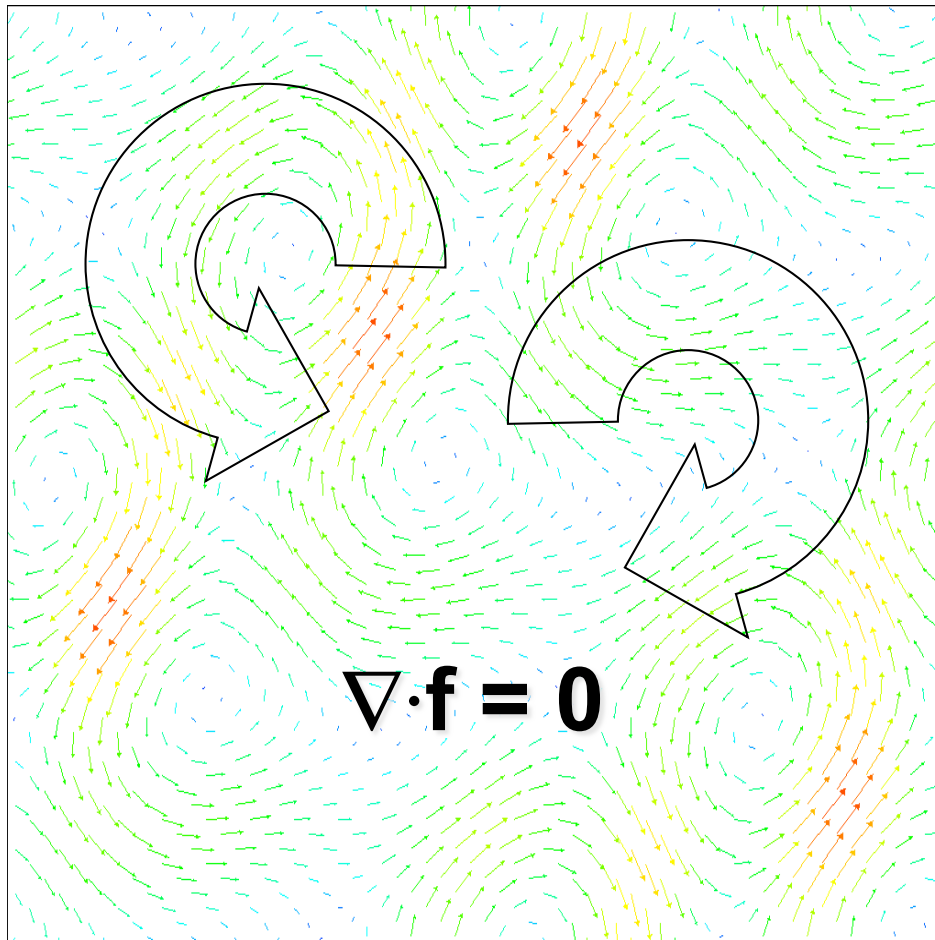
$$\frac{\partial \mathbf{v}}{\partial t} + (\mathbf{v} \cdot \nabla) \mathbf{v} = -\frac{1}{\rho} \nabla P - \nabla \Phi + \mathbf{f}$$
~~$$\frac{\partial}{\partial t} (\rho c) + \nabla \cdot [\mathbf{v}(\rho c + P)] = \rho \mathbf{v} \nabla \Phi + \rho \mathbf{v} \cdot \mathbf{f}$$

$$\Delta \Phi = 4\pi G \rho$$~~

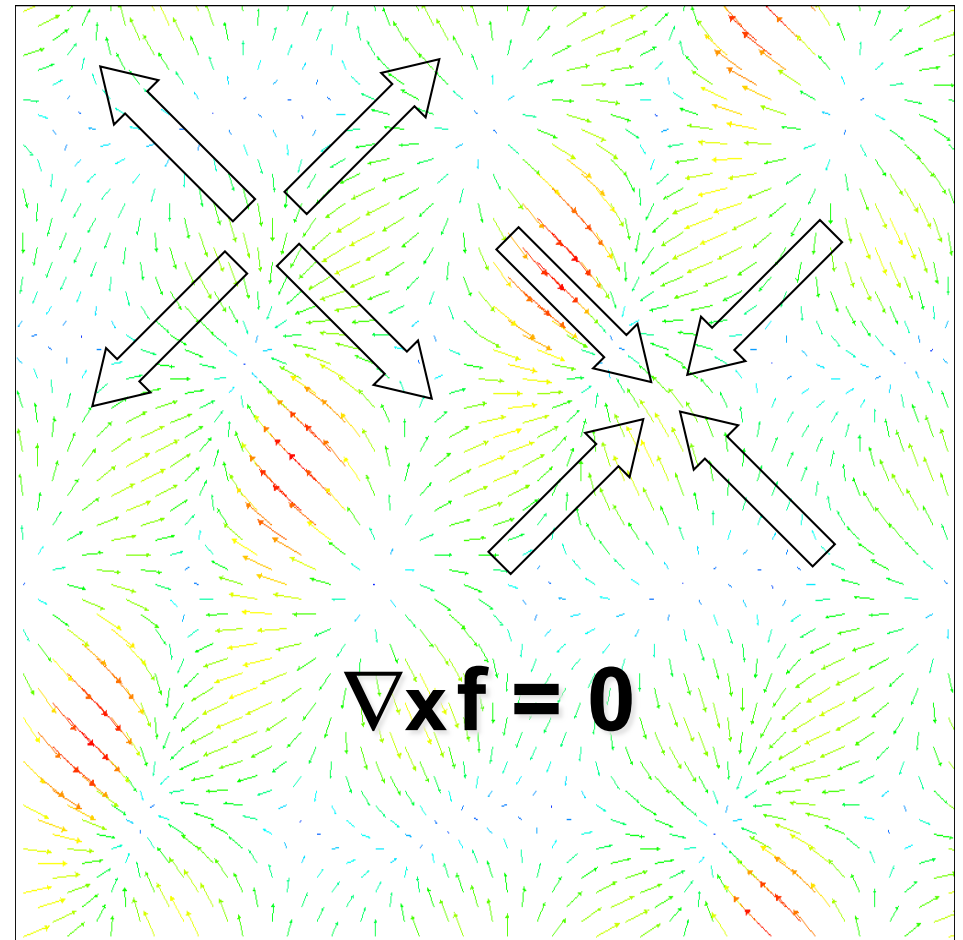
Turbulence forcing – solenoidal versus compressive

Ornstein-Uhlenbeck process (stochastic process with autocorrelation time)
→ **forcing varies smoothly in space and time**
following a well defined random process

Solenoidal forcing



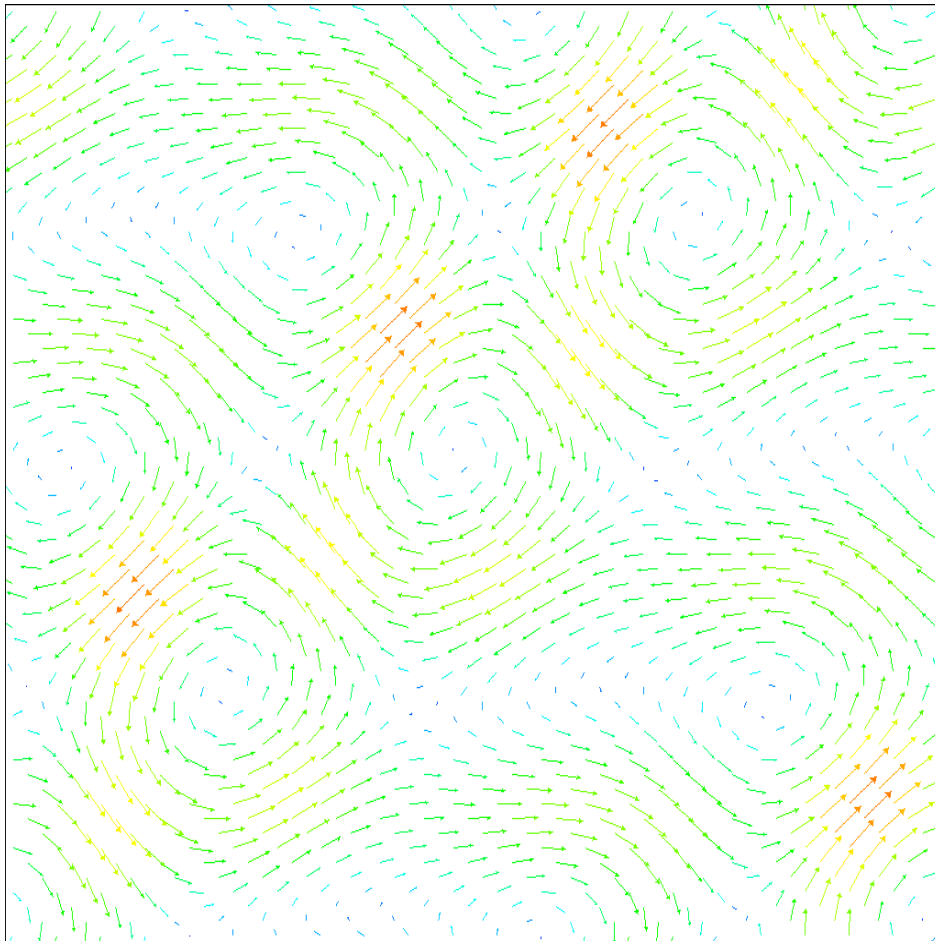
Compressive forcing



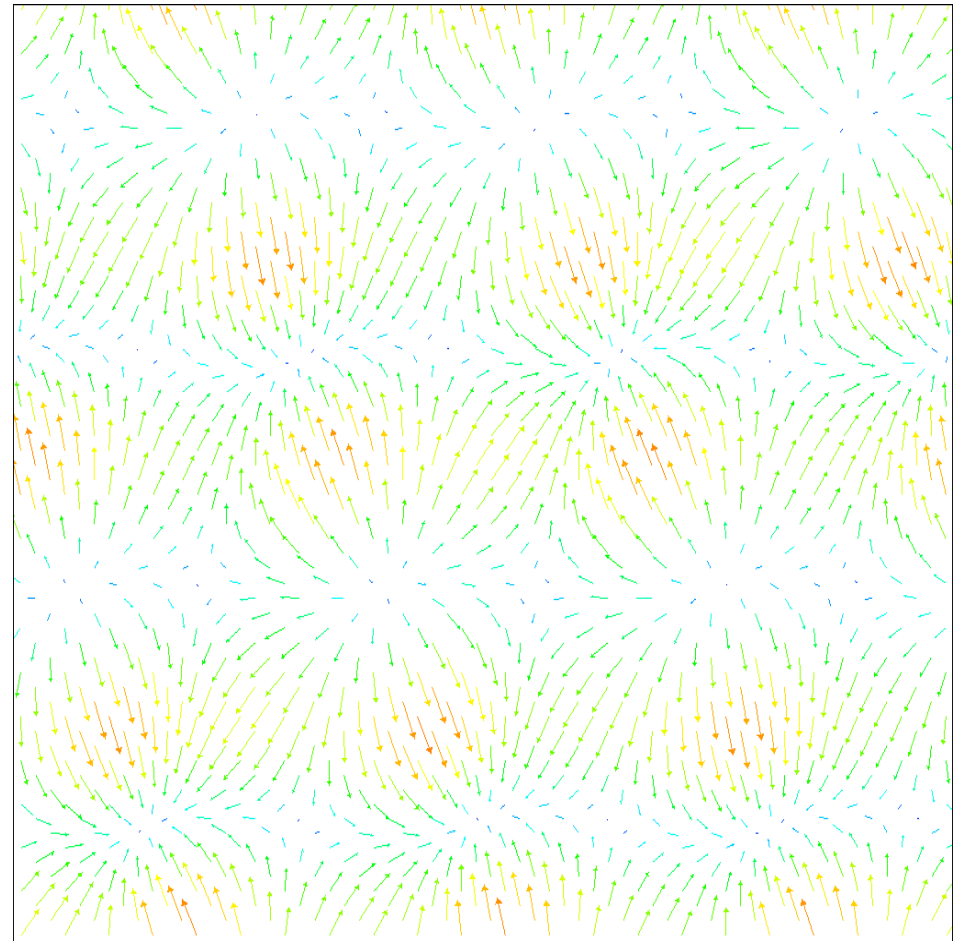
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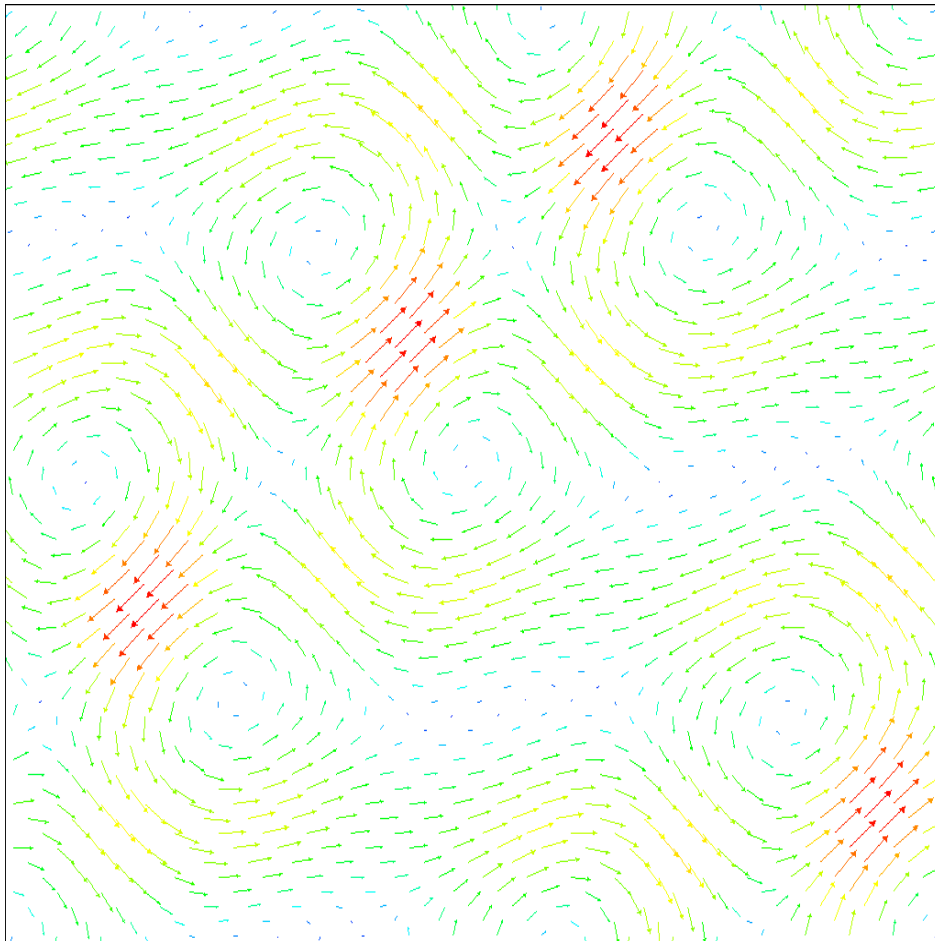
Compressive forcing



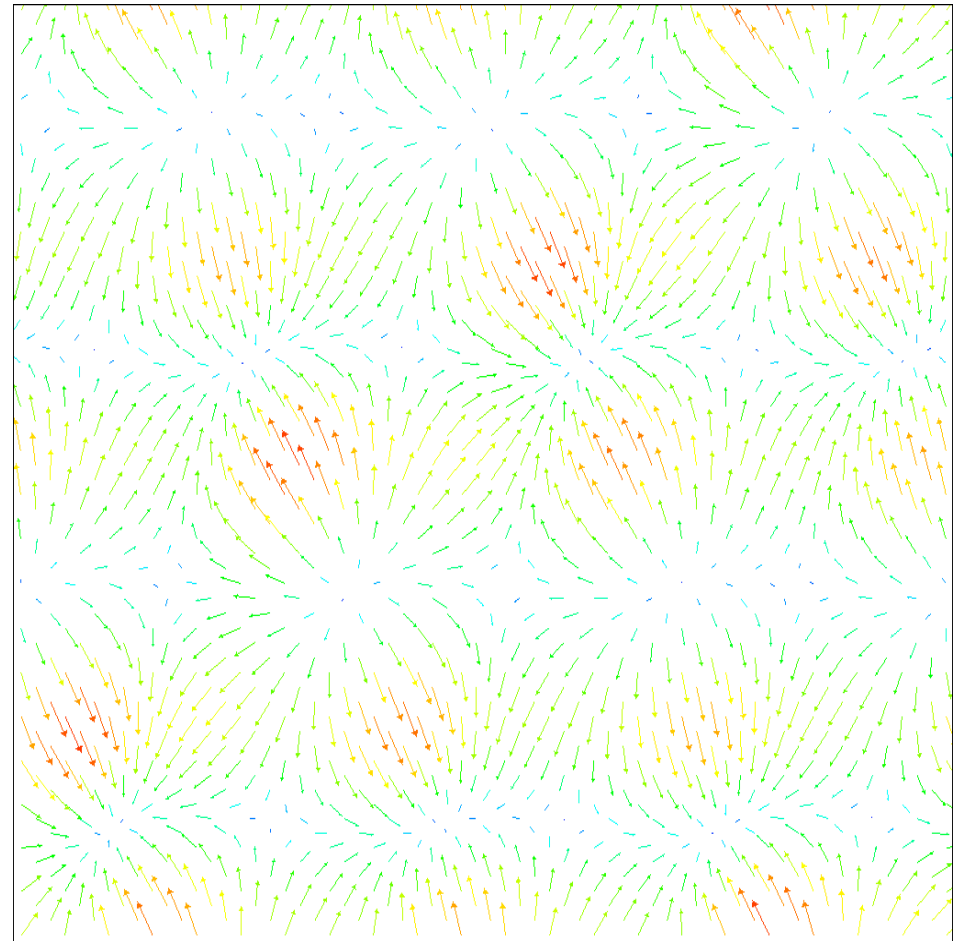
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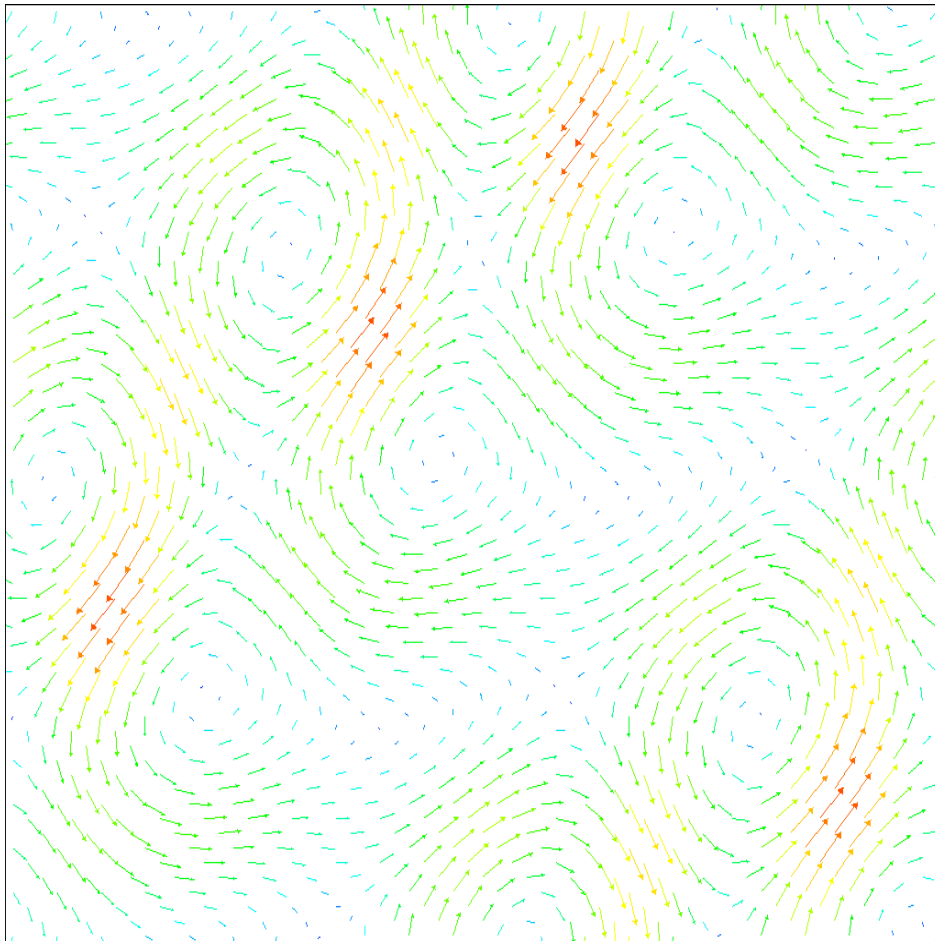
Compressive forcing



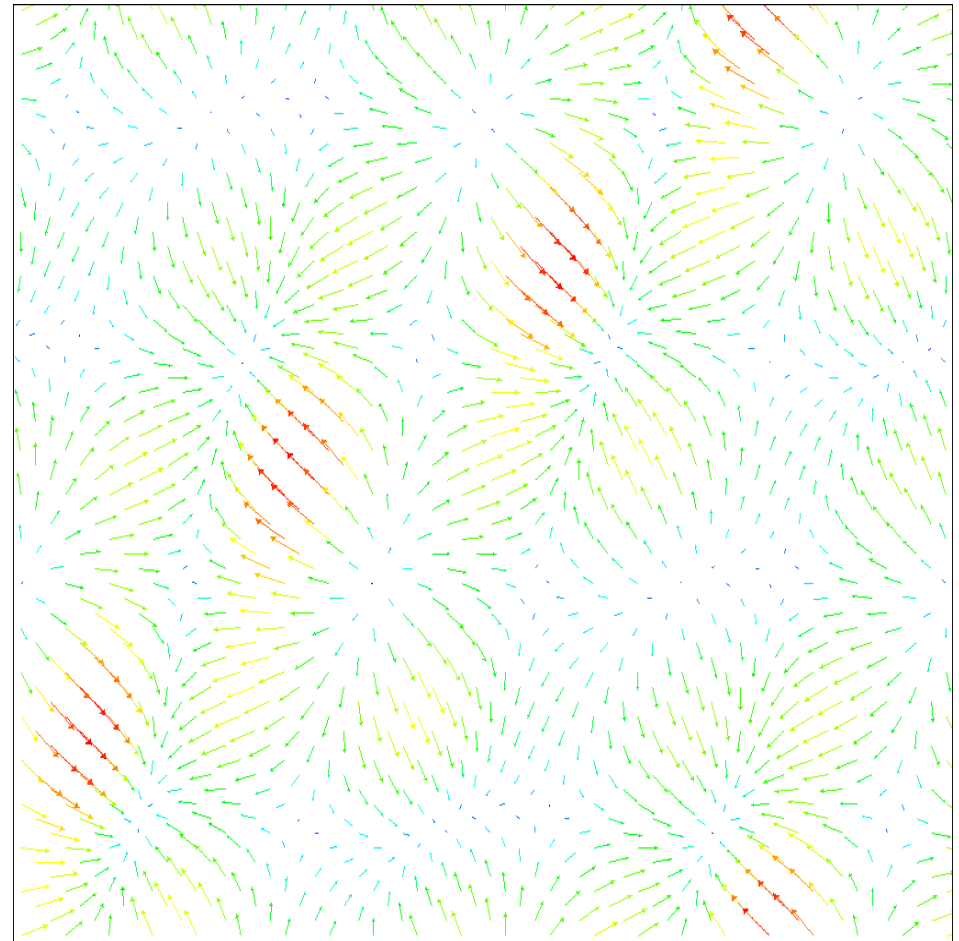
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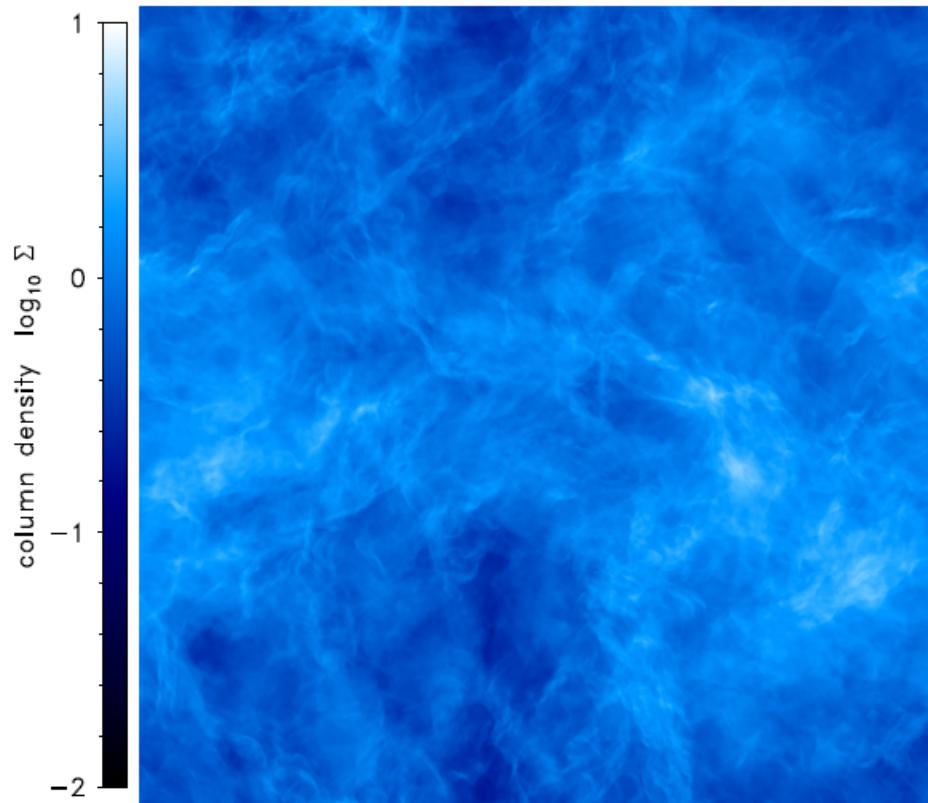
Turbulence driving experiments

Purely hydrodynamical, supersonic turbulence

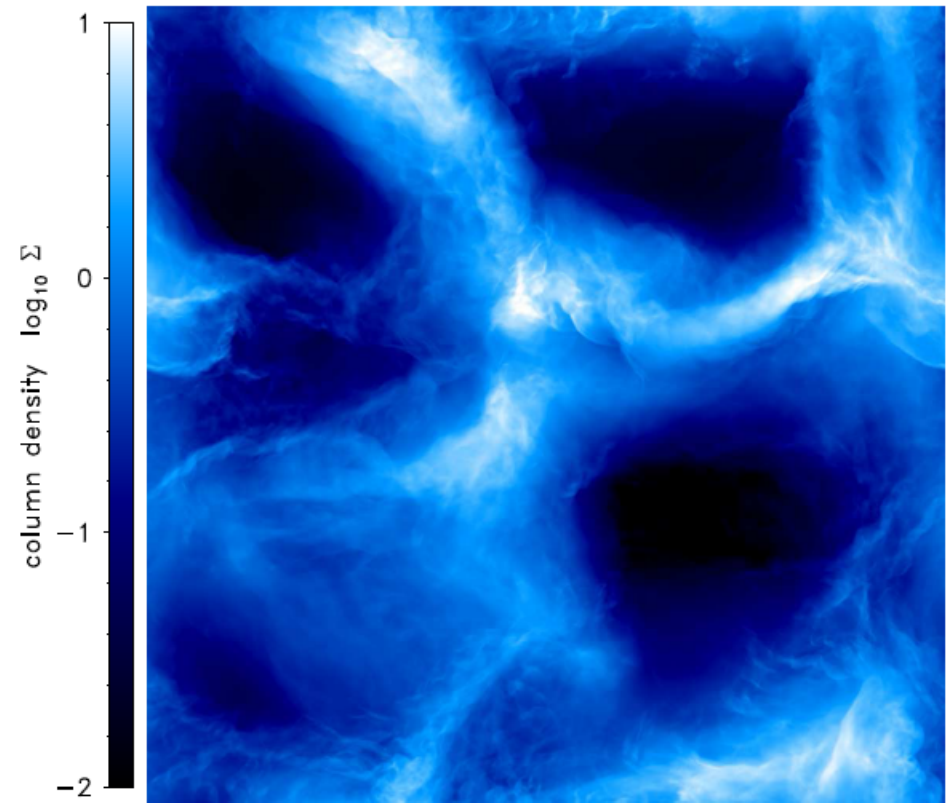
MHD turbulent dynamo (subsonic versus supersonic)

Turbulence forcing – solenoidal versus compressive

Solenoidal forcing



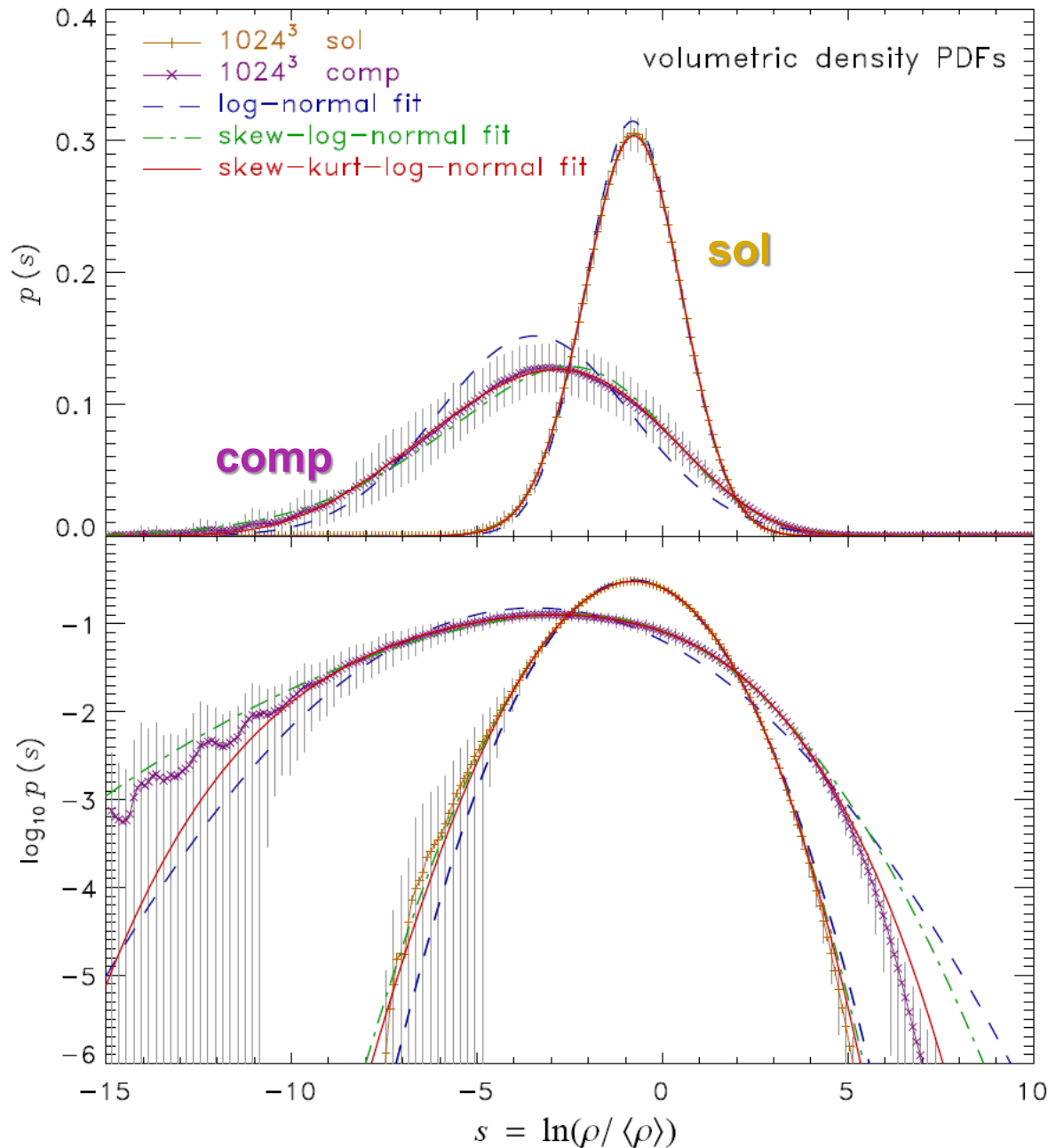
Compressive forcing



Compressive forcing yields 3 times larger density dispersion for the same Mach number

Federrath et al. (2008, 2009, 2010)

Turbulence forcing – solenoidal versus compressive



gas density PDF

PDFs are close to log-normals:

$$p_s ds = \frac{1}{\sqrt{2\pi\sigma_s^2}} \exp\left[-\frac{(s - \langle s \rangle)^2}{2\sigma_s^2}\right] ds$$

$$s = \ln(\rho / \langle \rho \rangle)$$

$$\sigma_\rho / \rho_0 = b\mathcal{M}$$



b = 1/3 (sol)

b = 1 (comp)

Federrath et al.(2008, 2010)

Turbulence forcing – solenoidal versus compressive

Density PDF is key for star formation

- **CMF / IMF** (Padoan & Nordlund 02, Hennebelle & Chabrier 08,09)
- **Star formation efficiency** (Elmegreen 08)
- **Kennicutt-Schmidt relation** (Elmegreen 02, Krumholz & McKee 05, Tassis 07)
- **Star formation rate** (Krumholz & McKee 05, Padoan & Nordlund 11)

All rely on integrals over the turbulent density PDF !

$$\text{SFR}_{\text{ff}} = \frac{\epsilon_{\text{core}}}{\phi_t} \int_{x_{\text{crit}}}^{\infty} xp(x) dx$$

e.g., Krumholz & McKee 2005, Hennebelle & Chabrier 2011

Turbulence forcing – solenoidal versus compressive

scaling relations,

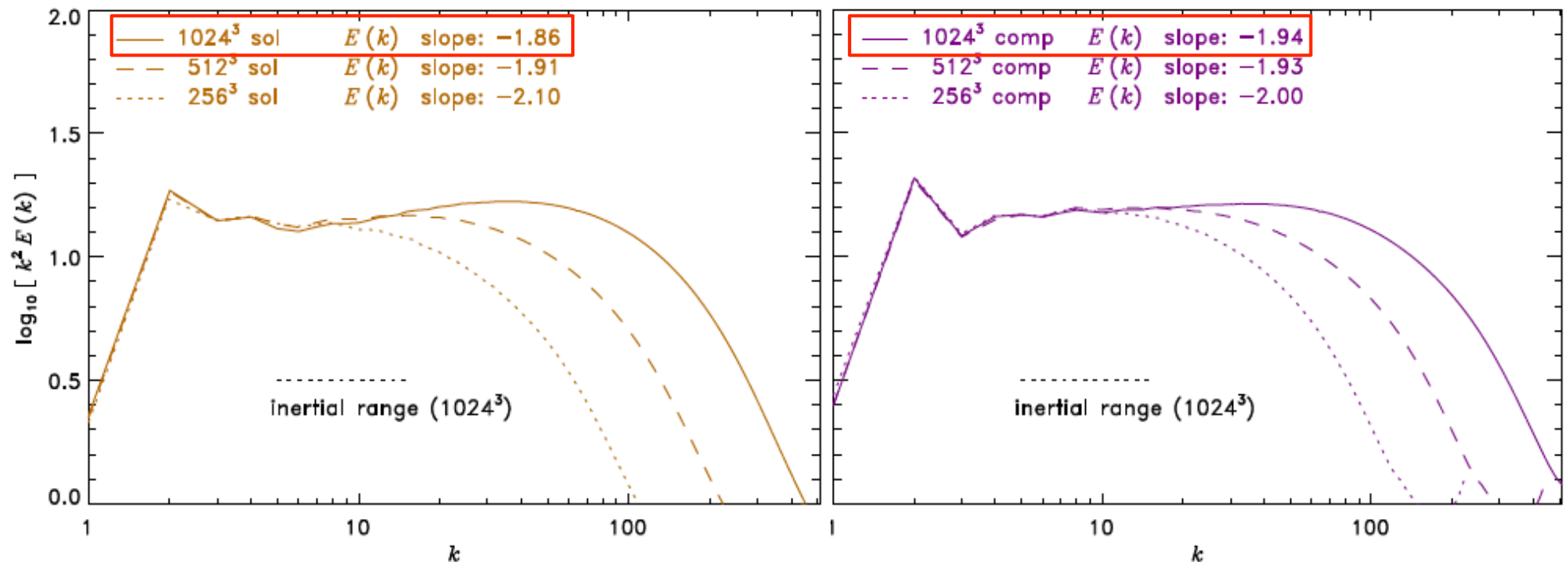
e.g., Larson law of turbulent clouds in the ISM:

$$P(k) \sim k^{-1.8-2.0}$$

(see e.g., Heyer & Brunt 2004)

Solenoidal forcing

Compressive forcing

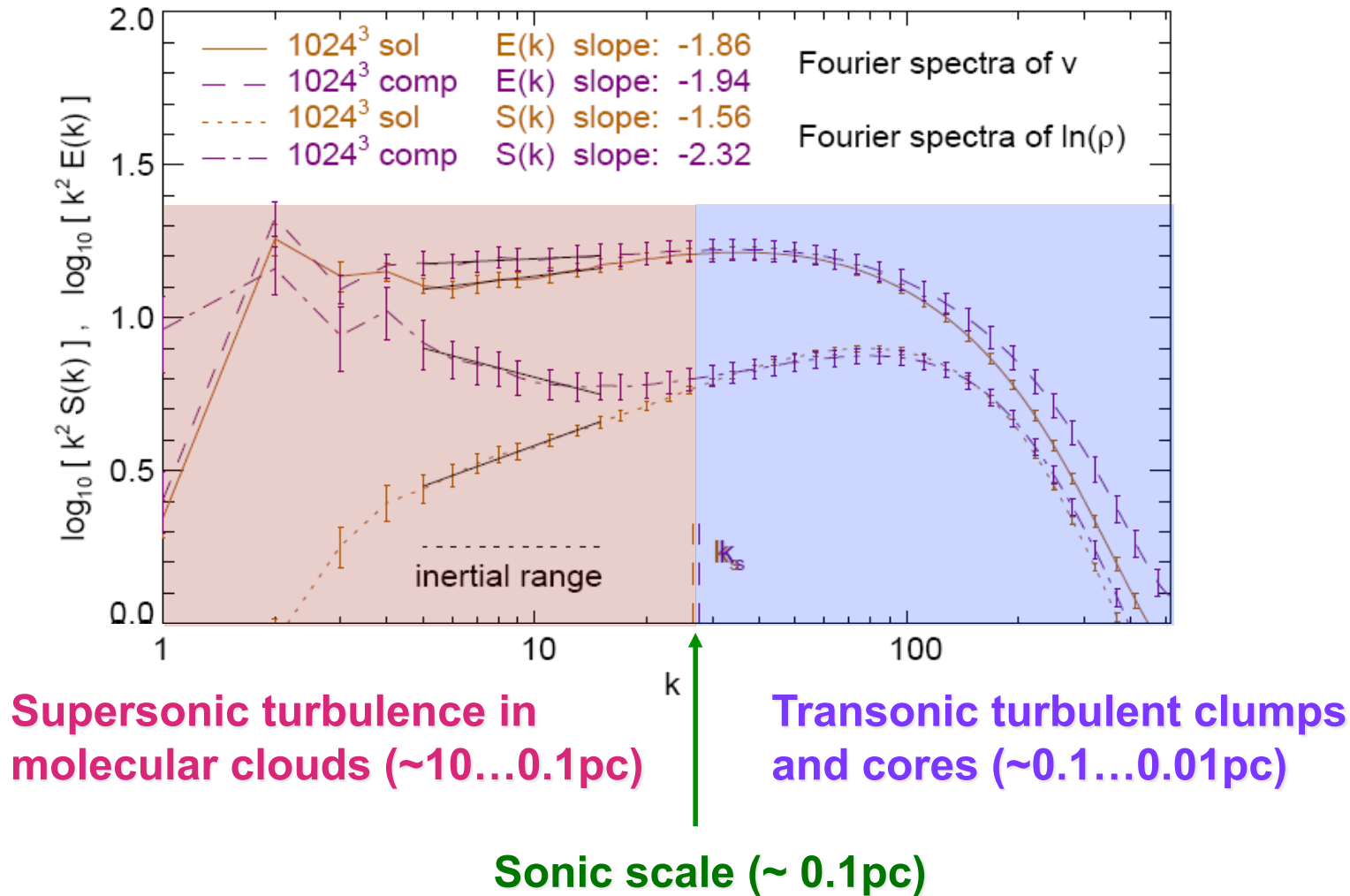


Federrath, Roman-Duval, Klessen, Schmidt & Mac Low, 2010, A&A 512, A81

Turbulence forcing – solenoidal versus compressive

Spatial correlations in the turbulent field (FTs, SFs, Δ -var, PCA)

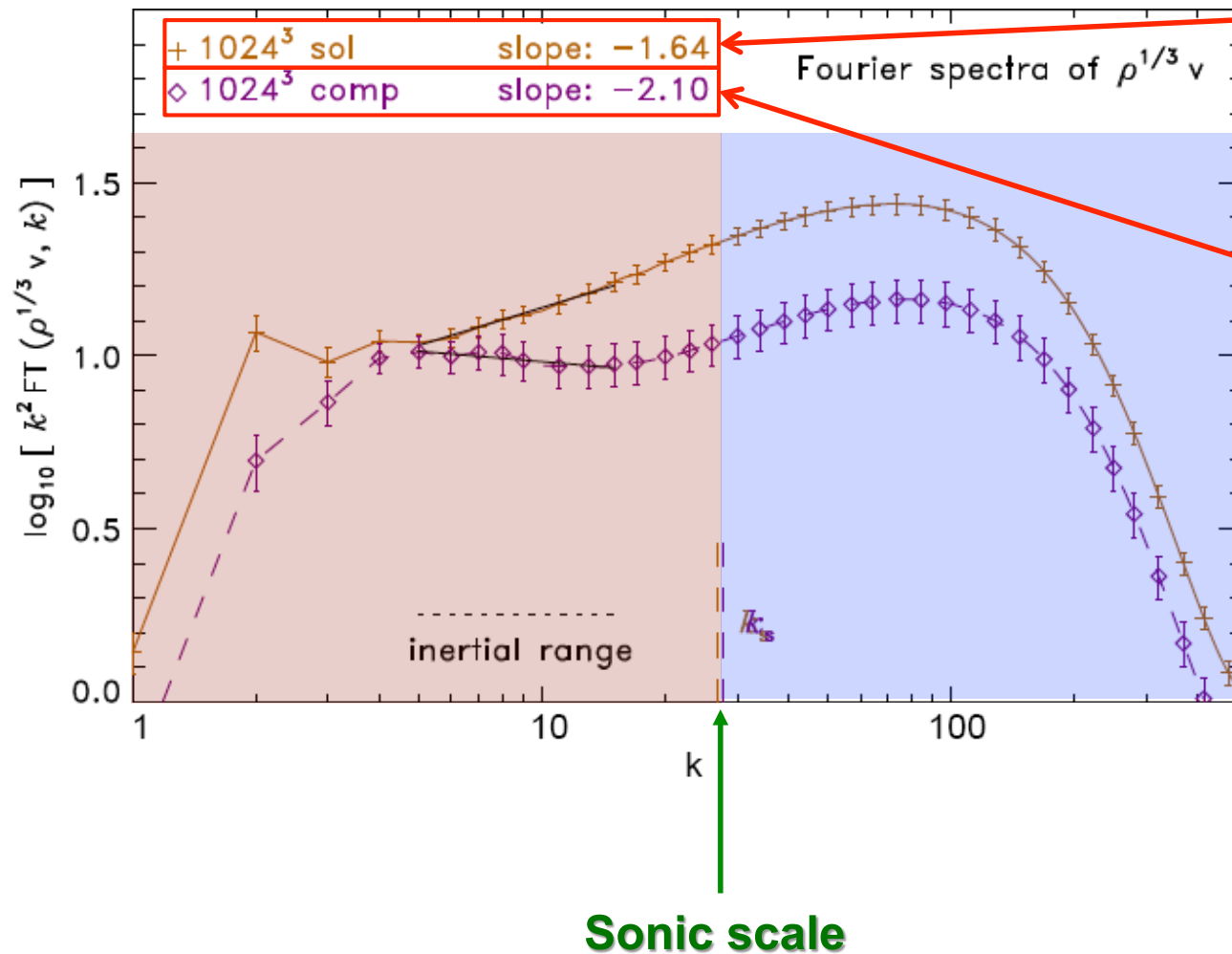
are consistent with observations (e.g., Larson 1981, Solomon+ 1987, Falgarone 1992, Heyer+ 2004)



Turbulence forcing – solenoidal versus compressive

$\rho^{1/3} v$ - spectrum

(suggested by Kritsuk et al. 2007;
see also Kowal et al. 2008; Schmidt et al. 2008; Aluie 2011)



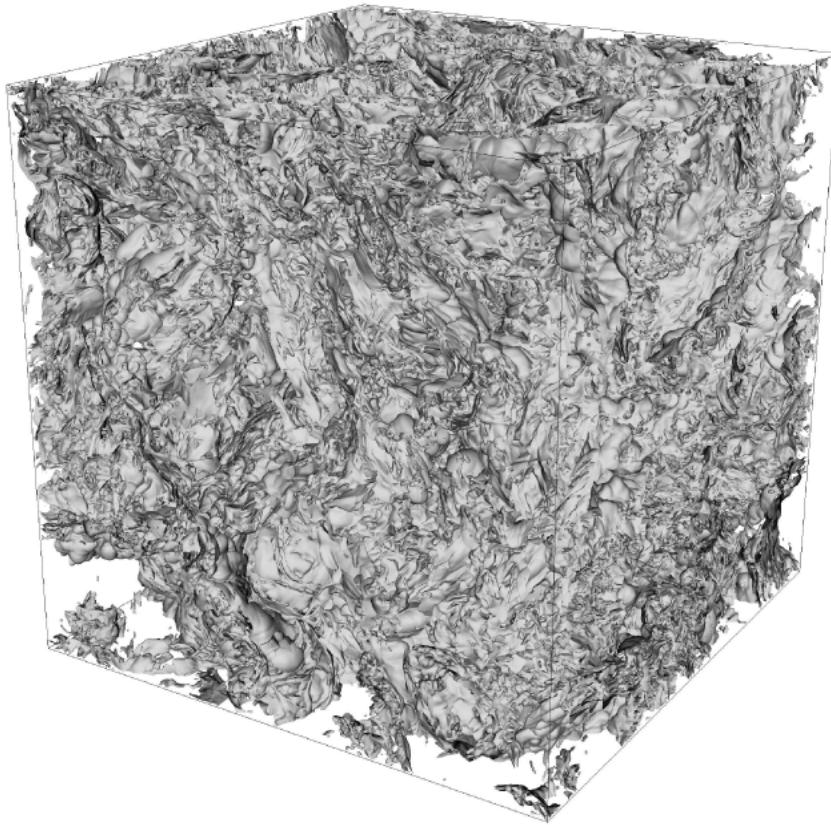
$\sim k^{-5/3}$
K41

$\sim k^{-19/9}$
predicted by
Galtier &
Banerjee
2011, PRL

On the fractal dimension of supersonic turbulence

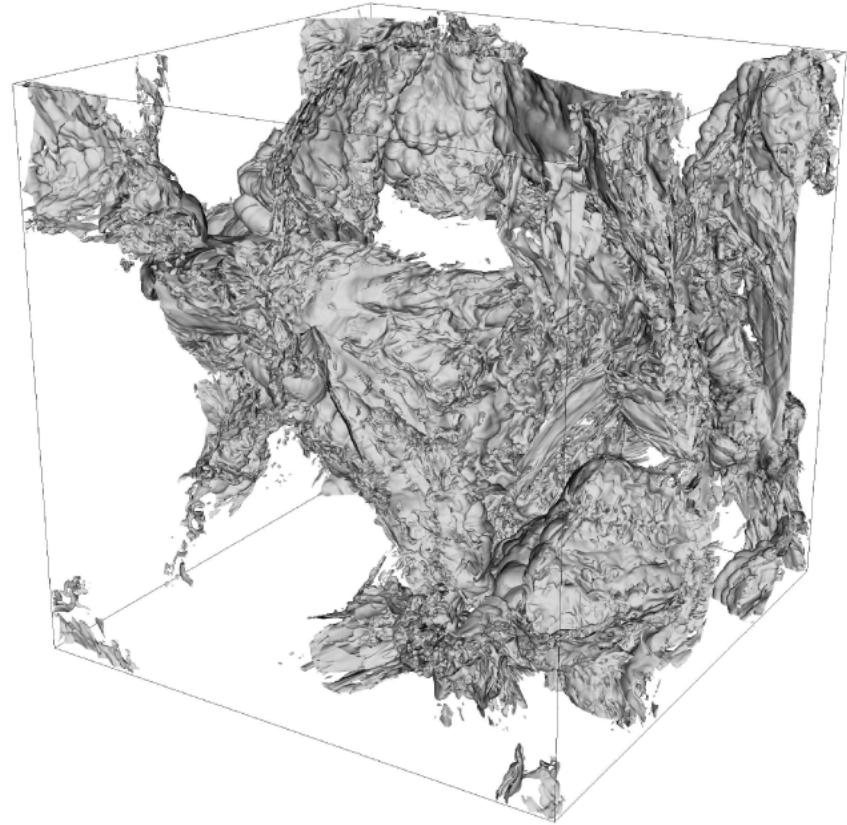
Fractal structures by box-counting, Δ -variance, perimeter-area methods:

Solenoidal forcing



$D = 2.6 \pm 0.1$

Compressive forcing



$D = 2.3 \pm 0.1$

Turbulence driving experiments

Purely hydrodynamical, supersonic turbulence

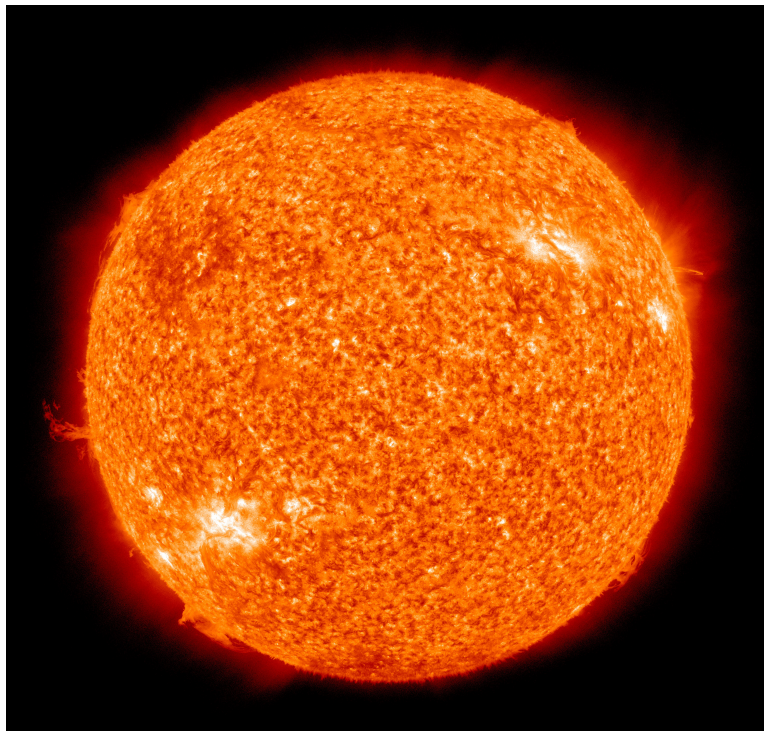
MHD turbulent dynamo (subsonic versus supersonic)

Turbulent dynamo – solenoidal versus compressive

Motivation:

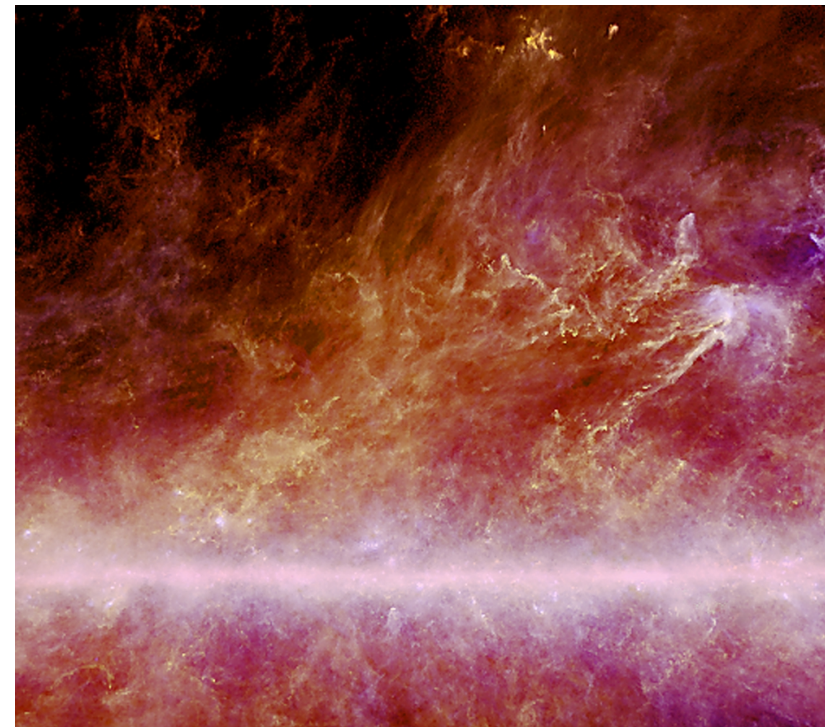
- ◆ Mach number dependence of turbulent dynamo unknown
- ◆ Effects of compression unclear
- ◆ Fundamental difference between

subsonic flows



NASA/SDO (AIA)

supersonic flows



ESA/HFI, IRAS 2010

Turbulent dynamo – solenoidal versus compressive

- ◆ Use idealized controllable turbulence box simulations
- ◆ purely solenoidal & purely compressive forcing
- ◆ Turbulence with Mach numbers in the range 0.02 – 20

Turbulent dynamo – solenoidal versus compressive

MHD equations solved on $128^3 - 512^3$ grid cells (FLASH v2.5, v4)

$$\partial_t \rho + \nabla \cdot (\rho \mathbf{u}) = 0,$$

$$\partial_t (\rho \mathbf{u}) + \nabla \cdot (\rho \mathbf{u} \otimes \mathbf{u} - \mathbf{B} \otimes \mathbf{B}) + \nabla p_\star = \nabla \cdot (2\nu \rho \mathcal{S}) + \rho \mathbf{F},$$

$$\begin{aligned} \partial_t E + \nabla \cdot [(E + p_\star) \mathbf{u} - (\mathbf{B} \cdot \mathbf{u}) \mathbf{B}] = \\ \nabla \cdot [2\nu \rho \mathbf{u} \cdot \mathcal{S} + \mathbf{B} \times (\eta \nabla \times \mathbf{B})], \end{aligned}$$

$$\partial_t \mathbf{B} = \nabla \times (\mathbf{u} \times \mathbf{B}) + \eta \nabla^2 \mathbf{B},$$

$$\nabla \cdot \mathbf{B} = 0,$$

Reynolds number, magnetic Prandtl number:

$$\text{Re} \approx 1500 \quad \text{Pm} \approx 2$$

Turbulent dynamo – solenoidal versus compressive

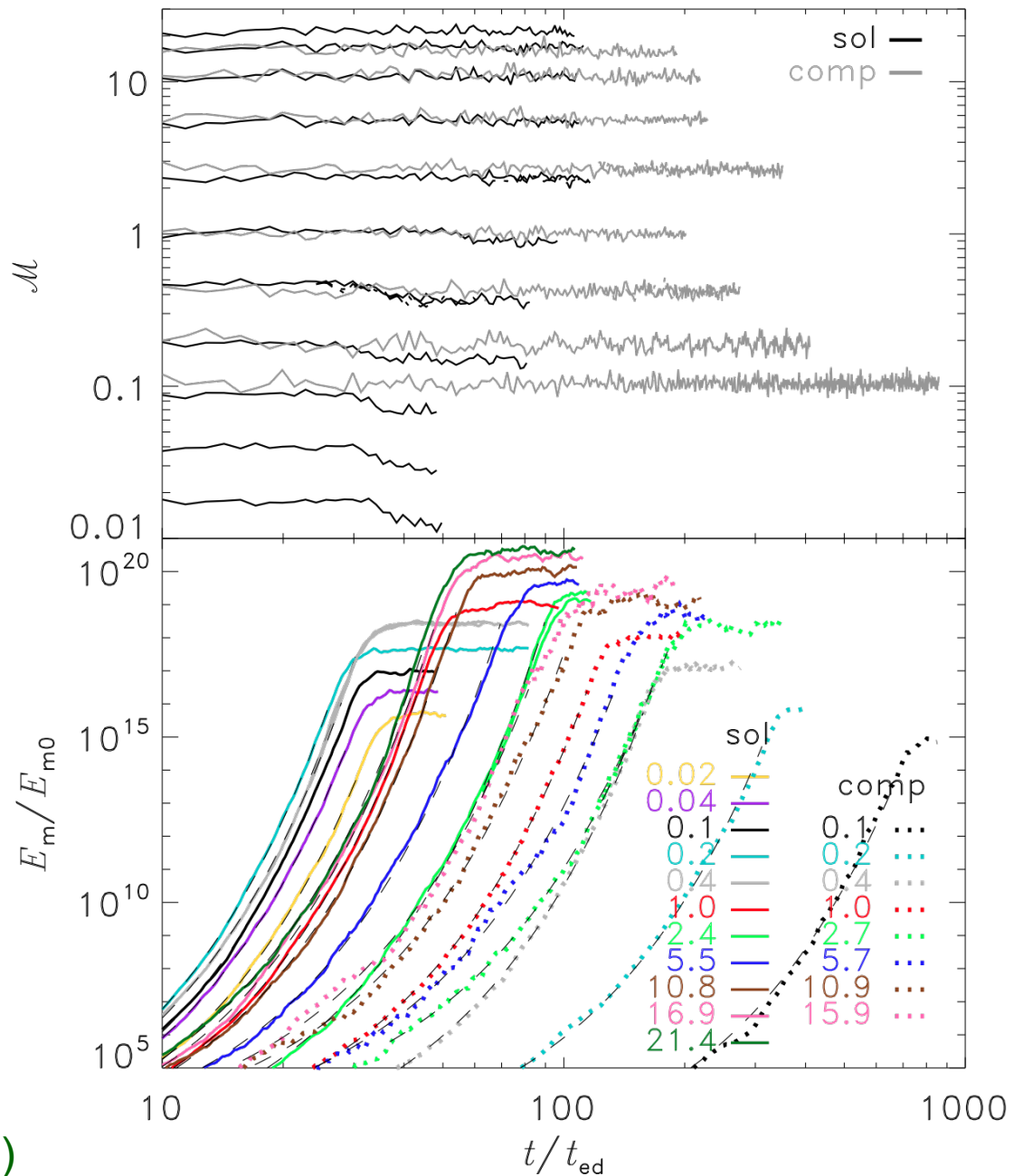
Time evolution

Mach number:

\mathcal{M}

Magnetic energy:

E_m/E_{m0}

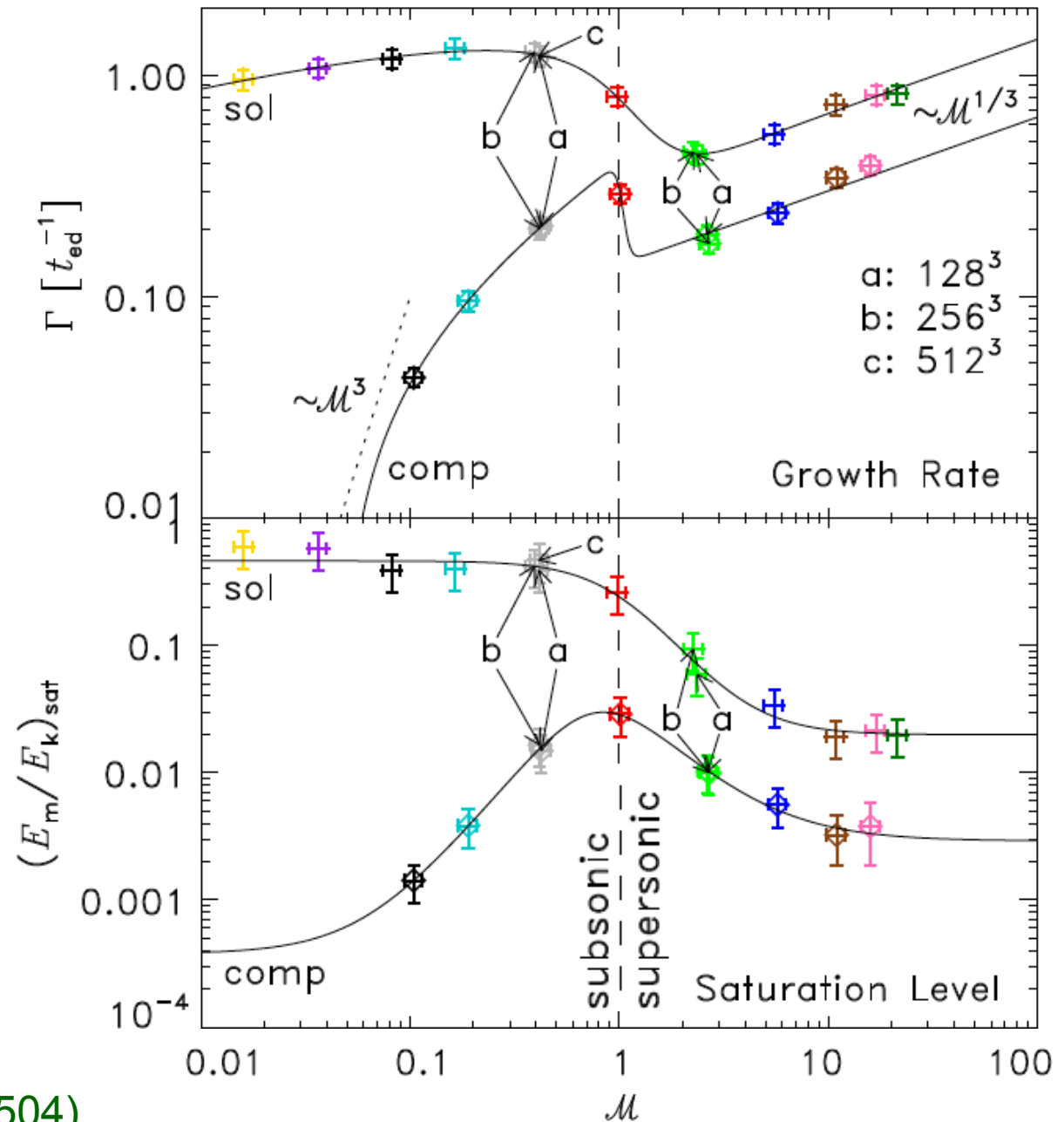


Turbulent dynamo – solenoidal versus compressive

Growth rates & saturation levels as function of Mach

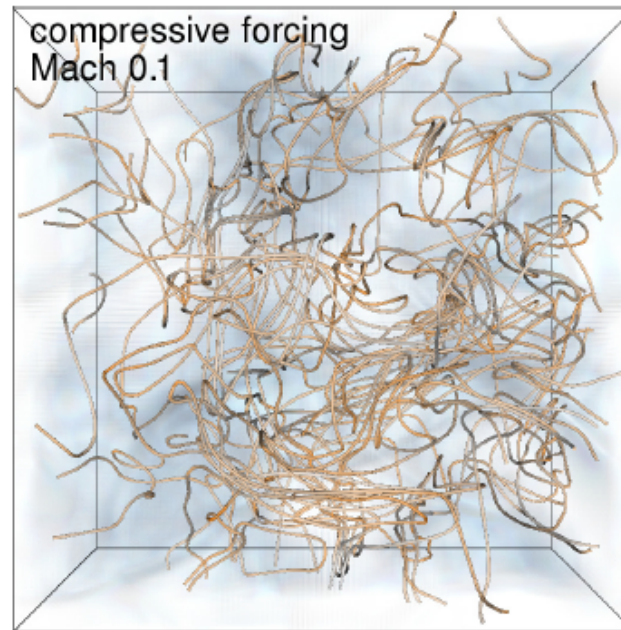
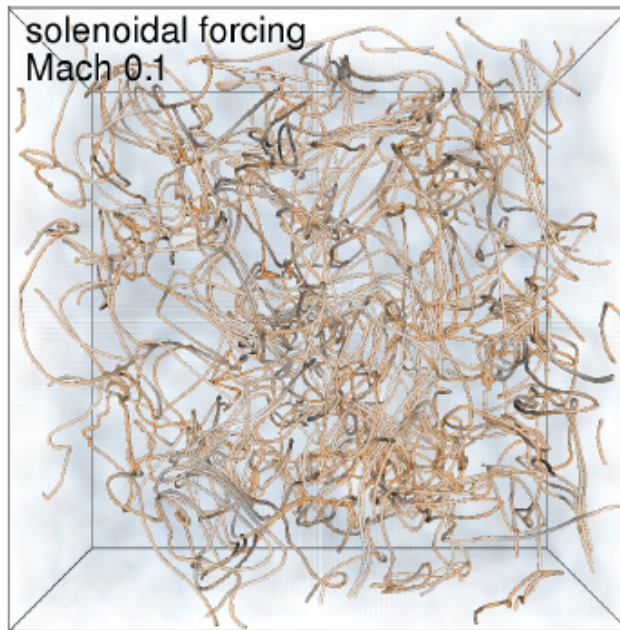
Growth rate:

Saturation level:
($E_{\text{mag}}/E_{\text{kin}})_{\text{sat}}$)



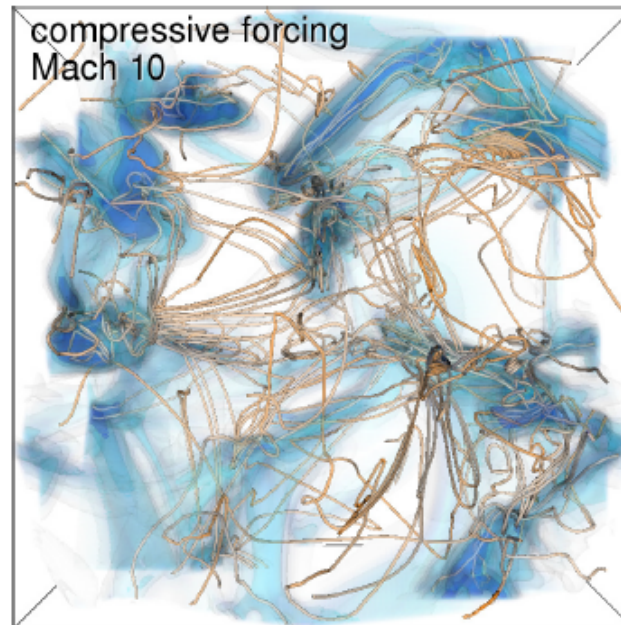
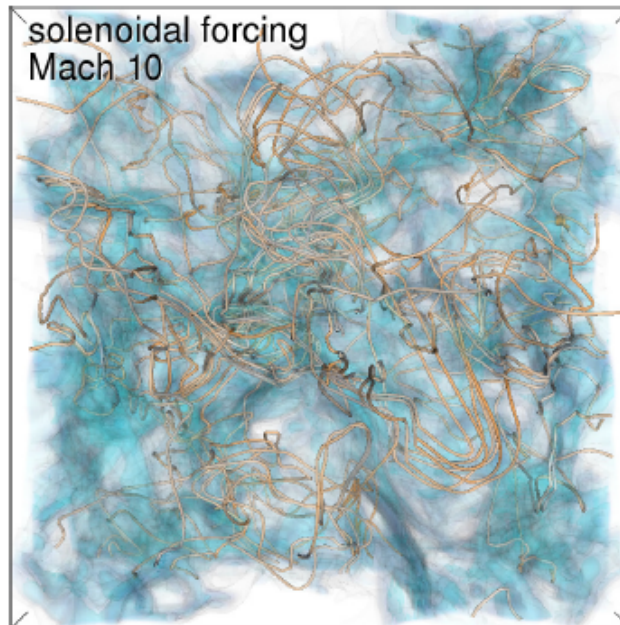
Turbulent dynamo – solenoidal versus compressive

**Mach 0.1
sol**



**Mach 0.1
comp**

**Mach 10
sol**



**Mach 10
comp**

Turbulent dynamo – solenoidal versus compressive

Generation of vorticity (see also Mee & Brandenburg 2006; Del Sordo & Brandenburg 2011):

$$\partial_t \boldsymbol{\omega} = \underbrace{\nabla \times (\mathbf{u} \times \boldsymbol{\omega})}_{\text{non-linear term (amplification!)}} + \underbrace{\nu \nabla^2 \boldsymbol{\omega}}_{\text{diffusion}} + \underbrace{\frac{1}{\rho^2} \nabla \rho \times \nabla p}_{\text{baroclinic term (zero for isothermal)}} + \underbrace{2\nu \nabla \times (\mathcal{S} \nabla \ln \rho)}_{\text{„anti-diffusion“ (only for compressible turbulence)}} + \underbrace{\nabla \times \mathbf{F}}_{\text{forcing term (zero for compressive forcing)}}$$

Strain tensor:

$$\mathcal{S}_{ij} = (1/2)(\partial_i u_j + \partial_j u_i) - (1/3)\delta_{ij} \nabla \cdot \mathbf{u}$$

Reynolds number in our simulations:

$$\text{Re} \approx 1500$$

Enough for the non-linear term to amplify small seeds generated by „anti-diffusion“ term

Turbulent dynamo – solenoidal versus compressive

Extending small-scale dynamo theory to compressive flows
(solving Kazantsev 1967 equation; see Subramanian 1999):

$$\frac{\partial M_L}{\partial t} = 2\kappa_{\text{diff}} M_L'' + 2 \left(\frac{4\kappa_{\text{diff}}}{r} + \kappa_{\text{diff}}' \right) M_L' + \frac{4}{r} \left(\frac{T_N}{r} - \frac{T_L}{r} - T_N' - T_L' \right) M_L$$

ansatz for solution:

$$\kappa_{\text{diff}}(r) = \eta + T_L(0) - T_L(r).$$

$$M_L(r, t) \equiv \frac{1}{r^2 \sqrt{\kappa_{\text{diff}}}} \psi(r) e^{2\Gamma t}.$$

ansatz for velocity correlations:

$$v(\ell) \propto \ell^\vartheta$$

$$T_L(r) = \frac{VL}{3} \left(1 - (r/L)^{\vartheta+1} \right)$$

Kolmogorov (theta=1/3): $T_N^K(r) = \frac{VL}{3} \left(1 - \frac{5}{3} \left(\frac{r}{L} \right)^{4/3} \right)$

Burgers (theta=1/2): $T_N^B(r) = \frac{VL}{3} \left(1 - \frac{2}{5} \left(\frac{r}{L} \right)^{3/2} \right)$

Turbulent dynamo – solenoidal versus compressive

Extending small-scale dynamo theory to compressive flows (solving Kazantsev 1967 equation; see Subramanian 1999):

ansatz for velocity correlations:

$$v(l) \propto l^\vartheta$$

Kolmogorov (theta=1/3):

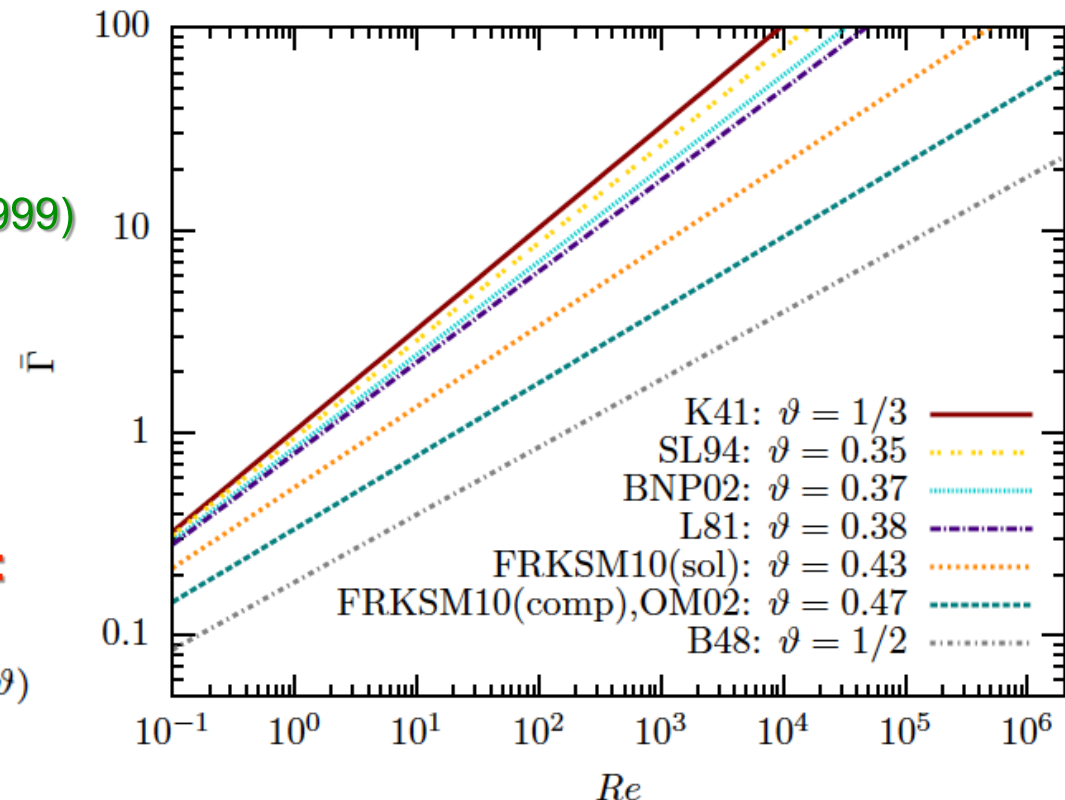
$$\Gamma^K \propto Re^{1/2} \quad (\text{see Subramanian 1999})$$

Burgers (theta=1/2):

$$\Gamma^B \propto Re^{1/3}$$

Growth rate as function of theta:

$$\Gamma = \frac{(163 - 304\vartheta) V}{60 L} Re^{(1-\vartheta)/(1+\vartheta)}$$



Compared turbulence statistics
for **solenoidal (divergence-free)** and **compressive (curl-free)** forcing

- Strong influence on gas **density statistics** (PDF → star formation)
- **Spectra** steeper than Kolmogorov
- **Sonic scale** as characteristic scale in compressible, supersonic turbulence
- **Dynamo growth rate and saturation level depend strongly on Mach and forcing**
- **Vorticity generation** is extremely different, in particular for subsonic flows

 **The forcing matters!**