

Proton acceleration and emission in galaxy clusters

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5th Korean Astrophysics Workshop

November 20th, 2009

Outline

- Galaxy Clusters as UHE-CR sources.
- A theoretical model for acceleration: diffusive shock acceleration.
- Energy losses and the maximum energy: including energy losses into shock acceleration self-consistently.
- ▶ p-γ interactions: primary protons → secondary electrons → non-thermal radiation.
- Conclusions.

Galaxy Clusters As Non-thermal Sources

Radio: ~1/3 of rich clusters (M ~ $10^{15} M_{\odot}$) synchrotron halos (e.g. Govoni & Feretti, 2004).

Non-thermal X-ray diffuse: excess at the level 10⁻¹¹ erg cm⁻²s⁻¹ in the range 10-100 keV (e.g. Fusco-Femiano, 1999; Eckert et al., 2008).

Gamma-rays: no detection \rightarrow upper limits on CR content (Reimer et al., 2003; Aharonian et al., 2009; Aleksic et al., 2009).



Radio + X-ray (assuming same population of electrons): B ~ 0.1-1 μ G Gamma: constraints on CR energy budget ~ few % thermal. For Coma $E_{_{CR}} \leq 10^{62}$ erg.

Particle Acceleration in Clusters

The ingredients:

Formation of (strong) accretion shock at the virial radius (Bertschinger, 1985; Kang et al., 1994);

- Magnetic field of order 0.1-1 μG;
- Shock velocity, free fall velocity:

$$v_{s} \sim \sqrt{\frac{2 G M_{cl}}{R_{cl}}} \simeq 2000 \left(\frac{M_{cl}}{10^{15} M_{\odot}} \right)^{1/2} \left(\frac{R_{cl}}{3 Mpc} \right)^{-1/2} \text{km/s}$$
 (Coma-like cluster)

 Acceleration time ~ Hubble time.
 High energy protons accumulate in the structure (Völk et al., 1996; Berezinsky et al., 1997).



A Self-consistent Approach

$$(Vannoni, Gabici, Aharonian; 2009)$$

$$\frac{\partial f(x, p, t)}{\partial t} + u \frac{\partial f(x, p, t)}{\partial x} - \frac{\partial}{\partial x} \left(D \frac{\partial f(x, p, t)}{\partial x} \right) - \frac{p}{3} \frac{\partial u}{\partial x} \frac{\partial f(x, p, t)}{\partial p} + \frac{1}{p^2} \frac{\partial}{\partial p} (p^2 L(x, p) f(x, p, t)) = Q(x, p)$$
loss term
Numerical formulation
kept general:
- protons or electrons
- general form of the loss term
L(x, p)
- general form of the diffusion
coefficient D(x, p)
- time evolution
Simplifying assumptions:
$$Q(x, p) = Q_0 \delta(x) \delta(p - p_0) \quad \text{injection at the shock}$$

$$D(p) = D. p^{\beta} \quad \text{independent of } x, \text{ Bohm: } D_0 = \frac{c^2}{3eB}, \beta = 1$$

$$L(x, p) = L(p) \quad \text{independent of } x$$

$$B_2 = \xi B_1 \quad \xi = 4, \text{ random magnetic field (Alfvén)}$$

 $\xi{=}4\,\text{,}$ random magnetic field (Alfvén waves)

Proton Energy Losses

Assumption: Spherical accretion shock.

pair production



Interaction with the CMB can become the limiting mechanism (Norman et al., 1995; Kang et al., 1997).

Pair production dominant energy loss channel for all realistic sets of parameters.

pion production

Time dependent calculation needed.

Proton Spectrum



non-exponential cut-off

(Vannoni et al., arXiv0910.5715)

No steady state.

Cut-off energy $\sim 7 \times 10^{18}$ eV



Production and Cooling of Secondary Electrons

$$\tau_{acc}=5$$
 Gyr.



Photopair production: electrons production spectra — detailed numerical calculation (Kelner & Aharonian, 2008).

Production and Cooling of Secondary Electrons

Cut-off in KN 1 1 1 1 111 1 1 1 1110 10-14 [erg⁻¹s⁻¹ 비나 10⁻¹⁵ 한 한 도 IC off CMB $B_1 = 0.3 \ \mu G$ $B_2=1 \mu G$ 10-16 F....... 1011 1012 1013 1014 1015 1016 1017 1018

E [eV]

Electrons rapidly cool via synchrotron and IC. IC cooling in Klein-Nishina regime.



Broadband SED, detectability



 Efficient proton acceleration up to 10¹⁸ eV.
 Efficient secondary electron production in pphoton interaction.
 Rapid electron cooling and energy transfer to non-thermal photons.

~100% of the energy lost by protons is found in non-thermal photons.

Broader cut-off, less steep than exponential.

Synchrotron: ~factor 10 enhancement downstream due to higher B. Peak energy ~100 keV.

Inverse Compton: same emission level up and downstream. Peak between 10 and 100 TeV, but absorption will reduce the gamma-ray flux.

Broadband SED, detectability



Shock Modification

If acceleration is efficient, dynamical reaction of particles expected on the shock structure.



Non-linear shock acceleration theory: (Drury & Völk 1981; Kang et al. 1996; Malkov 1997; Berezhko & Ellison 1999; Blasi 2002...)

\$ energy transferred by the shock to accelerated particles
is accumulated at energies close to the cut-off.

\$ shock modification is expected to translate in an enhancement of the non-thermal luminosity of the system.

Approximation: R=7, ξ =7

Proton Spectrum

For R=7 \rightarrow expected energy spectrum \propto E^{-1.5}





Integrated spectrum: all energy transferred at E around the cut-off

Broadband SED (Modified Shocks)



Conclusions

- Self-consistent calculation of particle shock acceleration including energy losses regime — a very general method.
 - spectral shape not predictable a priori,
 - time dependent approach needed.
- Clusters of galaxies can efficiently accelerate CRs.
- Not EHE-CR sources: maximum energy determined by p-CMB interactions, Emax~10¹⁹ eV.
- Pair production is the dominant energy loss channel.
- Secondary electrons cool rapidly via SY and IC (in KN regime). All the energy lost by protons is found in non-thermal radiation
- For Coma-like parameters: X-ray flux ~10⁻¹² erg cm⁻² s⁻¹ at ~100 keV γ-ray flux ~10⁻¹³ erg cm⁻² s⁻¹ at ~10 TeV, absorption at higher energies
- Possibility of enhanced emission if shocks are modified.

Detectable by next generation Cherenkov telescopes?