Saturation of CR-streaming instabilities at supernova shocks

Qinghuan Luo U. Sydney

Overview

- CR-streaming instabilities: kinetic theory
- Saturation mechanisms
 - -feedback effects constrain magnetic field amplification
- Implications for SNR forward shocks

Generation of magnetic fields

- Strong magnetic fields at SNR shocks
- Electromagnetic instabilities due to CR streaming
- Resonant
 - -Affvén wave instability (Skilling 1970, ...)
- Nonresonant instabilities
 - -CR-streaming instability $\omega \rightarrow 0$ (Bell 2004)
 - -Weibel instability $\omega \rightarrow 0$ (Weibel 1959)
 - -Whistler instability (e.g. Gary & Cairns 1999)
 - -Filamentation instability (FI) $\omega \rightarrow 0$ (Bell 2005)
- These instabilities can generate magnetic turbulence

CR-streaming instability

• Kinetic model: background plasma+CR current $J_{CR} = qn_{CR}v_{CR}$

Neutralisation conditions

$$e(n_e - n_p) = qn_{_{\mathrm{CR}}}, \quad e(n_e v_e - n_p v_p) = qn_{_{\mathrm{CR}}} v_{_{\mathrm{CR}}}$$

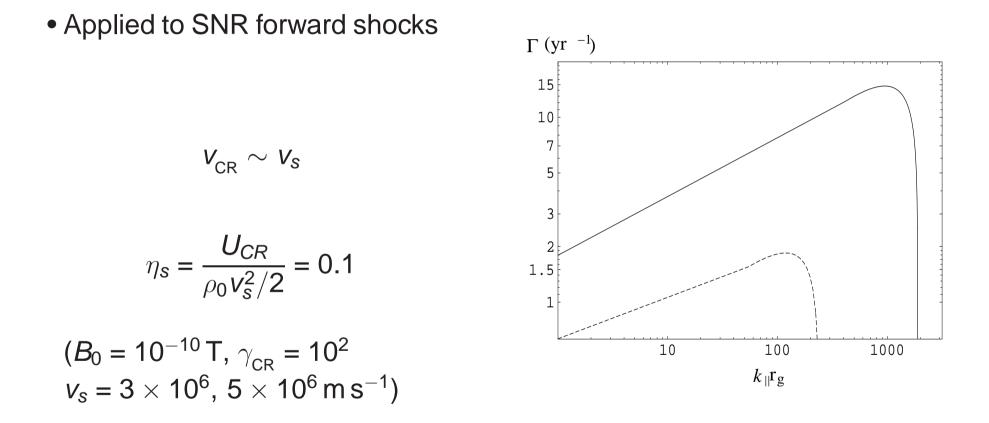
- \rightarrow compensating (or return) current $-J_{CR}$
- Purely growing mode due to the compensating current with condition

$$J_{CR} > rac{k_\parallel B_0}{\mu_0}$$

• Growth rate at $k_{\parallel}r_g \gg 1$

$$\Gamma pprox k_{\parallel} v_A \left(rac{2k_m}{k_{\parallel}}-1
ight)^{1/2}, \quad k_{\parallel} < 2k_m \equiv rac{\mu_0 J_{
m CR}}{B_0}$$

(Bell 2004; Reville et al. 2006; Amato & Blasi 2009; LM09)



Saturation

- Numerical simulations clearly show saturation but predict different *B* -MHD simulation (Bell 2004, 2005) -PIC simulation $B/B_0 \sim 1-30$ (Niemiec et al 2008; Riquelme & Spitkovsky 2009; Stroman et al. 2009)
- Limitation of numerical kinetic models by $n_{_{
 m CR}}/n_{
 m p}\ll 1$
- Determining of saturation important for the relevance of the instability
- Saturation mechanisms
 - -Stabilizing due to feedback effects (LM09)
 - -Magnetisation of CR beam (Medvedev & Loeb 1999)
 - -Magnetic trapping (Davidson 1972; Achterberg et al 2007)

Saturation due to diffusion

- Feedback on CR current treated as quasilinear diffusion (QLD)
 -General formalism (e.g. Shapiro & Shevchenko 1964)
 -App to Weibel instability (Davidson et al 1972)
- Nonresonant/resonant diffusion reduces anisotropy–changes $v_{CR}(t)$
- Nonresonant diffusion, similar to the firehose instability
- Resonant diffusion, pitch angle scattering The case $kr_g \sim$ 1 well discussed (e.g. Melrose 1980)
- For $k_{\parallel}r_{g0} \gg 1$, only small fraction of CRs with $\alpha \sim \pi/2$ strongly affected (LM09)
- Pitch angle scattering dominant in reducing $V_{CR}(t)$

A streaming model

CR distribution

$$f(p,\alpha) = n_{CR} \left(1 + 3\frac{V_{CR}}{v}\cos\alpha\right) \frac{g(p)}{4\pi p^2}$$
$$g(p) = \begin{cases} \frac{b-1}{p_1} \left[1 - \left(\frac{p_1}{p_2}\right)^{b-1}\right]^{-1} \left(\frac{p}{p_1}\right)^{-b}, \ p_1 \le p \le p_2\\ 0 & \text{otherwise} \end{cases}$$

• Streaming velocity $\langle v_{\parallel} \rangle = v_{CR}$

• Evolution of CR treaming motion

$$\frac{dv_{CR}}{dt} = \frac{\int p_{\perp} dp_{\perp} dp_{\parallel} v_{\parallel} (dF/dt)}{\int p_{\perp} dp_{\perp} dp_{\parallel} F}$$

• Scattering time

$$t_{\rm s} \sim \frac{\gamma_0}{\Omega} \left(\frac{B_0}{B}\right)^2 k_0 r_{g0}$$

compared to growth time $1/\Gamma$

• Saturated magnetic field

$$\frac{B}{B_0} \sim \xi (k_0 r_{g0})^{3/4} (k_m r_{g0})^{1/4} \beta_A^{1/2} \le 2^{3/4} \xi k_m r_{g0} \beta_A^{1/2}$$

• For SNR shocks, $r_{g0} \leq R_s$, $k_0 < 2k_m$, $v_A = 10$ km s⁻¹

$$rac{B}{B_0}pprox$$
 16, $R_s=3\,\mathrm{pc},$ k_mpprox 300 $\mathrm{pc}^{-1},$ $\eta_s=0.1$

- Amplification modest, similar to recent PIC simulations (e.g. Riquelme & Spitkovsky 2009; Stroman et al. 2009)
- Prediction of scaling $B \propto v_s^{\alpha}$, α depends on how k_0 evolves

Magnetisation of CR beam

- Deflection due to self-generated *B* (Medvedev & Loeb 1999; Riquelme & Spitkovsky 2009)
- Larmor radius $\sim 1/k_m$
- Saturation

$$\frac{B}{B_0}\approx k_m r_{g0}$$

• Higher than the limit by QLD mechanism

$$\frac{B_m}{B_d} \sim \left(\frac{k_0}{k_m}\right)^{1/4} \left(\frac{c}{v_A}\right)^{1/2} \gg 1$$

Summary

- The nonresonant instability can grow efficiently to nonlinear regime, due to the return current
- Saturation due to diffusion occurs well before magnetisation of CRs
- Application to SNR forward shocks, predicting modest amplification, consistent with the result from recent numerical kinetic models (Riquelme & Spitkovsky 2009; Stroman et al. 2009)
- Analytical result not limited by $n_{\rm CR}/n_p \ll 1$
- Limitation of QLD, but as the feedback mechanism generic it may still provide reasonable estimate