

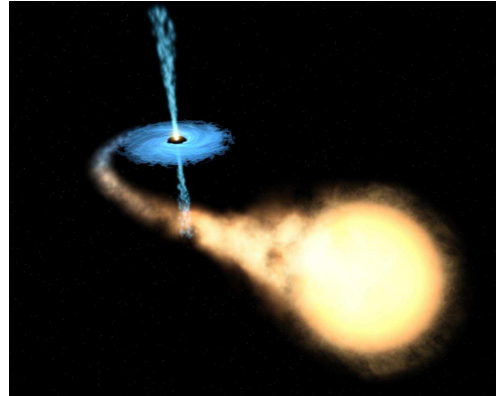
Simulation of relativistic outflows in astrophysics

Indranil Chattopadhyay

indra@canopus.cnu.ac.kr

Dept of Astronomy & Space Science,
Chungnam National University,
Daejeon, South Korea.

Collaborator: Prof. Dongsu Ryu (CNU)



- A short introduction on relativistic flow & Equation of State.
- A brief look on the code.
- Relativistic outflows.
- Concluding Remarks.

Introduction

- We encounter relativistic fluid in relativistic astrophysical jets (galactic/extragalactic), GRBs or say matter falling onto compact objects.

Introduction

- We encounter relativistic fluid in relativistic astrophysical jets (galactic/extragalactic), GRBs or say matter falling onto compact objects.
- A fluid can be relativistic on account of (a) Bulk velocity ($v \sim c$) or (b) temperature (thermal energy \gtrsim rest energy).

Introduction

- We encounter relativistic fluid in relativistic astrophysical jets (galactic/extragalactic), GRBs or say matter falling onto compact objects.
- A fluid can be relativistic on account of (a) Bulk velocity ($v \sim c$) or (b) temperature (thermal energy \gtrsim rest energy).
- Examples: (a) Accretion on to BH: At $r \rightarrow$ large, $v \sim 0$ & $T_p \sim$ small ; At $r \sim$ few $10r_g$, $v \gtrsim 0.1c$, $T_p \sim 10^{12} K$; The adiabatic index (c_p/c_v) $\gamma \rightarrow 5/3 \rightarrow 4/3$.

Introduction

- We encounter relativistic fluid in relativistic astrophysical jets (galactic/extragalactic), GRBs or say matter falling onto compact objects.
- A fluid can be relativistic on account of (a) Bulk velocity ($v \sim c$) or (b) temperature (thermal energy \gtrsim rest energy).
- Examples: (a) Accretion on to BH: At $r \rightarrow$ large, $v \sim 0$ & $T_p \sim$ small ; At $r \sim$ few $10r_g$, $v \gtrsim 0.1c$, $T_p \sim 10^{12} K$; The adiabatic index (c_p/c_v) $\gamma \rightarrow 5/3 \rightarrow 4/3$.
- (b) Jets: Originates at $r \sim$ few $10r_g$ with $v \sim 0$, and $T_p \sim 10^{12} K$ ($\gamma \sim 4/3$) ; at $r \sim 1000r_g$ $v \gtrsim 0.9c$, $T_p \sim$ moderate ($5/3 \gtrsim \gamma \gtrsim 4/3$) : As the jet hits ambient medium, $v \sim$ small, but T_p increases ($\gamma \sim 4/3$).

Introduction

- We encounter relativistic fluid in relativistic astrophysical jets (galactic/extra galactic), GRBs or say matter falling onto compact objects.
- A fluid can be relativistic on account of (a) Bulk velocity ($v \sim c$) or (b) temperature (thermal energy \gtrsim rest energy).
- Examples: (a) Accretion on to BH: At $r \rightarrow$ large, $v \sim 0$ & $T_p \sim$ small ; At $r \sim$ few $10r_g$, $v \gtrsim 0.1c$, $T_p \sim 10^{12} K$; The adiabatic index (c_p/c_v) $\gamma \rightarrow 5/3 \rightarrow 4/3$.
- (b) Jets: Originates at $r \sim$ few $10r_g$ with $v \sim 0$, and $T_p \sim 10^{12} K$ ($\gamma \sim 4/3$) ; at $r \sim 1000r_g$ $v \gtrsim 0.9c$, $T_p \sim$ moderate ($5/3 \gtrsim \gamma \gtrsim 4/3$) : As the jet hits ambient medium, $v \sim$ small, but T_p increases ($\gamma \sim 4/3$).
- It is important to use correct γ , influences the value of $c_s \Rightarrow$ would affect T distribution, as well as structure propagations: Are we equipped to do so?

A brief account on EoS

- The equations of motion are ($c = 1$):

$$(\rho u^\nu)_{;\nu} = 0, \quad T^{\mu\nu}_{;\nu} = 0$$

ρ , u^μ are proper density and 4-velocity, $T^{\mu\nu} = \rho h u^\mu u^\nu + p g^{\mu\nu}$; h = specific enthalpy = $(e + p)/\rho$, e total proper energy density and p pressure.

- Number of eqs = 5; Number of variables = 6.

A brief account on EoS

- To resolve, one assumes an EoS *i.e.*, $h \equiv h(p, \rho)$ or $e \equiv e(p, \rho)$.
- In non-relativistic regime the internal energy density is given by

$$e_{int} = \frac{p}{(\gamma - 1)}$$

- the most favoured EoS in RHD is (hereafter ID) is

$$h = 1 + \frac{\gamma \Theta}{\gamma - 1}, \quad \text{or} \quad e = \rho + \frac{p}{\gamma - 1}$$

$\Theta = p/\rho$ a temperature like variable.

A brief account on EoS

- The general definition of polytropic index and sound speed:

$$n = \rho \frac{\partial h}{\partial p} - 1, \quad c_s^2 = -\frac{\rho}{nh} \frac{\partial h}{\partial \rho}.$$

A brief account on EoS

- The general definition of polytropic index and sound speed:

$$n = \rho \frac{\partial h}{\partial p} - 1, \quad c_s^2 = -\frac{\rho}{nh} \frac{\partial h}{\partial \rho}.$$

- For ID, n & c_s :

$$n = \frac{1}{\gamma - 1}, \quad c_s^2 = \frac{\gamma \Theta (\gamma - 1)}{\gamma \Theta + \gamma - 1},$$

$\Theta \rightarrow \infty$, and for $\gamma > 2$, $c_s > 1$!!!

A brief account on EoS

- The general definition of polytropic index and sound speed:

$$n = \rho \frac{\partial h}{\partial p} - 1, \quad c_s^2 = -\frac{\rho}{nh} \frac{\partial h}{\partial \rho}.$$

- For ID, n & c_s :

$$n = \frac{1}{\gamma - 1}, \quad c_s^2 = \frac{\gamma \Theta (\gamma - 1)}{\gamma \Theta + \gamma - 1},$$

$\Theta \rightarrow \infty$, and for $\gamma > 2$, $c_s > 1$!!!

- Not surprising! ID is hijacked from Newtonian physics!

A brief account on EoS

- The general definition of polytropic index and sound speed:

$$n = \rho \frac{\partial h}{\partial p} - 1, \quad c_s^2 = -\frac{\rho}{nh} \frac{\partial h}{\partial \rho}.$$

- For ID, n & c_s :

$$n = \frac{1}{\gamma - 1}, \quad c_s^2 = \frac{\gamma \Theta (\gamma - 1)}{\gamma \Theta + \gamma - 1},$$

$\Theta \rightarrow \infty$, and for $\gamma > 2$, $c_s > 1$!!!

- Not surprising! ID is hijacked from Newtonian physics!
- (a) In ID γ or n has no temperature dependence
- (b) Gives unphysical wave speeds !

A brief account on EoS

- The general definition of polytropic index and sound speed:

$$n = \rho \frac{\partial h}{\partial p} - 1, \quad c_s^2 = -\frac{\rho}{nh} \frac{\partial h}{\partial \rho}.$$

- For ID, n & c_s :

$$n = \frac{1}{\gamma - 1}, \quad c_s^2 = \frac{\gamma \Theta (\gamma - 1)}{\gamma \Theta + \gamma - 1},$$

$\Theta \rightarrow \infty$, and for $\gamma > 2$, $c_s > 1$!!!

- Not surprising! ID is hijacked from Newtonian physics!
- (a) In ID γ or n has no temperature dependence
- (b) Gives unphysical wave speeds !
- Are we free to use any other arbitrary EoS?

A brief account on EoS

- Relativistic kinetic theory imposes a strong constrain on EoS, known as Taub's inequality,

$$(h - \Theta)(h - 4\Theta) \geq 1$$

- ID do not satisfy TAUB's inequality (TI) for all values of Θ and $\gamma \Rightarrow$ ID is relativistically wrong!

A brief account on EoS

- Chandrashekhar (1934), Synge (1957) calculated the exact EoS for simple gas (RP)

$$h = \frac{K_3(1/\Theta)}{K_2(1/\Theta)},$$

K_2 and K_3 are modified Bessels function of second kind & of order 2 & 3.

- This accuracy comes at expense of extra computational cost as shown by Falle & Komissarov (1996)! \Rightarrow Not suitable for Numerical RHD!
- What do we do? Try and propose a new equation of state

- New EoS must satisfy:
 - (a) Taub's Inequality for all Θ !
 - (b) Should give $\gamma \rightarrow 5/3$ (or $n \rightarrow 3/2$) for $\Theta \rightarrow 0$ & $\gamma \rightarrow 4/3$ (or $n \rightarrow 3$) $\Theta \rightarrow \infty$!
 - (c) Good fit to RP!

- Our proposed EoS (RC) is,

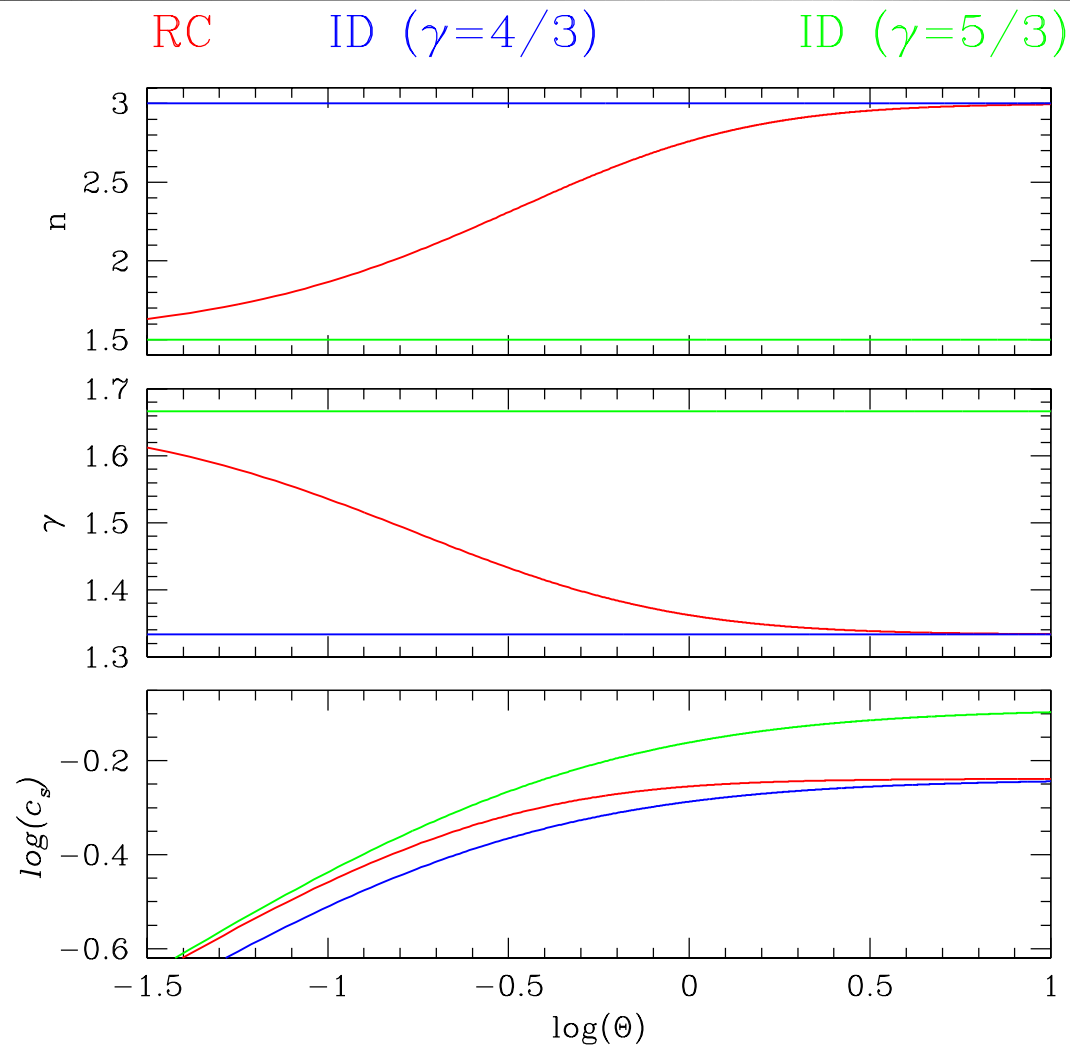
$$\frac{p}{e - \rho} = \frac{3p + 2\rho}{9p + 3\rho}; \quad \text{or} \quad h = 2 \frac{6\Theta^2 + 4\Theta + 1}{3\Theta + 2}. \quad (iv)$$

- Expression of n & c_s with RC EoS:

$$n = 3 \frac{9\Theta^2 + 12\Theta + 2}{(3\Theta + 2)^2}, \quad c_s^2 = \frac{\Theta(3\Theta + 2)(18\Theta^2 + 24\Theta + 5)}{3(6\Theta^2 + 4\Theta + 1)(9\Theta^2 + 12\Theta + 2)}.$$

- RC is extremely accurate!

$$\frac{|h_{\text{RP}} - h_{\text{RC}}|}{h_{\text{RP}}} \lesssim 0.8\%$$



- Comparing various quantities top to bottom (i) n , (ii) γ , (iii) c_s ; for RC, ID($\gamma = 4/3$), ID($\gamma = 5/3$).

Highlights of the new code

- We have implemented RC and developed a new code (based on TVD scheme)!!
- The full eigen structure has been derived
- The **Jacobian matrix and the eigen structure is derived without assuming any exact form of h (EoS)**
- This **makes the code ideal to compare** the effect of **various EoS**
- For the details of the code and its test runs please see [Ryu, Chattopadhyay, Choi \(2006\)](#)

Highlights of the new code

- The equations of motions in conserved form $[(\rho u^\nu)_{;\nu} = 0, T_{;\nu}^{\mu\nu} = 0]$:

$$\frac{\partial \vec{q}}{\partial t} + \frac{\partial \vec{F}^j}{\partial x^j} = 0, \quad (v)$$

- The conserved quantities are:

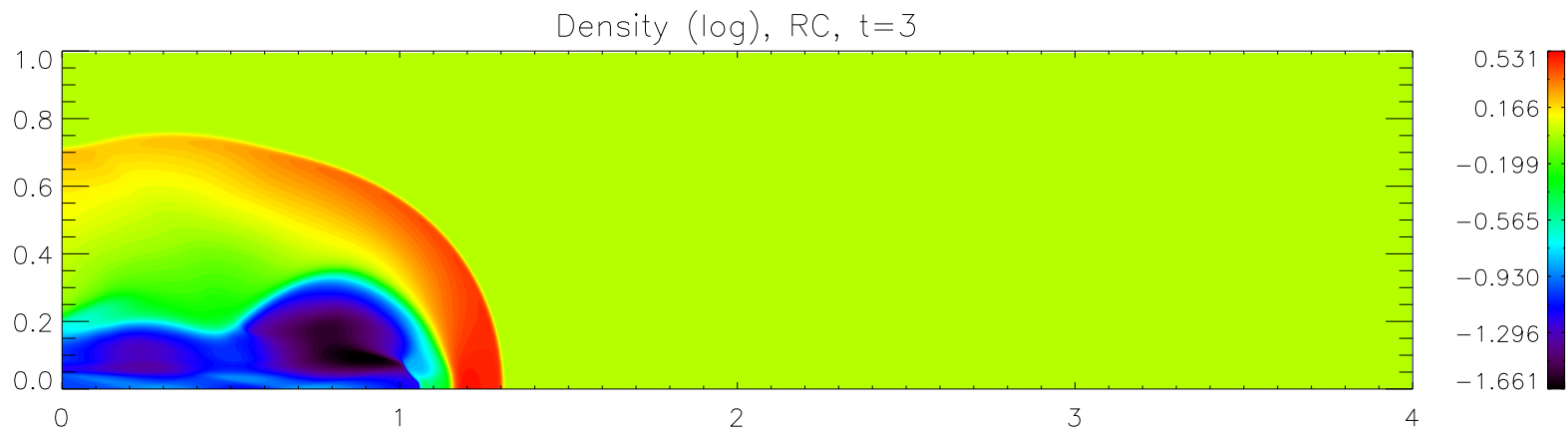
$$\vec{q} = [D \quad M^i \quad E]^T, \quad \vec{F}^j = [Dv^j \quad M^i v^j + p\delta^{ij} \quad (E + p)v^j]^T.$$

- And the transformation relation betⁿ observer frame to fluid rest frame are:

$$D = \Gamma\rho, \quad M^i = \Gamma^2\rho h v^i, \quad E = \Gamma^2\rho h - p, \quad (vi)$$

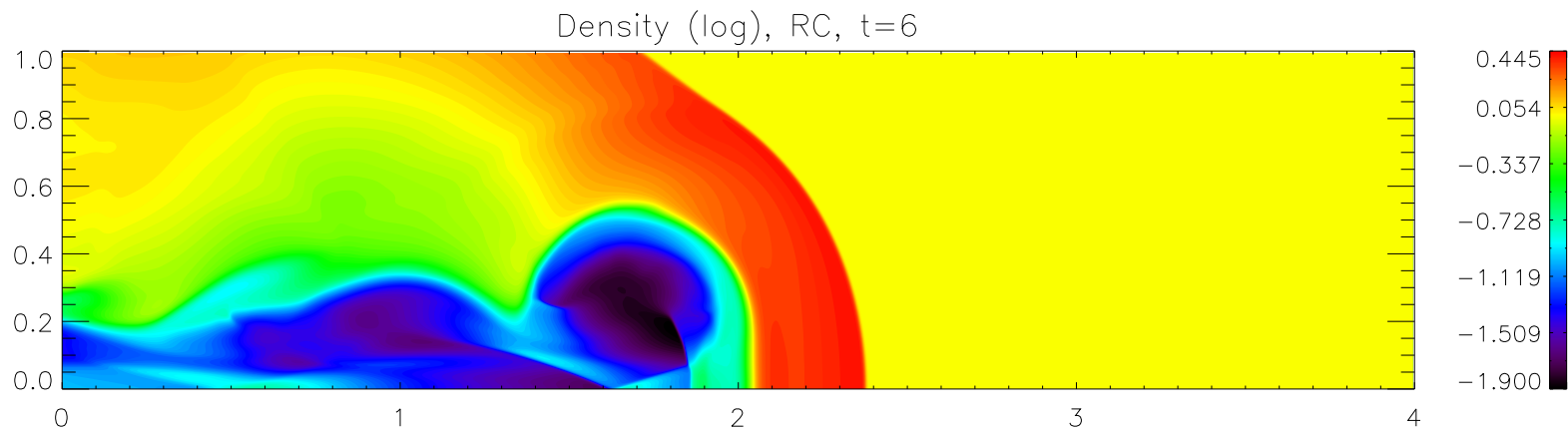
$$\Gamma = (1 + u_i u^i)^{1/2} = 1/(1 - v^2)^{1/2}$$

Relativistic JETS



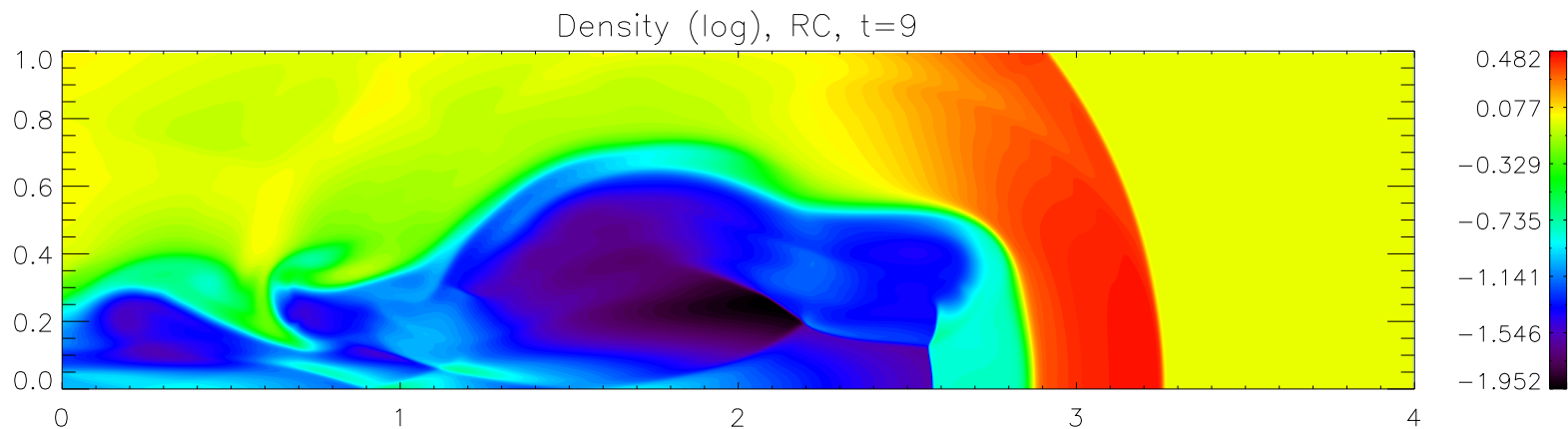
- Jet **RC**— $\log(\rho)$: t=3; Res:1024×256; $Y = 4X$; $X = 21.3 \times \text{jet radii}(r_d)$; Init cond:
 $v_x = 0.9$, $\rho_j = 0.1$, $\rho_a = 1$, $p = 0.01$; Initial jet length = r_d .

Relativistic JETS



- Jet **RC**— $\log(\rho)$: t=6;

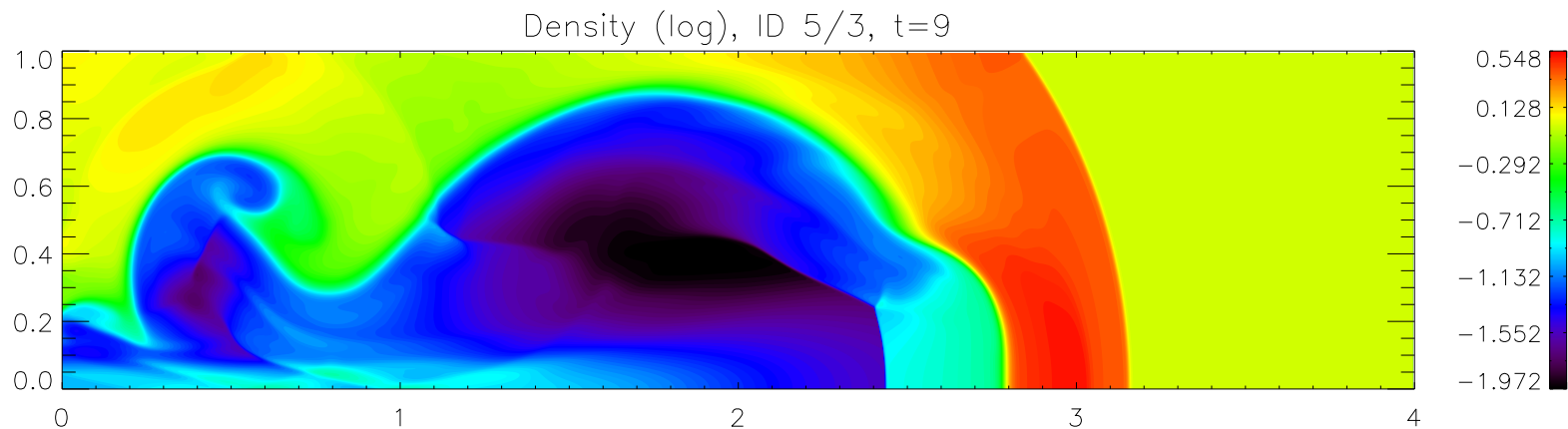
Relativistic JETS



- Jet **RC**— $\log(\rho)$: t=9;

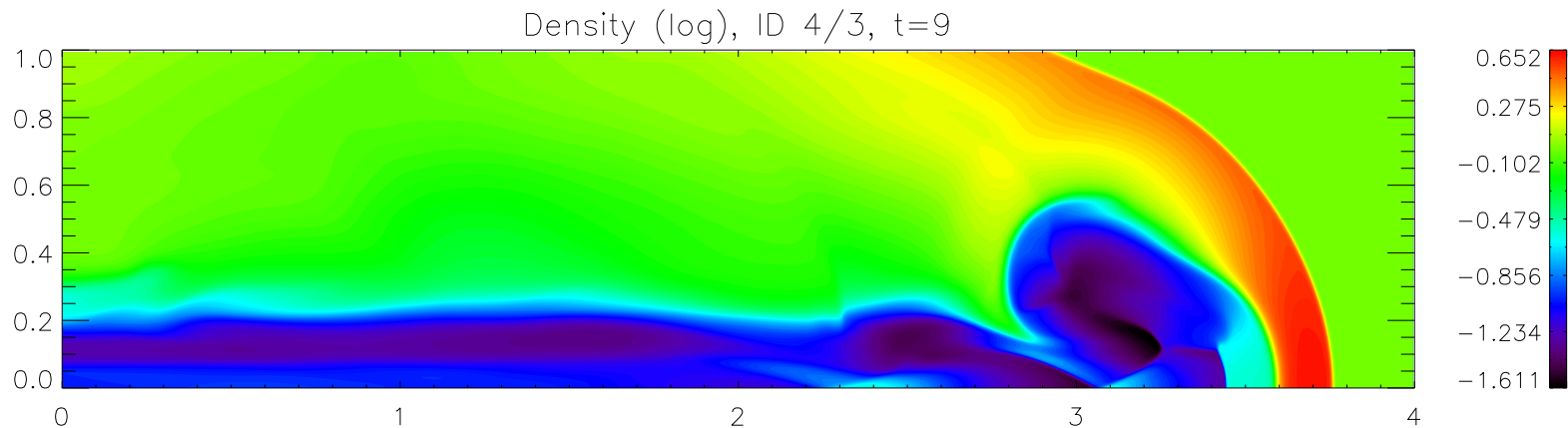
How does solutions with ID EoS look like at this time?

Relativistic JETS



- Jet ID $\gamma = 5/3$ — $\log(\rho)$: t=9; Res:1024×256; $Y = 4X$; $X = 21.3 \times \text{jet radii}(r_d)$;
Init cond: $v_x = 0.9$, $\rho_j = 0.1$, $\rho_a = 1$, $p = 0.01$; Initial jet length = r_d .

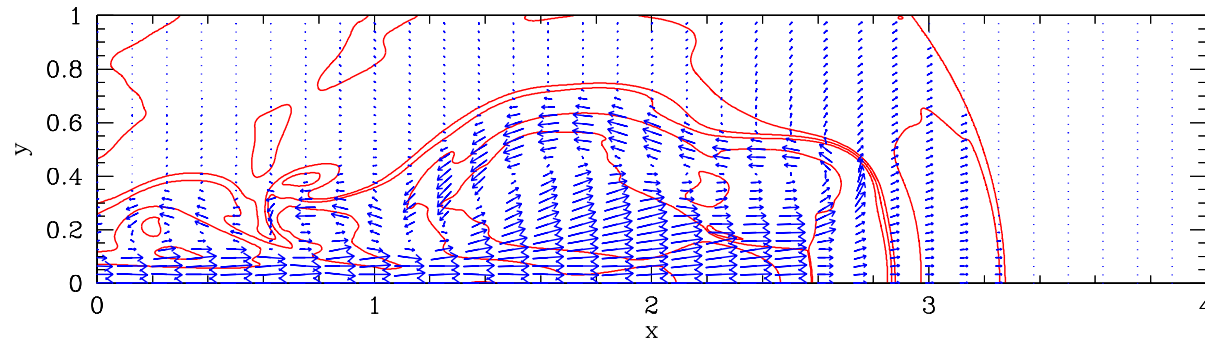
Relativistic JETS



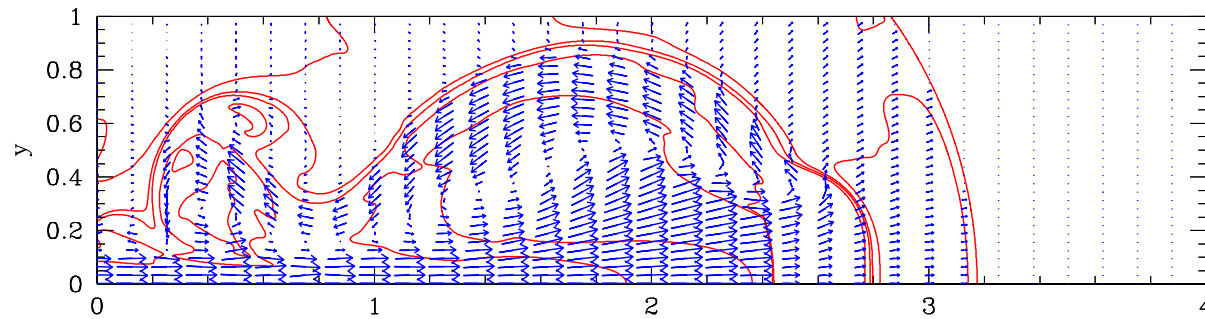
- Jet ID $\gamma = 4/3$ — $\log(\rho)$: t=9; Res:1024×256; $Y = 4X$; $X = 21.3 \times \text{jet radii}(r_d)$; Init cond: $v_x = 0.9$, $\rho_j = 0.1$, $\rho_a = 1$, $p = 0.01$; Initial jet length = r_d .
ID $\gamma = 4/3$ definitely is quite different, but why?

Relativistic JETS

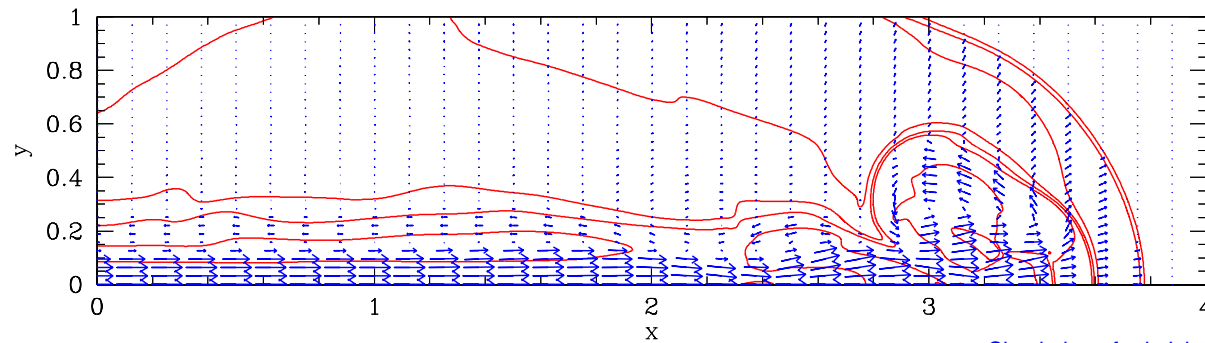
RC; $v_{xi}=0.9$; time=9



ID $\gamma=5/3$; $v_{xi}=0.9$; time=9



ID $\gamma=4/3$; $v_{xi}=0.9$; time=9

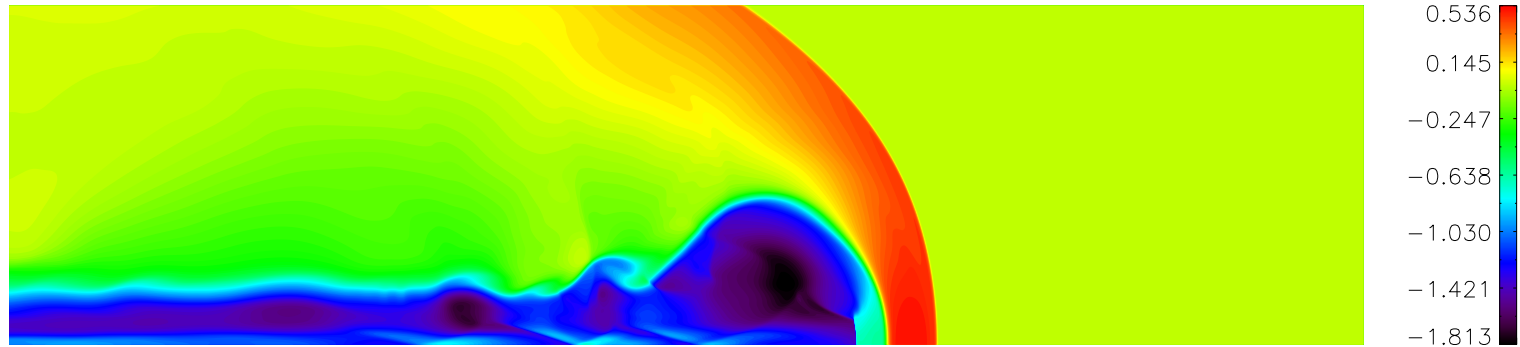


Relativistic JETS

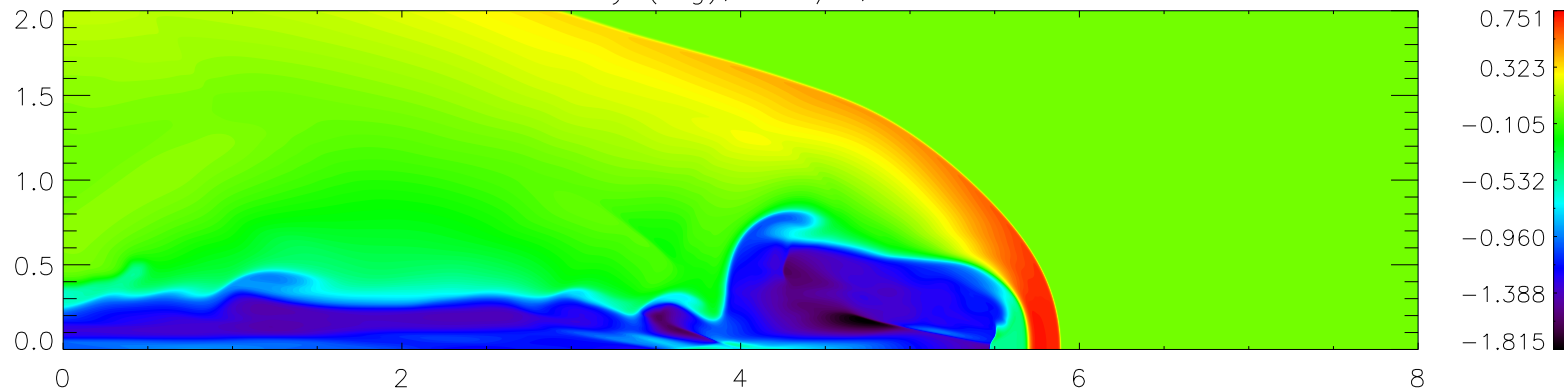
- As jet hits ambient medium it heats up
- A part of it flows back.
- With **ID** $\gamma = 5/3$ c_s is higher, it puffs up more, interacts with back flowing material, instabilities sets in
- With $\gamma = 4/3$, remains thin, limited back reaction, beam structure almost intact!
- With **RC** correct c_s gives a correct mode of interaction
- The bow shock is in median position between ID $\gamma = 5/3$ and $\gamma = 4/3$
- With $\gamma = 4/3$, even moderate injection Γ , produces near perfect beam structure
- Let us **jack up injection velocity**

Relativistic JETS

Density (log) RC, $t=12$



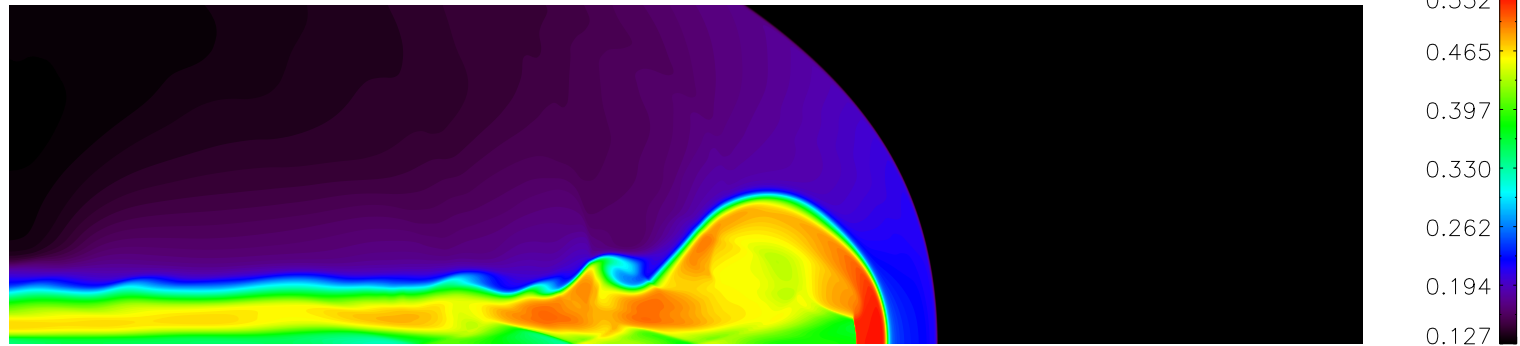
Density (log), ID $4/3$, $t=12$



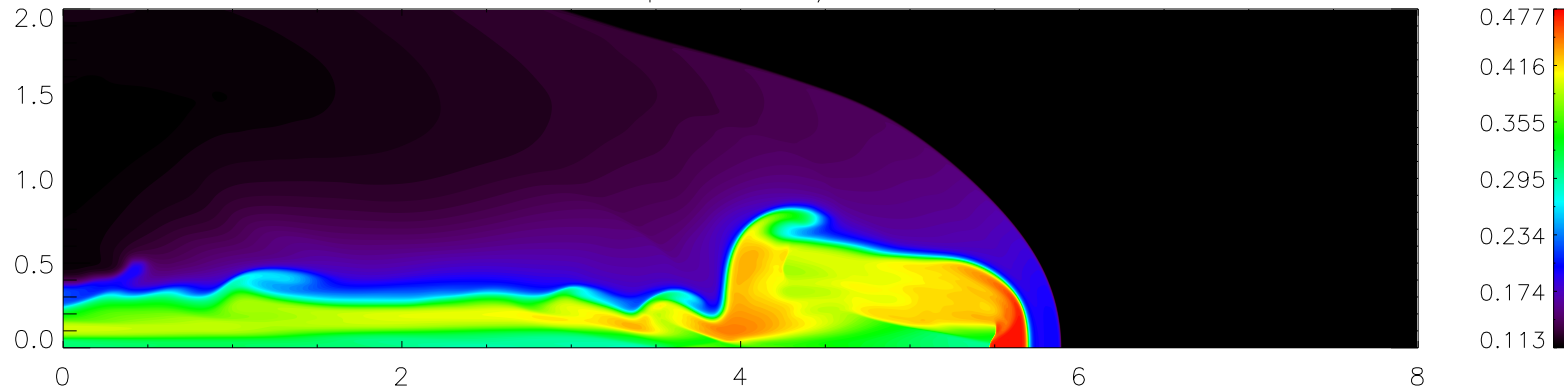
- Density contours at $t = 12$. Upper panel **RC**, lower panel **ID** ($\gamma = 4/3$), $v_{xi} = 0.95$.

Relativistic JETS

Sound speed, RC, $t=12$



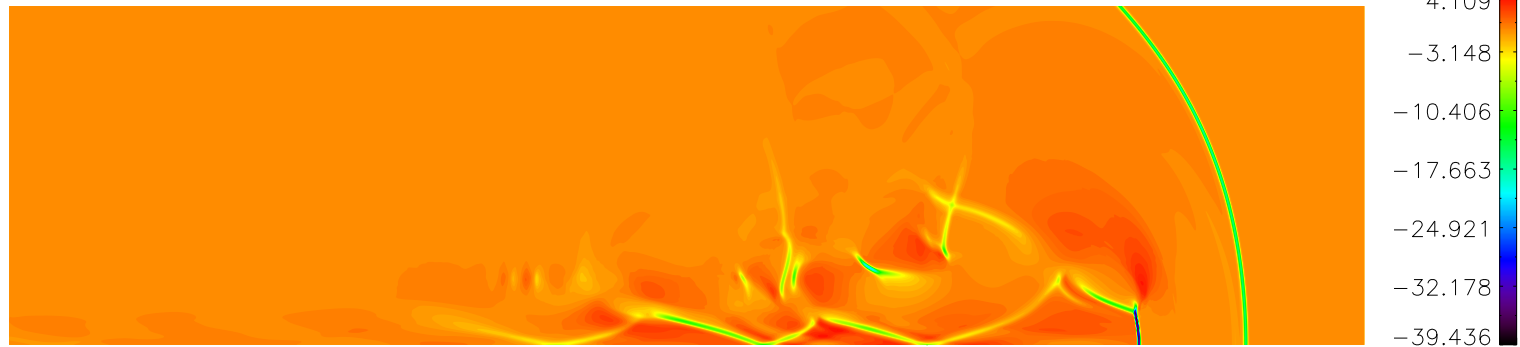
Sound speed, ID $4/3$, $t=12$



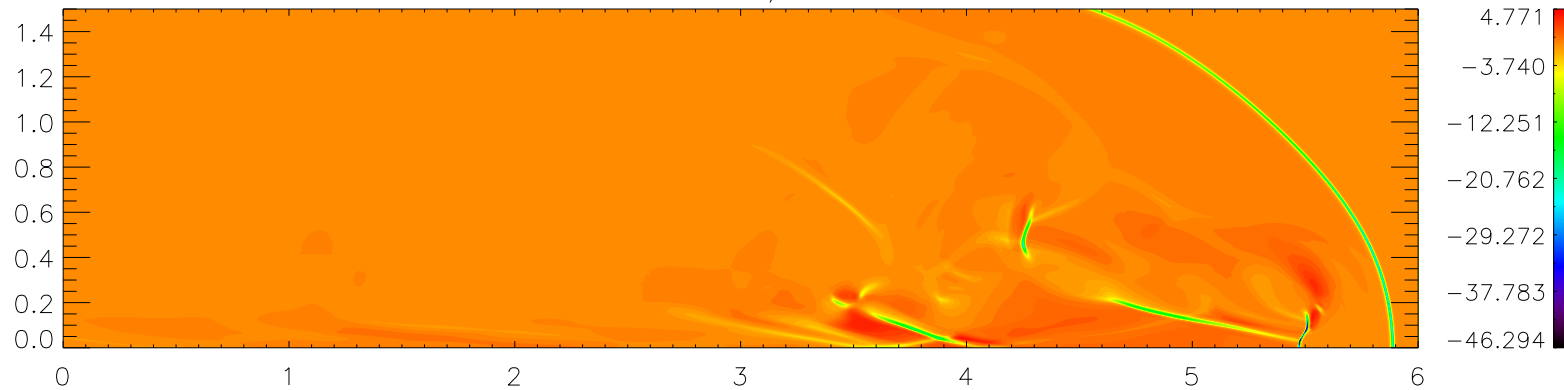
- Sound speed at $t = 12$. Upper panel **RC**, lower panel **ID** ($\gamma = 4/3$), $v_{xi} = 0.95$.

Relativistic JETS

Shocks, RC, t=12

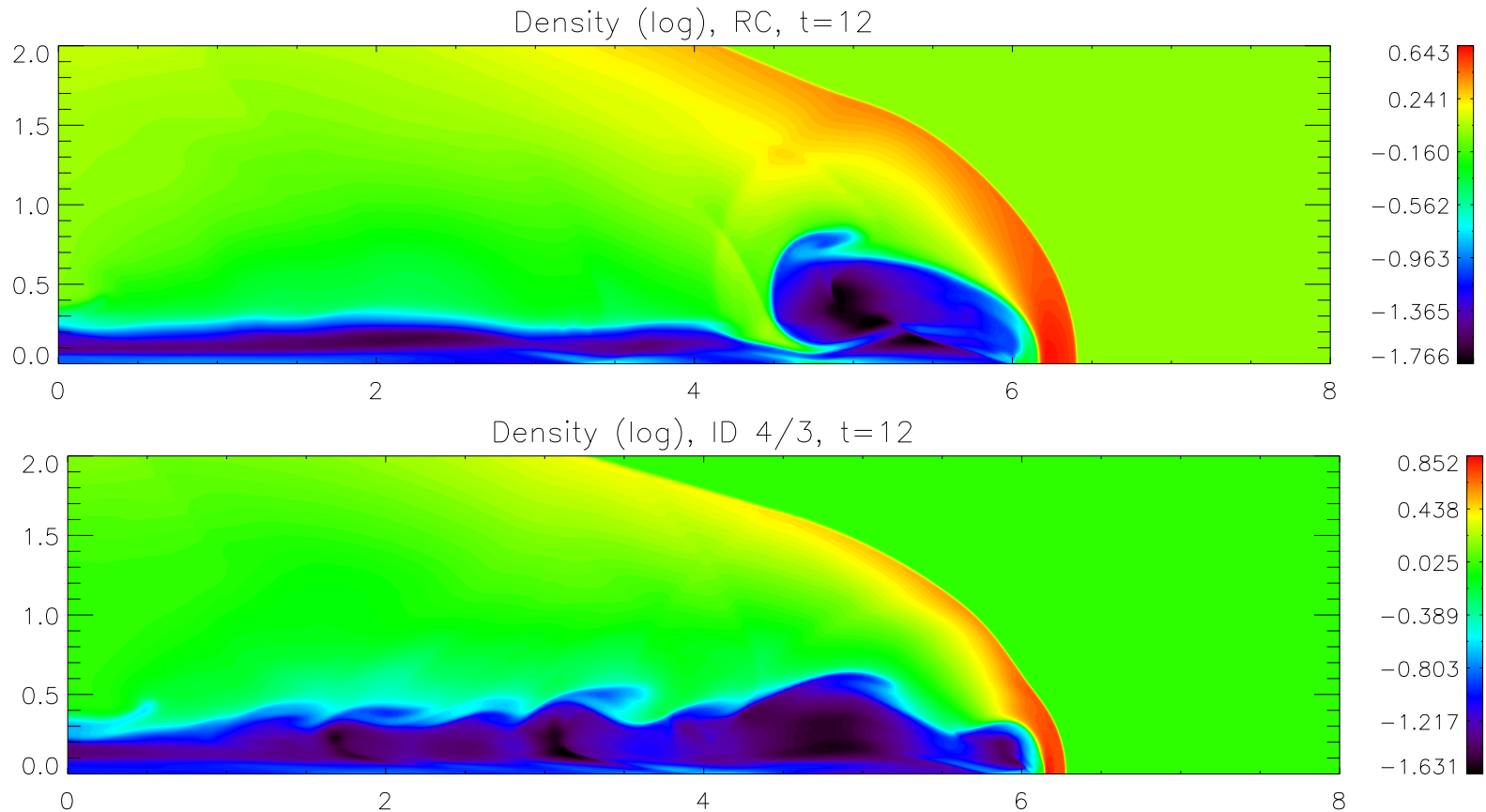


Shocks, ID 4/3, t=12

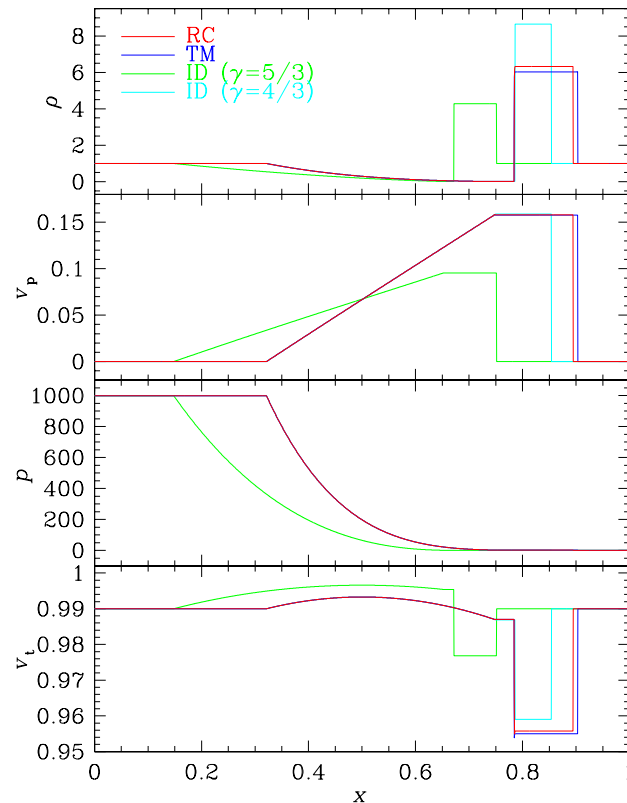


- Shocks at $t = 12$. Upper panel **RC**, lower panel **ID** ($\gamma = 4/3$), $v_{xi} = 0.95$. Should affect emitted radiation.

Relativistic JETS



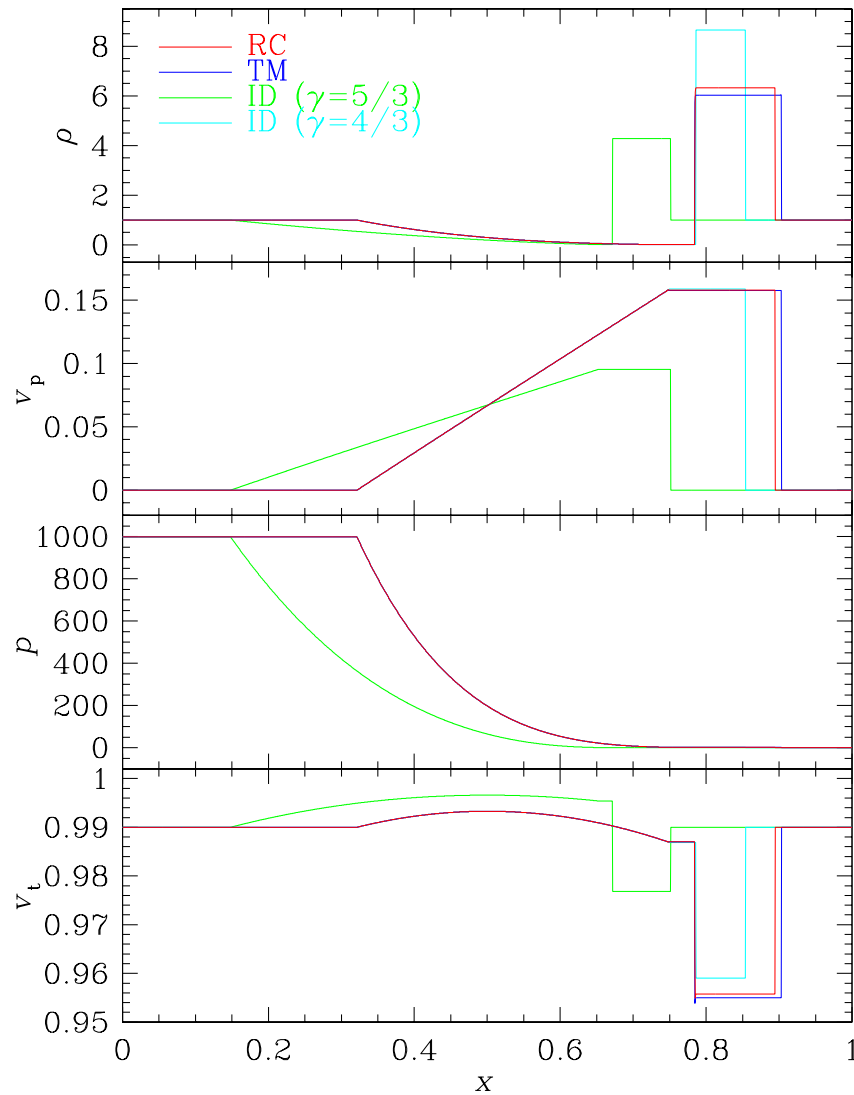
- $t = 12$; Density: $v_{xi} = 0.96$;
solution with RC is faster!!!



- Strong-Shock Tube — $t=0.18$; Init cond:

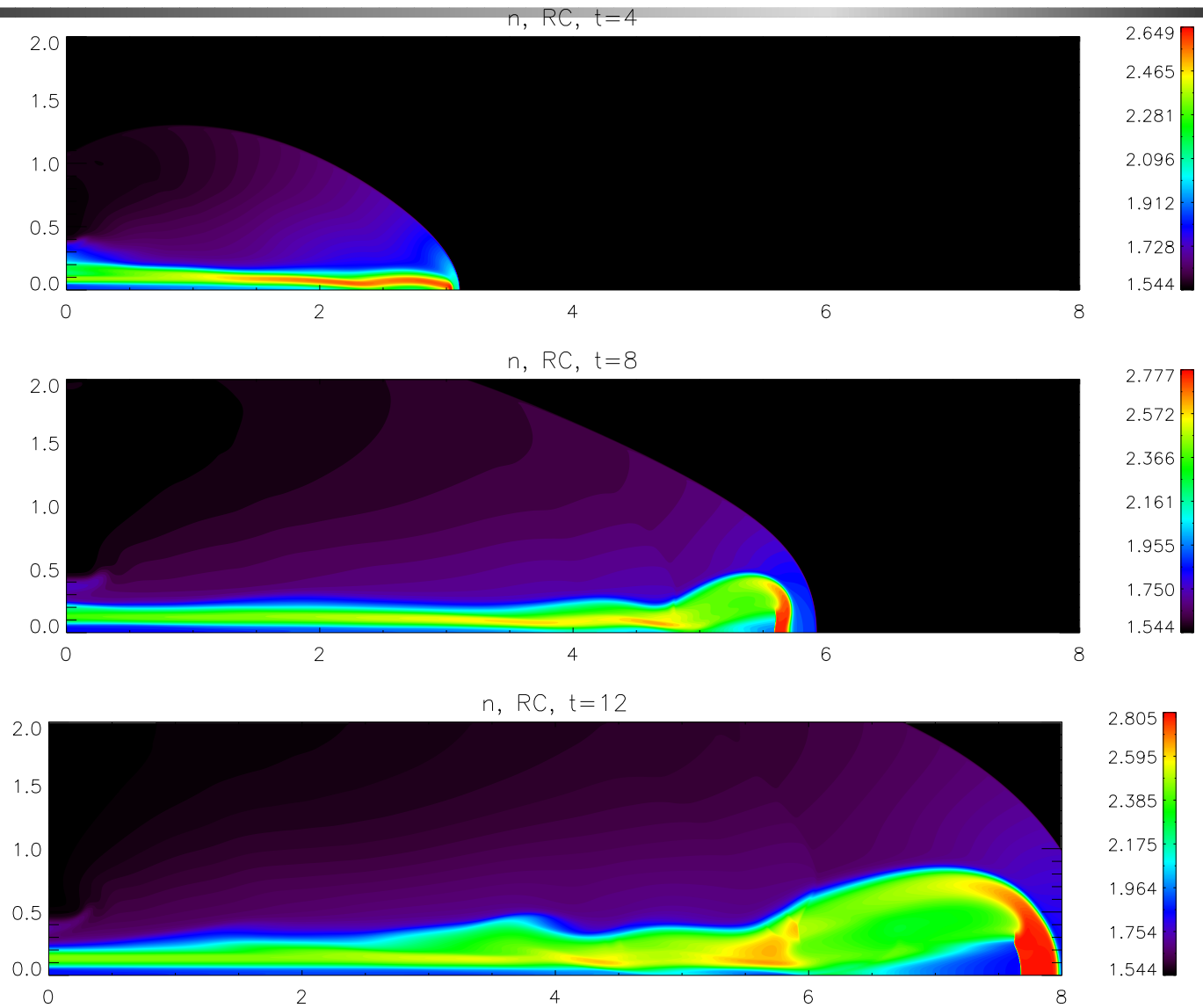
$$\begin{aligned}
 (\rho, v^x, v^y, v^z, p) &= (1, 0, 0.99, 0, 10^3); 0 \leq x \leq 0.5 \\
 &= (1, 0, 0.99, 0, 10^{-2}); 0.5 < x \leq 1.
 \end{aligned}$$

Relativistic JETS



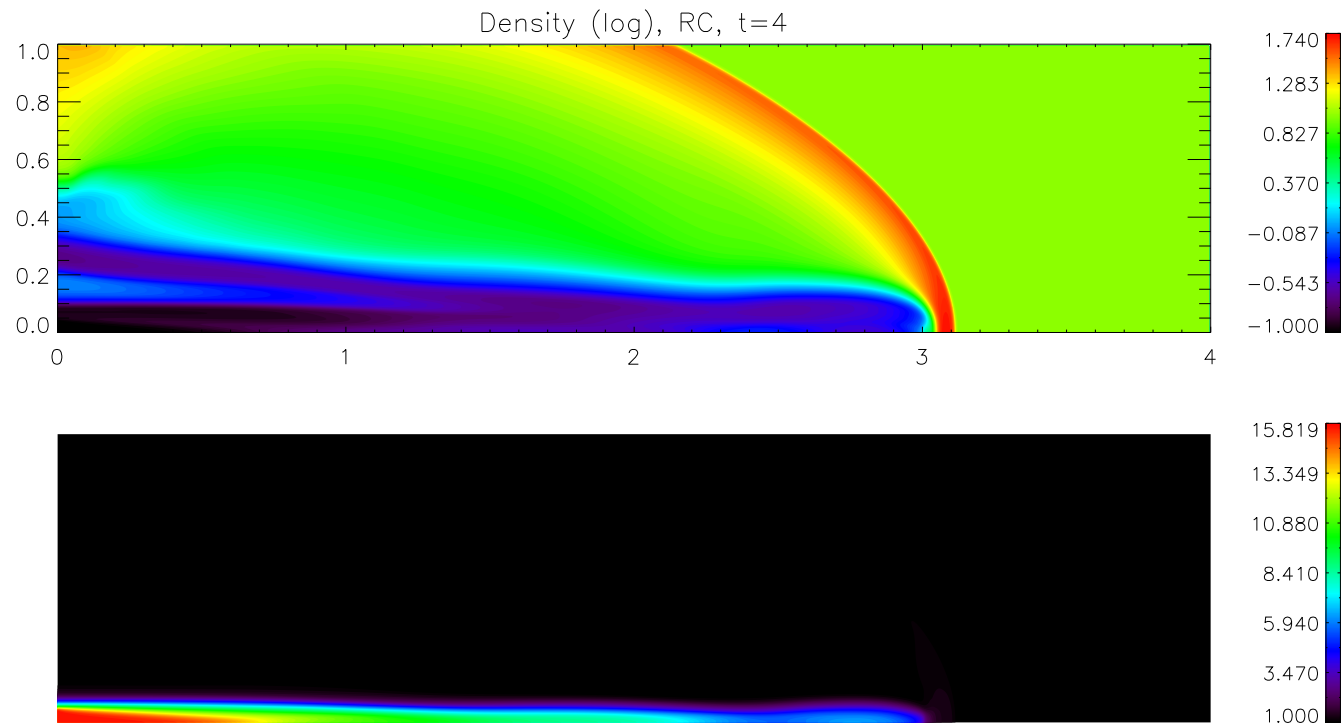
- Using ID EoS and some intermediate values of γ is not good, since γ generated by the interaction of the outflow with the the ambient medium

Relativistic JETS



- RC; $v_{ri} = 0.99$; n is variable and is neither 3 nor 3/2

Relativistic JETS



- RC $v_{xi} = 0.998$, upper panel density and lower panel Γ

Concluding Remarks

- ID EoS is relativistically wrong.
- Considering temperature dependent adiabatic/polytropic index is important.
- Internal structure produced by RC is different.
- These internal shocks will produce a different spectrum.
- Injection variables derived by analyzing observational data will also be faulty if wrong EoS is chosen.
- The jet head moves slower or faster than that modeled by ID EoS.
- Adiabatic index depends on the thermodynamic properties of the jet, which in turn is determined by the complicated interaction with the ambient medium, **thus cannot be assigned a priori**.
- Our preliminary investigation suggests, it will be interesting to investigate propagation of jets through non uniform medium.