The Role of the Magnetic Field in Black Hole-Neutron Star Electromagnetic Counterparts By Mew-Bing Wan Asia-Pacific Center for Theoretical Physics, South Korea

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Outline

- Background and motivations
- System
- Methods
- Results
- Ongoing work

Coalescing compact binaries

are the strongest gravitational wave (GW) emitters

detectable by ground-based detectors,

e.g., LIGO, VIRGO, GEO, KAGRA, IndIGO

 \Rightarrow strong tool for verifying the existence of GWs as a prediction of general relativity (GR)

• Accompanying electromagnetic wave (EM) signature

 \Rightarrow effectively localizes the source of GWs + its environment

 \Rightarrow improves confidence of GW detection

• Kilonova emission in conjuction with sGRB 130603B

 \Rightarrow compact binary progenitor hypothesis for short gamma-ray bursts (sGRBs) reinforced

⇒ sGRBs are strong EM counterparts for compact binaries



(Nissanke, *et al*, 2012)

TABLE 2

Representative GW Network Scenarios for Detectable Samples of NS-NS and NS-BH mergers. The notation med. and max. Refer to the median and maximum values of parameter distributions.

Feature	Lower Bound Scenario	Upper Bound Scenario Coherent 5 detector: Net5b 138 yr ⁻¹ (NS-NS) 17 yr ⁻¹ (NS-BH)	
Relative fractions & rates	Coincident 3 detector: Net3a 19 yr ⁻¹ (NS-NS) $ $ 3 yr ⁻¹ (NS-BH)		
Detectable distance	Coincident 3 detector: Net3a 220–400 Mpc (med-max; NS-NS) 350–600 Mpc (med-max; NS-BH)	Coherent 5 detector: Net5b 390–750 Mpc (med-max; NS-NS) 650–1250 Mpc (med-max; NS-BH)	
Sky Area Errors	Coherent 3 detector: Net3b $55-180 \text{ deg}^2$ (med-max; NS-NS) $50-170 \text{ deg}^2$ (med-max; NS-BH)	Coincident 5 detector: Net5a 7–120 deg ² (med-max; NS-NS) 6-65 deg ² (med-max; NS-BH)	

TABLE 1

Relative fractions $\times 10^{-4}$ of NS–NS and NS–5 M_{\odot} BH Mergers with collimated (denoted 'B') and isotropic (denoted 'I') emission Detectable in GWs Using Three Different Selection Criteria with Four GW Networks. The notation 'OS' represents optical squeezing in the LIGO interferometers. The range given represents the 1- σ statistical error of our simulation.

GW Network		Net3	Net4I	Net4K	Net5
		B I	B I	B I	B I
Coincident "a"	NS-NS	$0.3 \pm 0.3 11 \pm 2$	$0.3 \pm 0.3 17 \pm 2$	$0.3 \pm 0.3 17 \pm 2$	$0.3 \pm 0.3 23 \pm 2$
	$\rm NS-5M_{\odot}BH$	$0.7\pm0.5 50~{\pm}4$	$1.0 \pm 0.6 79 \pm 5$	$1.3 \pm 0.7 77 \pm 4$	$2.3 \pm 0.9 104 \pm 6$
Coherent "b"	NS-NS	$0.8 \pm 0.4 36 \pm 3$	$0.8 \pm 0.4 57 \pm 4$	$0.8 \pm 0.4 \mid \! 59 \pm 4$	$1.5 \pm 0.6 78 \pm 4$
	$\rm NS-5M_{\odot}BH$	$2.3\pm0.9 170\pm7$	$3.7 \pm 1.1 251 \pm 9$	$4.0\pm1.2 243\pm9$	$4.7 \pm 1.2 323 \pm 10$
EM precursor	NS-NS	$0.8 \pm 0.4 54 \pm 4$	$1.0 \pm 0.5 80 \pm 4$	$1.8 \pm 0.6 81 \pm 4$	$2.0 \pm 0.7 113 \pm 5$
	$\rm NS\text{-}5M_{\odot}BH$	$3.0\pm1.0 244\pm9$	$4.0\pm1.2 350\pm\!11$	$4.3\pm1.2 350\pm11$	$6.3 \pm 1.5 464 \pm 12$

• Can a black hole-neutron star (BHNS) merger be an sGRB progenitor?/ What are the possible EM counterparts that can be generated by a BHNS merger?

• What is the role of the NS magnetic field in generating these EM counterparts?

• How does the NS magnetic field amplitude and configuration evolve under various amplification mechanisms?

Numerical relativity

• Etienne+ '12:

BHNS mergers can power *jets* along funnel above and below equatorial plane with a strong enough poloidal field threading the accretion disk

- Giacomazzo+ '14:
- small-scale dynamo acting

inside hypermassive neutron star (HMNS) formed from BNS mergers;

dissipation of amplified B field can generate sGRB precursor flares

• Paschalidis+ '14:

BHNS mergers can power *jets* along funnel above and below equatorial plane via the action of NS magnetosphere;

role of the MRI in the accretion disk is not clear

due to insufficient resolution

- Topology of B field configuration in NS
- ⇒ generation of a large-scale poloidal field in merger remnant + effective *jet collimation*
- Magneto-rotational instability (MRI) in merger remnant
- \Rightarrow amplification of toroidal + poloidal field in merger remnant
- Our way:
- Can the MRI work to generate a large-scale poloidal field in the accretion disk + surrounding corona after a BHNS merger?

System



BH mass = 4 X NS massNS with APR EOS





NS

~10¹⁵ G

BH





• angular momentum transport • magnetic field winding: poloidal field in NS \rightarrow toroidal field in accretion disk





 central potential → outwardly decreasing angular momentum
 magnetic tension = tension of elastic coil connecting 2 orbiting particles
 → resisting stretching → stabilizing
 → resisting shearing → de-stabilizing

System

linear solution of the wave equation

$$\ddot{\xi}_{R} - 2\Omega \dot{\xi}_{\phi} = -\left(\frac{d\Omega^{2}}{d \ln R} + (\mathbf{k} \cdot \mathbf{u}_{A})^{2}\right) \xi_{R}$$

$$\ddot{\xi}_{\phi} + 2\Omega \dot{\xi}_{R} = -(\mathbf{k} \cdot \mathbf{u}_{A})^{2} \xi_{\phi}$$

exact non-linear solution of the full axisymmetric hydrodynamic equations \equiv "channel" solution for accretion disk (Goodman & Xu, '94) 1.46 2.43 1.70 2.68 1.95 2.92 2.19

(Figure: Balbus & Hawley, 1998)

3.1



• Instability criteria a) $\frac{d\Omega^2}{d \ln R} < 0$ b) $(\mathbf{k} \cdot \mathbf{u}_{\mathbf{A}})^2 < -\frac{d\Omega^2}{d \ln R}$



• Maximum unstable growth rate a) axisymmetric analysis

$$\left|\omega_{max}\right| = \frac{1}{2} \left|\frac{d\Omega}{d \ln R}\right|, \left(\mathbf{k} \cdot \mathbf{u}_{\mathbf{A}}\right)_{max}^{2} = -\left(\frac{1}{4} + \frac{\kappa^{2}}{16\Omega^{2}}\right)\frac{d\Omega^{2}}{d \ln R}$$

b) non-axisymmetric analysis $k_Z \gg k_R \sim k_R(0) \gg m \frac{d \ln \Omega}{dR} \rightarrow \text{axisymmetric limit}$

 $\Rightarrow \text{ poloidal field magnitude determines maximum growth rate}$ $\Rightarrow wavelength of maximum growth mode, <math>\lambda_{\text{MRI}} \sim \left(\frac{2\pi}{\Omega}\right) \left(\frac{B_i e_k^i}{\sqrt{4\pi \sigma}}\right)$



System



high-velocity *matter jet*/ collimated by B field

luminous short burst / of gamma rays

> initially matter-poor region along BH rotation axis

> > hot dense massive

accretion disk

(Image: daviddarling.info/encyclopaedia)





(Image: Kiuchi, 2013)





grids preserving ∇.B=o well by Balsara's method ('09):
WENO5 volume and surface reconstruction for magnetic field
offset in the electric field during restriction
all GRMHD variables evolved at grid vertices

• Time-scale of sGRB observed by BATSE



 \Rightarrow long-term simulation in the time scale of sGRB (~100 ms)



Evolution of BH mass

Evolution of BH spin



• Disk mass 0.35 M_s \Rightarrow massive disk

• Angular momentum of BH increased due to accreted mass

Results



Results







- Ram velocity of matter in funnel \rightarrow fallback
- Wind on disk edges starts to develop at t~25ms after merger
- \rightarrow net mass outflow

Results



- B field advected by in-falling matter
- Turbulent cascade
- in accretion disk
- near funnel above and below merger remnant





- Amplification via MRI of both toroidal and poloidal field
- by ~1 order of magnitude
- Saturation reached at t~42ms after merger
- Resolution of wavelength of fastest mode
- a) toroidal direction: ~10 points
- b) poloidal direction: ~1 point



E_{toroidal} > E_{poloidal}
 by 1 order of magnitude
 saturation:

 $E_{poroidal}/KE_{plasma} \sim 0.1-0.2\%$ $E_{toroidal}/KE_{plasma} \sim 0.4-0.6\%$





Comparison with non-axisymmetric linear analysis (Balbus & Hawley, '92): % difference averaged from 40 points on equatorial plane along both toroidal and radial directions $13 \le R \le 107$ km



 both poloidal and toroidal fields amplified more readily by MRI in linear regime further away from black hole/low density regions of disk
 ⇒ large-scale poloidal field may emerge starting from outer edges of disk

Ongoing work

- Thompson & Duncan '93:
- direction of dipole in NS
- not aligned with axis of its rotation axis
- \Rightarrow vary configuration of magnetic field inside NS
- Expectations:
- a) increase in resolution of wavelength of fastest MRI mode in poloidal field
- b) saturation \rightarrow equipartition
- c) jet collimation \rightarrow sGRB?

Ongoing work

tilted embedded NS dipole

