

Hydrodynamic Simulations of the Central Molecular Zone with a Realistic Galactic Potential

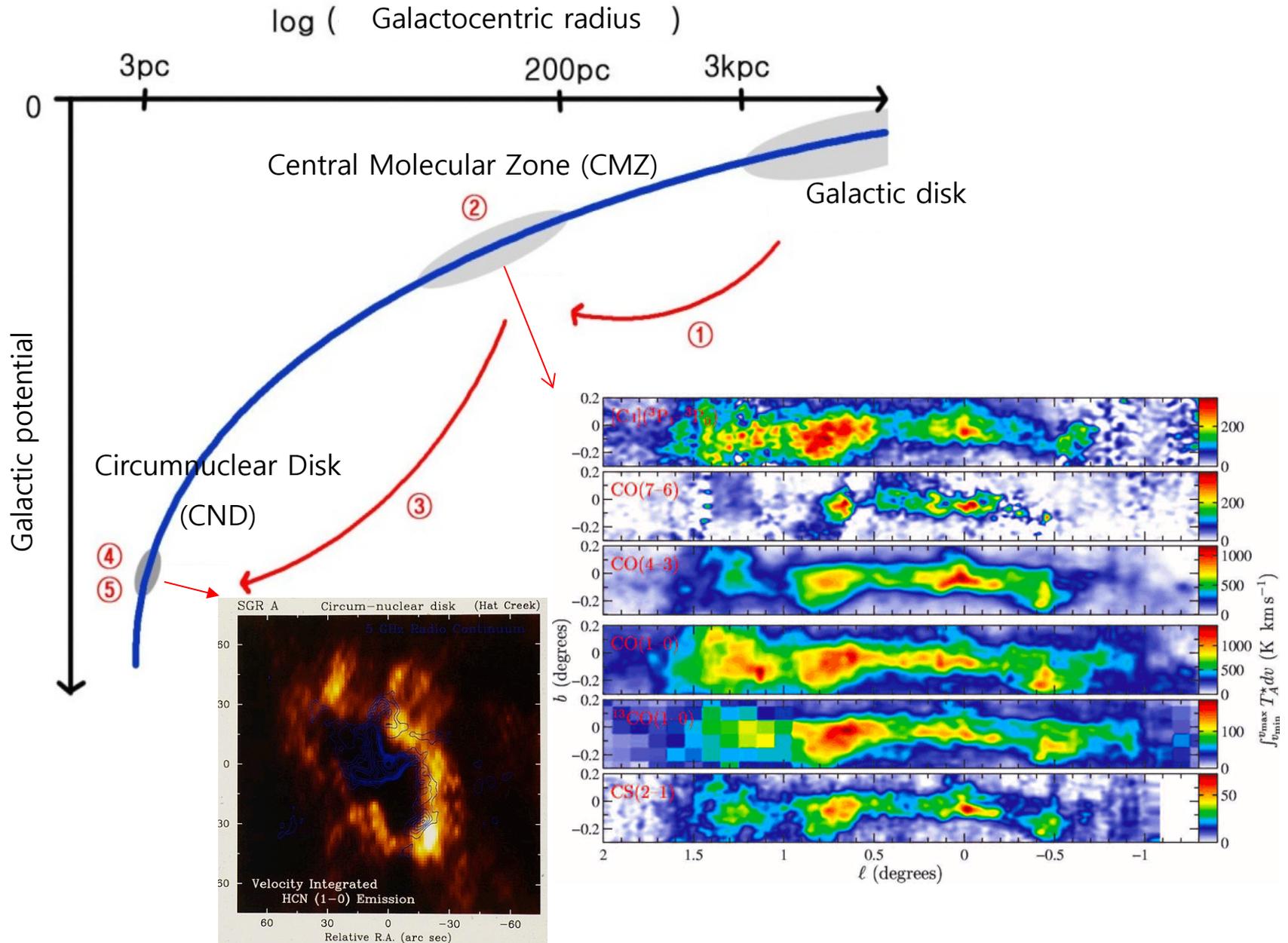
Jihye Shin (Korea Institute for Advanced Study)

Collaborators:

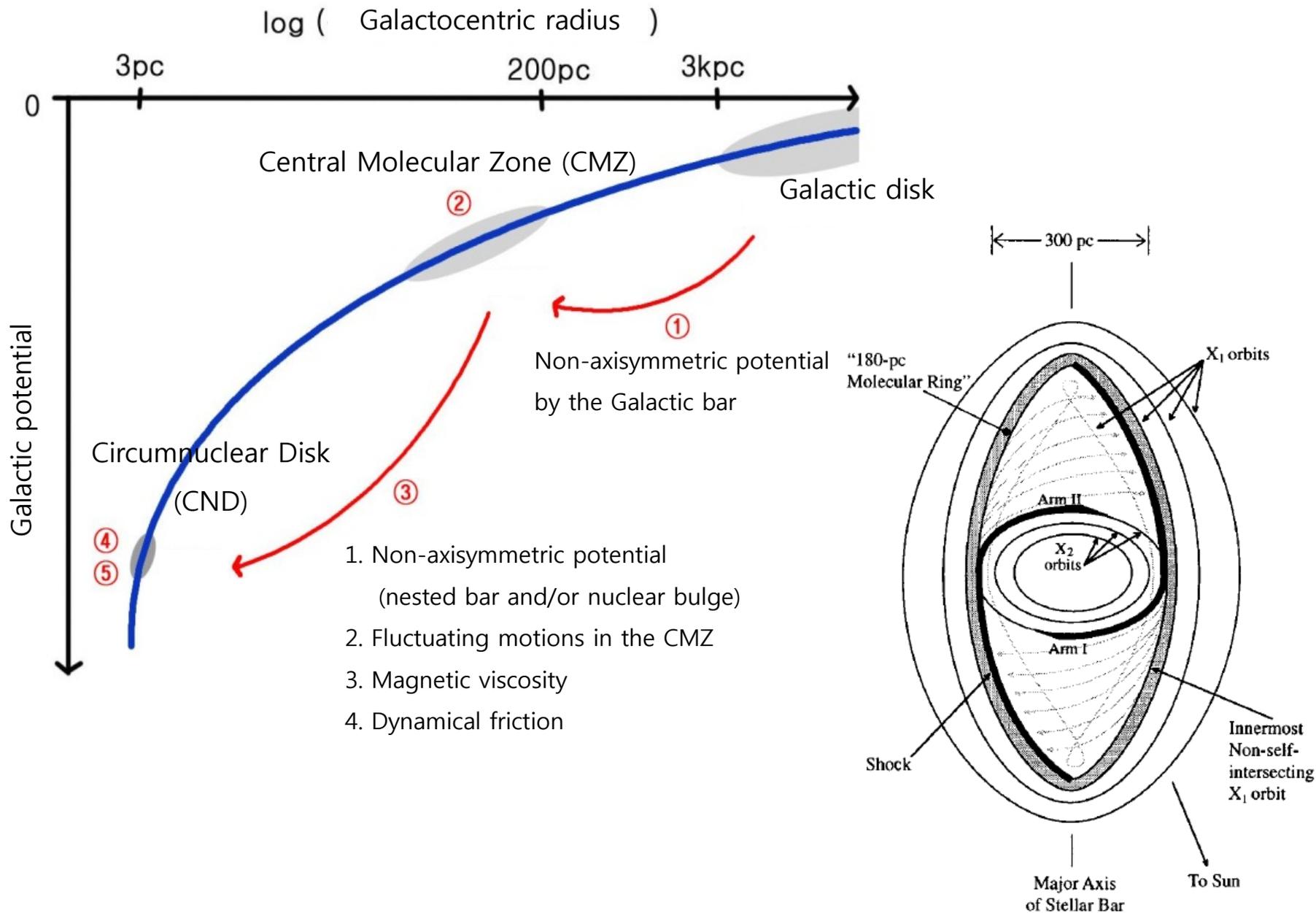
Sungsoo S. Kim¹, Junichi Baba², Takayuki Saitoh², Jeong-Sun Hwang¹, Kyungwon Chun¹, Shunsuke Hozumi³

(¹Kyung Hee University, Korea, ²Tokyo Institute for Technology, Japan, ³Shiga University, Japan)

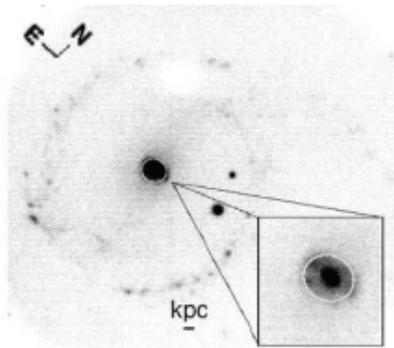
Migration of gas toward the Galactic center



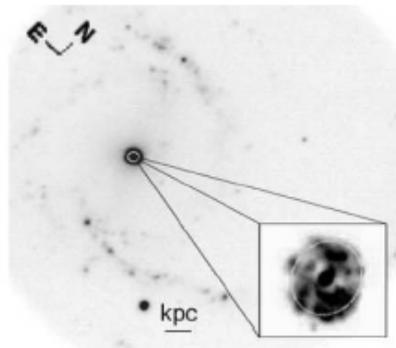
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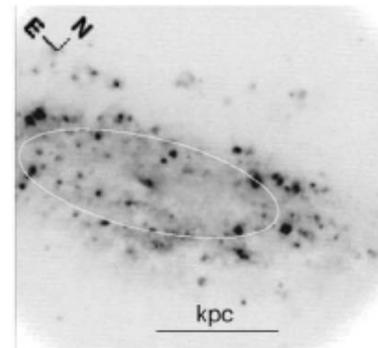
Nuclear ring in the external barred galaxies



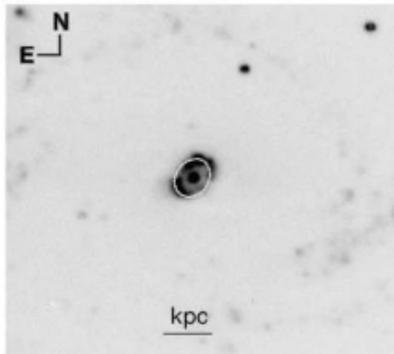
(m) NGC 5945



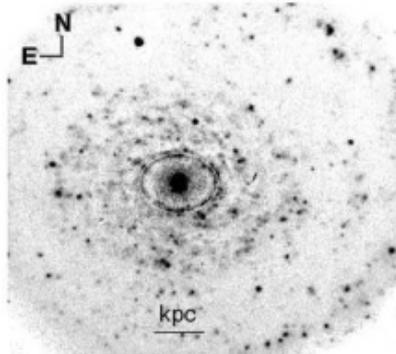
(n) NGC 5953



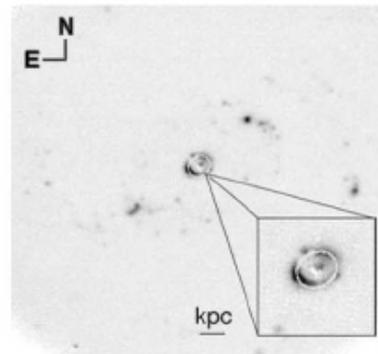
(o) NGC 6503



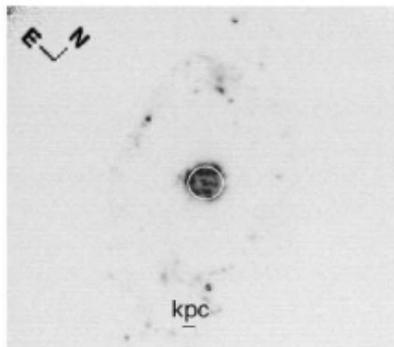
(p) NGC 6951



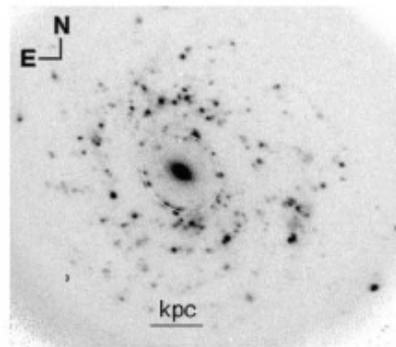
(q) NGC 7217



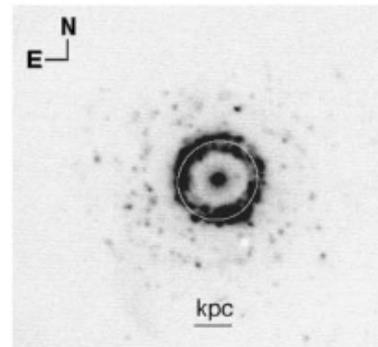
(r) IC 1438



(s) NGC 7570



(t) NGC 7716



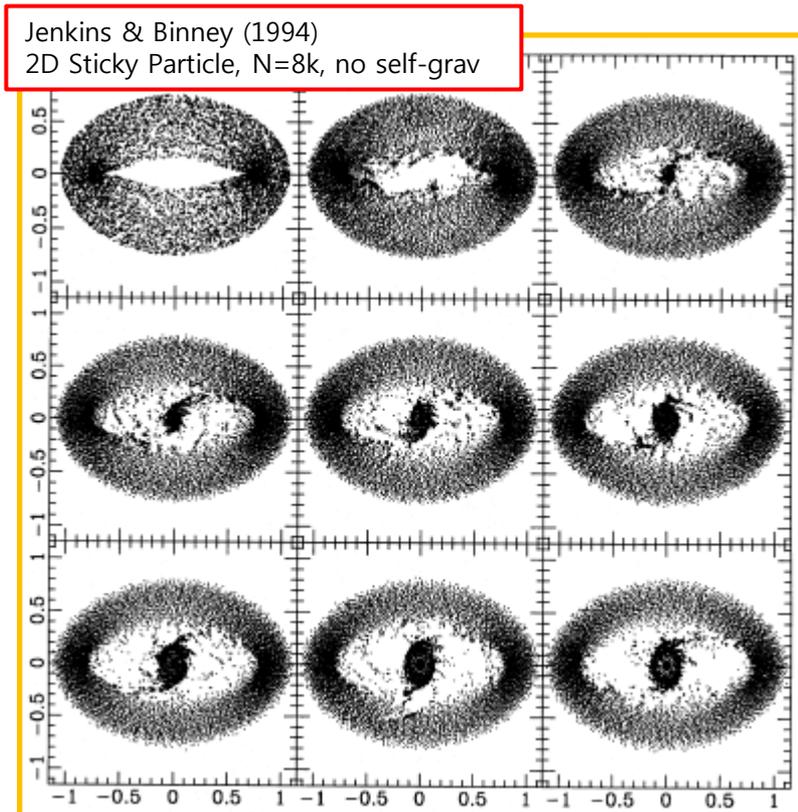
(u) NGC 7742

Numerical studies

1. Jenkins & Binney (1994) with 2D sticky particle simulation

Lee et al. (1999), Englmarier & Gerhard (1999) with 2D SPH simulation

-> transition of gas motion from X_1 to X_2 orbits takes places in a bar potential (Binney et al. 1991)



Numerical studies

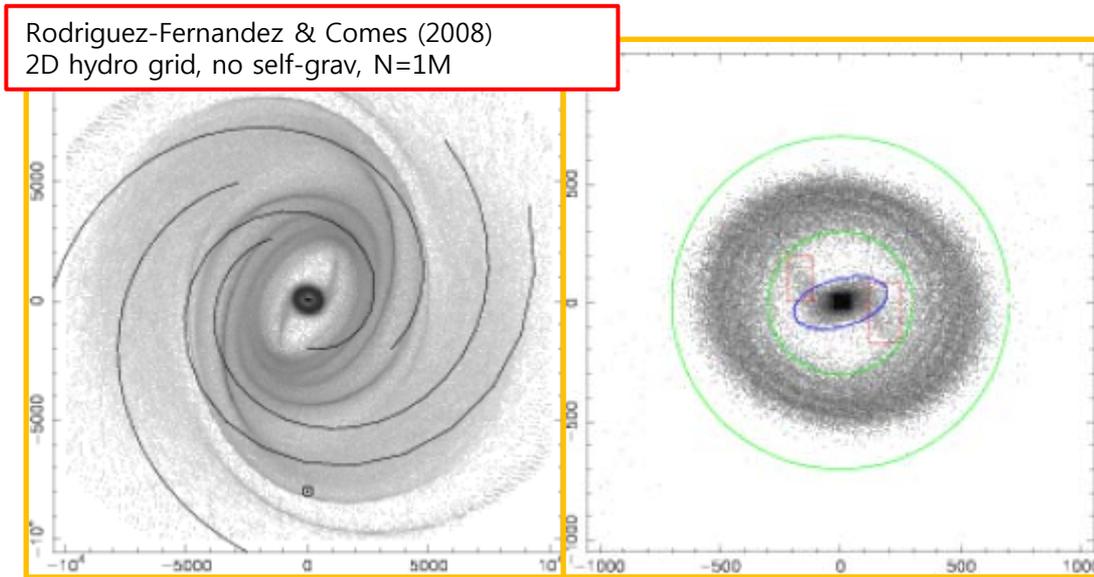
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-> including realistic mass distribution of the Galactic bulge using 2MASS star count map



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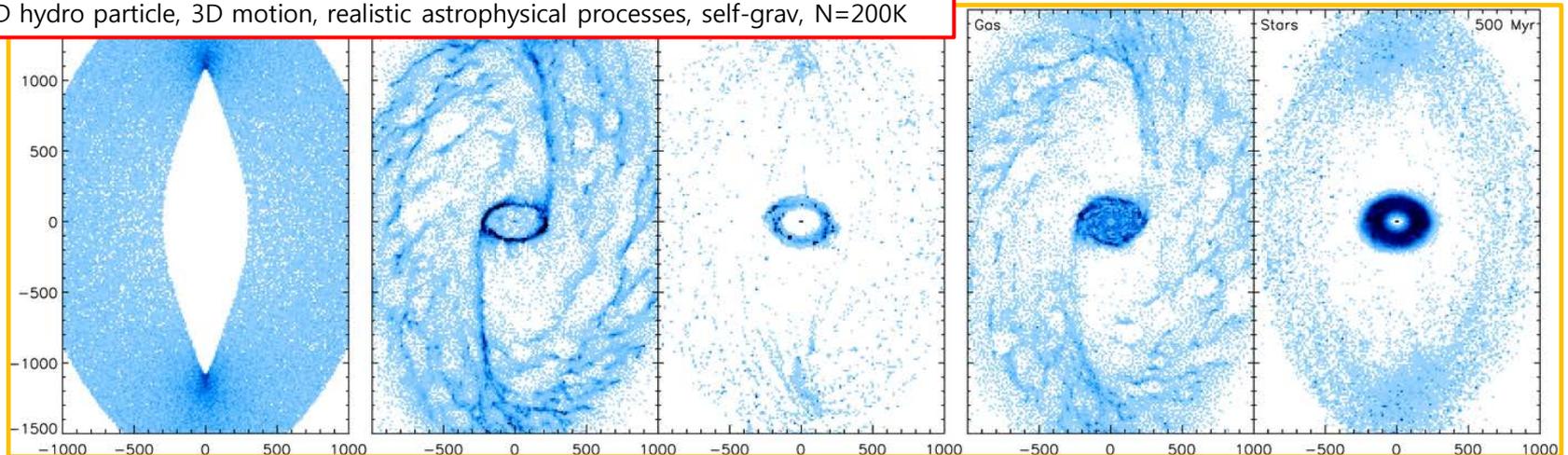
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3. Kim et al. (2011) with 3D SPH simulation

-> including realistic astrophysical processes for temperature evolution and star formation activities

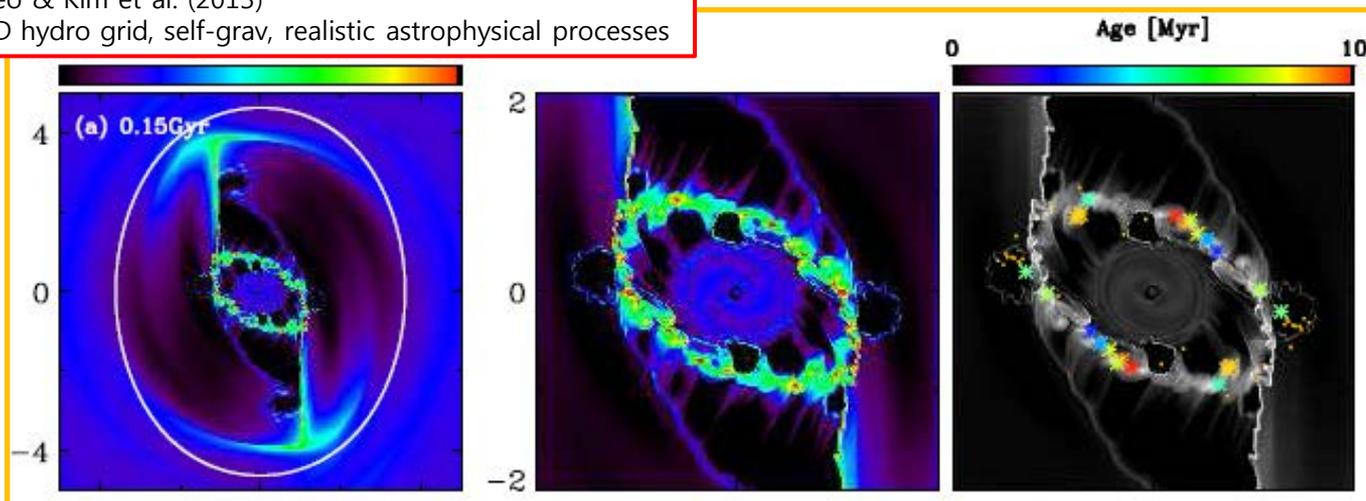
Kim et al. (2011)
3D hydro particle, 3D motion, realistic astrophysical processes, self-grav, N=200K



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4. Kim, Seo, & Kim (2012), Seo & Kim (2013, 2014) with grid-based 2D hydrodynamics
-> **physical properties and temporal evolution of nuclear ring** of barred and barred-spiral galaxies

Seo & Kim et al. (2013)
2D hydro grid, self-grav, realistic astrophysical processes



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Limitations of these works

- realistic galactic potential
- various astrophysical processes in the gas
- vertical motion to the disk

In this study, we aim to trace **the realistic motion of gas clouds** from the Galactic disk of $R_g > 3$ kpc to the ~ 200 pc CMZ under **the realistic Galactic mass distribution**.

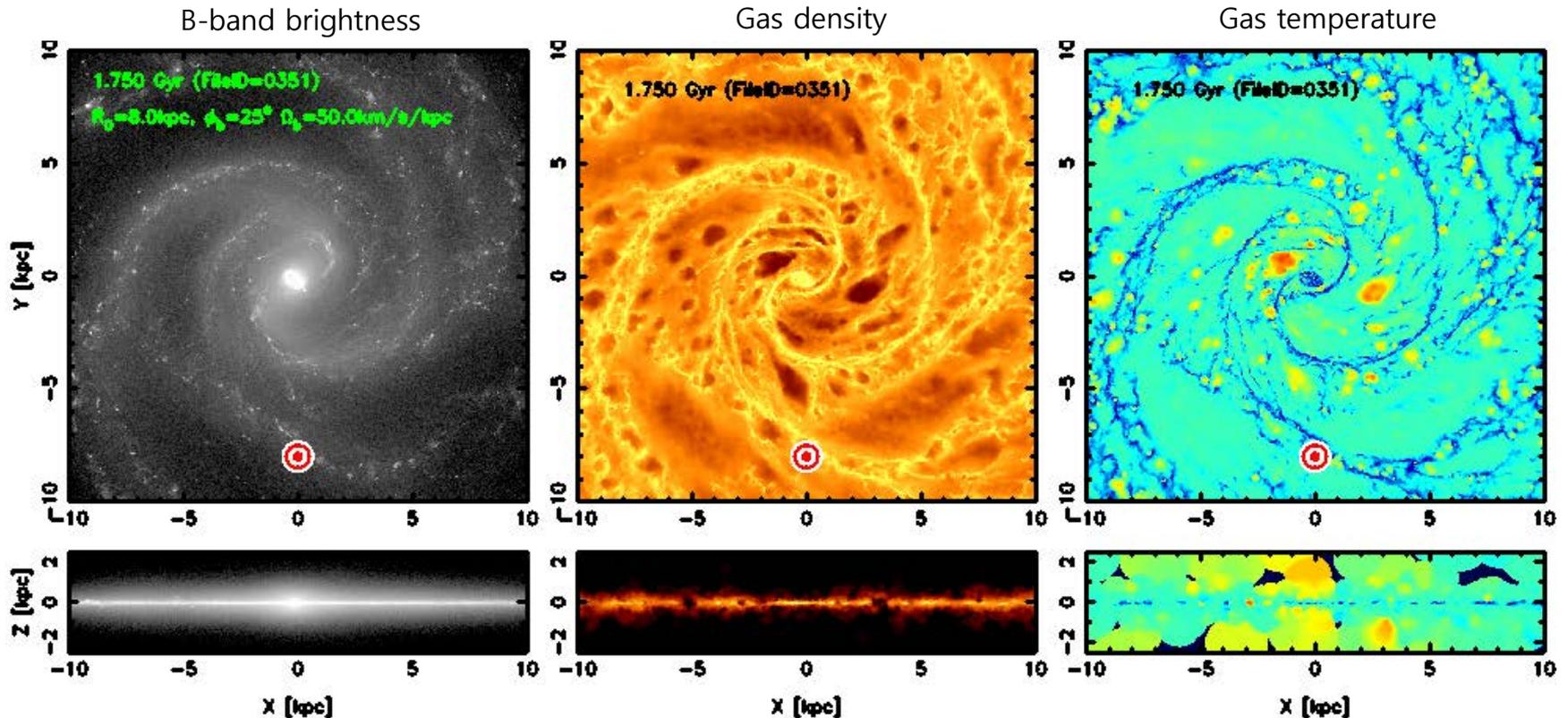
The Galaxy model

The standard compound Galaxy model = bulge + disk + halo

Non-axisymmetric structure = bar + spiral arms

We adopt a simulation snapshot of [Baba \(2015\)](#)

that reproduces the 3D structures of the stellar disk, grand-design spiral arms, bar, and also atomic/molecular gas layers of the Milky Way.



How to include the realistic 3D structures

1. Re-run the simulation?

-> more than one month using ~1000 cores.

2. Live particles?

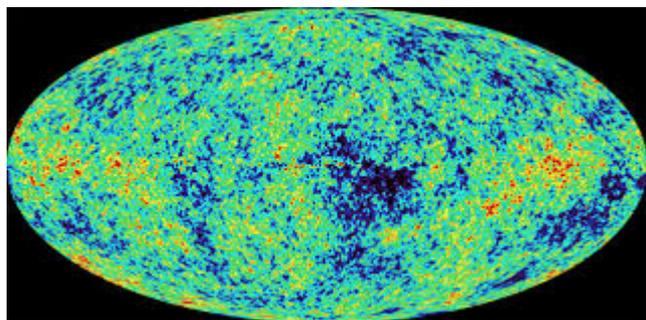
-> too heavy (~10 million particles for the Galactic structures)

3. Long-range force with latticed particle distribution?

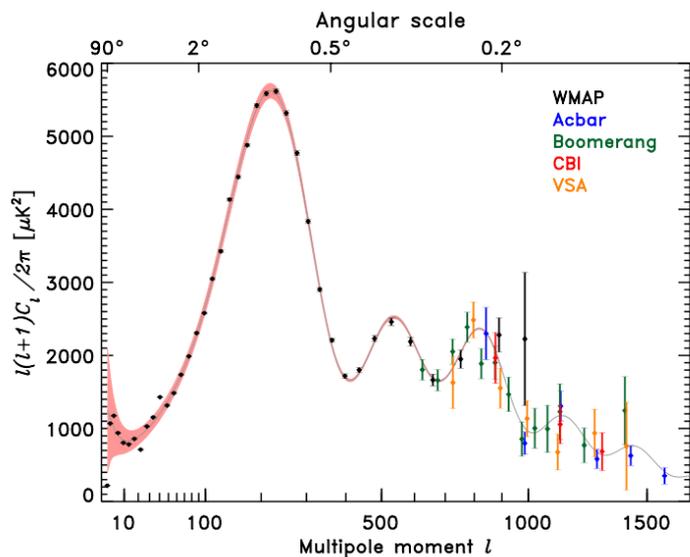
-> fast, but not accurate

How to include the realistic 3D structures

We describe the realistic 3D structures using multipole expansions (MEs).

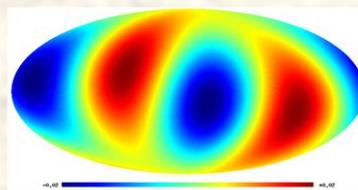


CMB (WMAP)

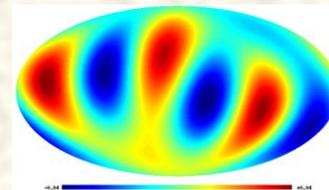


Multipole expansion

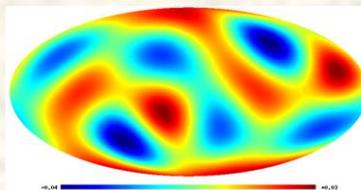
$$\Delta T(\theta, \phi) = \sum_{l=2}^{\infty} \sum_{m=-l}^{m=l} a_{l,m} Y_{l,m}(\theta, \phi)$$



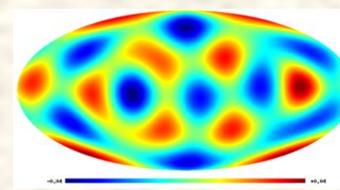
L = 2 (quadrupole)



L = 3 (octupole)



L = 4



L = 5

Multipole expansions to describe the Galactic structure

Bi-orthogonal basis set for density-potential pair

$$\rho(\vec{r}) = \sum_{nlm} A_{nlm} \rho_{nlm}(\vec{r}),$$
$$\text{PO } \Phi(\vec{r}) = \sum_{nlm} A_{nlm} \Phi_{nlm}(\vec{r}),$$

-> one-to-one relationship for ρ_{nlm} - Φ_{nlm} , satisfying

1. 3-dimensional, spherical model (Hernquist & Ostriker 1992) : **3D-spherical model**
2. 3-dimensional, cylindrical model (Earn 1996) : **3D-disk model**
3. 2-dimensional, polar model (Aoki & Iye 1978) : **2D-disk model**

Multipole expansions to describe the Galactic structure

1. 3-dimensional, spherical model (Hernquist & Ostriker 1992) : 3D-spherical model

$$\rho(r, \theta, \phi) = \sum_{nlm} A_{nlm} \rho_{nlm}(r, \theta, \phi),$$

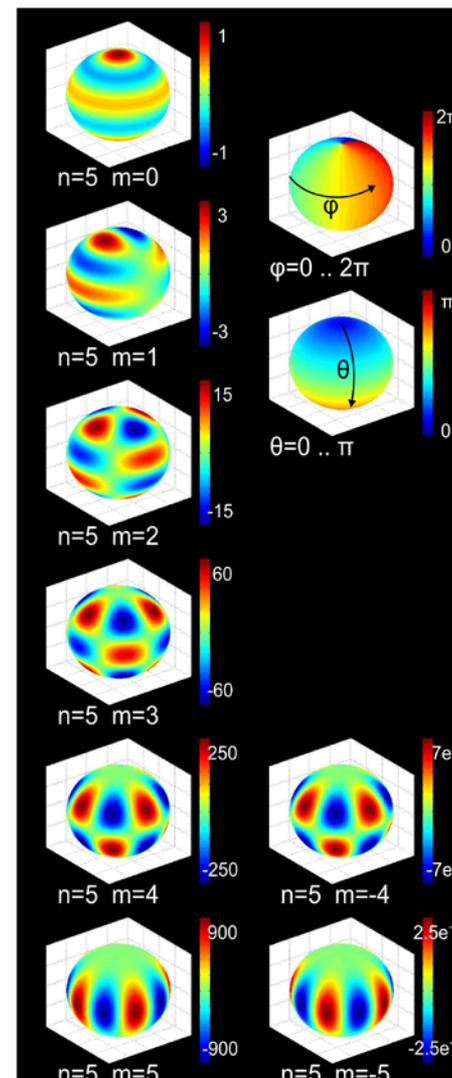
$$\Phi(r, \theta, \phi) = \sum_{nlm} A_{nlm} \Phi_{nlm}(r, \theta, \phi),$$

$$\rho_{nlm}(r, \theta, \phi) = \frac{K_{nl}}{2\pi} \frac{r^l}{r(1+r)^{2l+3}} C_n^{(2l+3/2)} \left(\frac{r-1}{r+1} \right) \sqrt{4\pi} Y_{lm}(\theta, \phi),$$

$$\Phi_{nlm}(r, \theta, \phi) = -\frac{r^l}{(1+r)^{2l+1}} C_n^{(2l+3/2)} \left(\frac{r-1}{r+1} \right) \sqrt{4\pi} Y_{lm}(\theta, \phi),$$

Gegenbauer Polynomial

Spherical harmonic



Multipole expansions to describe the Galactic structure

2. 3-dimensional, cylindrical model (Earn 1996) : **3D-disk model**

$$\rho(R, \phi, z) = \sum_{kmh} B_{kmh} \rho_{kmh}(r, \phi, z),$$

$$\Phi(R, \phi, z) = \sum_{kmh} B_{kmh} \Phi_{kmh}(r, \phi, z),$$

$$\rho_{kmh}(r, \phi, z) = \frac{1}{4\pi G} J_m(kR) e^{im\phi} k^2 e^{-k|z|},$$

$$\Phi_{kmh}(r, \phi, z) = -\frac{1}{2} J_m(kR) e^{im\phi} (1 + k|z|) e^{-k|z|},$$

Cylindrical Bessel function.

3. 2-dimensional, polar model (Aoki & Iye 1978) : **2D-disk model**

$$\mu_{nm}(R, \phi) = \frac{M_a}{R_a^3} \frac{(2n+1)}{2\pi} \exp(im\phi) P_{nm}(r)$$

$$\Psi_{nm}(R, \phi) = -\frac{GM}{R_a} \exp(im\phi) P_{nm}(r)$$

Legendre function

How to use the multipole expansions

A. 3D SPHERICAL MULTIPOLE EXPANSION MODEL OF HERNQUIST & OSTRIKER (1992)

For a computational convenience, density ρ and potential Φ in a position of (r, θ, ϕ) are rewritten as

$$\rho(r, \theta, \phi) = \sum_{l=0}^{\infty} \sum_{m=0}^l P_{lm}(\cos \theta) [A_{lm}(r) \cos m\phi + B_{lm}(r) \sin m\phi],$$

$$\Phi(r, \theta, \phi) = \sum_{l=0}^{\infty} \sum_{m=0}^l P_{lm}(\cos \theta) [C_{lm}(r) \cos m\phi + D_{lm}(r) \sin m\phi],$$

where P_{lm} is a Legendre function, and A_{lm} , B_{lm} , C_{lm} , and D_{lm} are

$$\begin{bmatrix} A_{lm}(r) \\ B_{lm}(r) \\ C_{lm}(r) \\ D_{lm}(r) \end{bmatrix} = N_{lm} \sum_{n=0}^{\infty} \tilde{A}_{nl} \begin{bmatrix} \tilde{\rho}_{nl}(r) \\ \tilde{\rho}_{nl}(r) \\ \tilde{\Phi}_{nl}(r) \\ \tilde{\Phi}_{nl}(r) \end{bmatrix} \begin{bmatrix} \Sigma_{nlm}(\cos) \\ \Sigma_{nlm}(\sin) \\ \Sigma_{nlm}(\cos) \\ \Sigma_{nlm}(\sin) \end{bmatrix}.$$

Here, N_{lm} , \tilde{A}_{nl} , $\tilde{\rho}_{nl}$, and $\tilde{\Phi}_{nl}$ are

$$N_{lm} = \frac{2l+1}{4\pi} (2 - \delta_{m0}) \frac{(l-m)!}{(l+m)!},$$

$$\tilde{A}_{nl} = -\frac{2^{8l+6}}{K_{nl}} \frac{n! (n+2l+3/2) [\Gamma(2l+3/2)]^2}{4\pi \Gamma(n+4l+3)},$$

$$\tilde{\rho}_{nl} = \frac{K_{nl}}{2\pi} \frac{r^l}{r(1+r)^{2l+3}} C_n^{(2l+3/2)} \left(\frac{r-1}{r+1} \right) \sqrt{4\pi},$$

$$\tilde{\Phi}_{nl} = -\frac{r^l}{(1+r)^{2l+1}} C_n^{(2l+3/2)} \left(\frac{r-1}{r+1} \right) \sqrt{4\pi},$$

where C_n^α is a Gegenbauer polynomial, and K_{nl} is a normalization constant of

$$K_{nl} = n(n+4l+3)/2 + (l+1)(2l+1).$$

Multipole expansions to describe the Galactic structure

The expansion coefficients of $\Sigma_{nlm}(\cos)$ and $\Sigma_{nlm}(\sin)$ are calculated with a collection of k th particles by

$$\begin{bmatrix} \Sigma_{nlm}(\cos) \\ \Sigma_{nlm}(\sin) \end{bmatrix} = \sum_k m_k \tilde{\Phi}_{nl}(r_k) P_{lm}(\cos \theta_k) \begin{bmatrix} \cos m\phi_k \\ \sin m\phi_k \end{bmatrix}.$$

-> making tables of $\Sigma_{nlm}(\cos)$ and $\Sigma_{nlm}(\sin)$

-> reading the tables when a simulation run starts

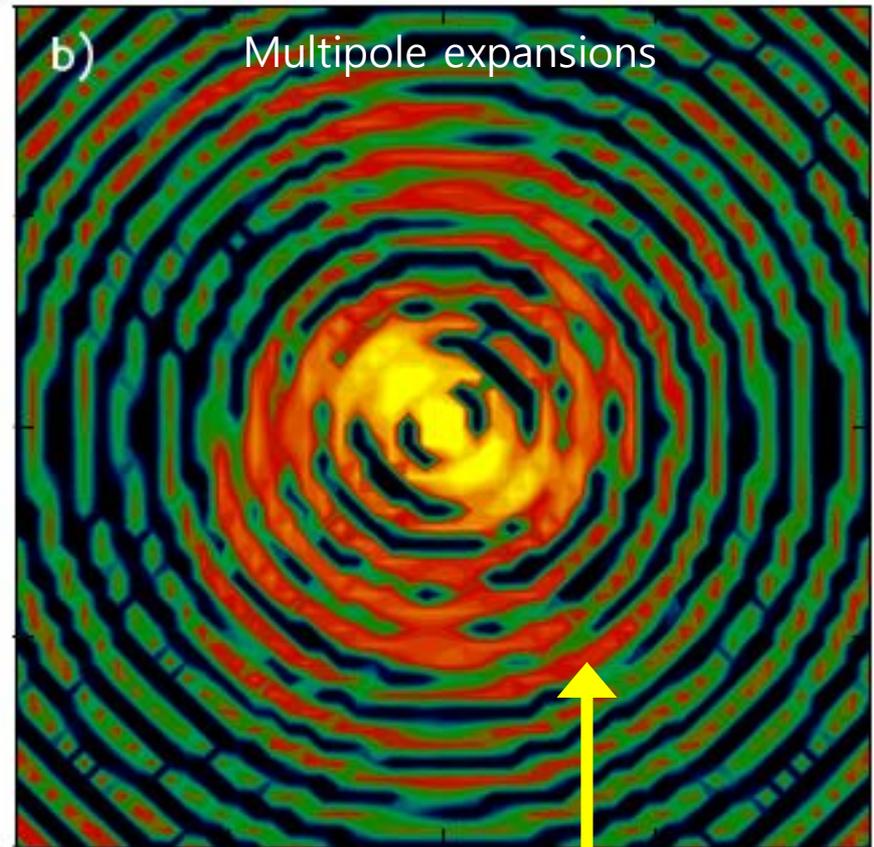
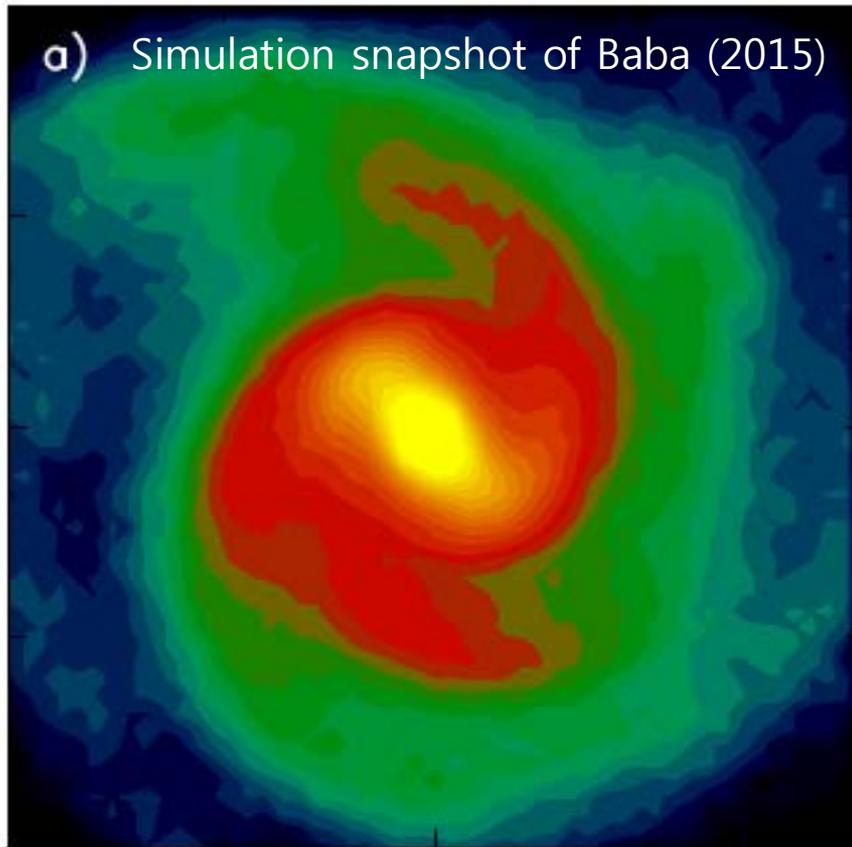
Similar to ρ and Φ , accelerations a_r , a_θ , and a_ϕ are derived by

$$\begin{aligned} a_r(r, \theta, \phi) &= - \sum_{l=0}^{\infty} \sum_{m=0}^l P_{lm}(\cos \theta) [E_{lm}(r) \cos m\phi + F_{lm}(r) \sin m\phi], \\ a_\theta(r, \theta, \phi) &= - \frac{1}{r} \sum_{l=0}^{\infty} \sum_{m=0}^l \frac{dP_{lm}(\cos \theta)}{d\theta} [C_{lm}(r) \cos m\phi + D_{lm}(r) \sin m\phi], \\ a_\phi(r, \theta, \phi) &= - \frac{1}{r} \sum_{l=0}^{\infty} \sum_{m=0}^l \frac{mP_{lm}(\cos \theta)}{\sin \theta} [D_{lm}(r) \cos m\phi - C_{lm}(r) \sin m\phi], \end{aligned}$$

where E_{lm} and F_{lm} are

$$\begin{bmatrix} E_{lm}(r) \\ F_{lm}(r) \end{bmatrix} = N_{lm} \sum_{n=0}^{\infty} \tilde{A}_{nl} \frac{d}{dr} \tilde{\Phi}_{nl}(r) \begin{bmatrix} \Sigma_{nlm}(\cos) \\ \Sigma_{nlm}(\sin) \end{bmatrix}.$$

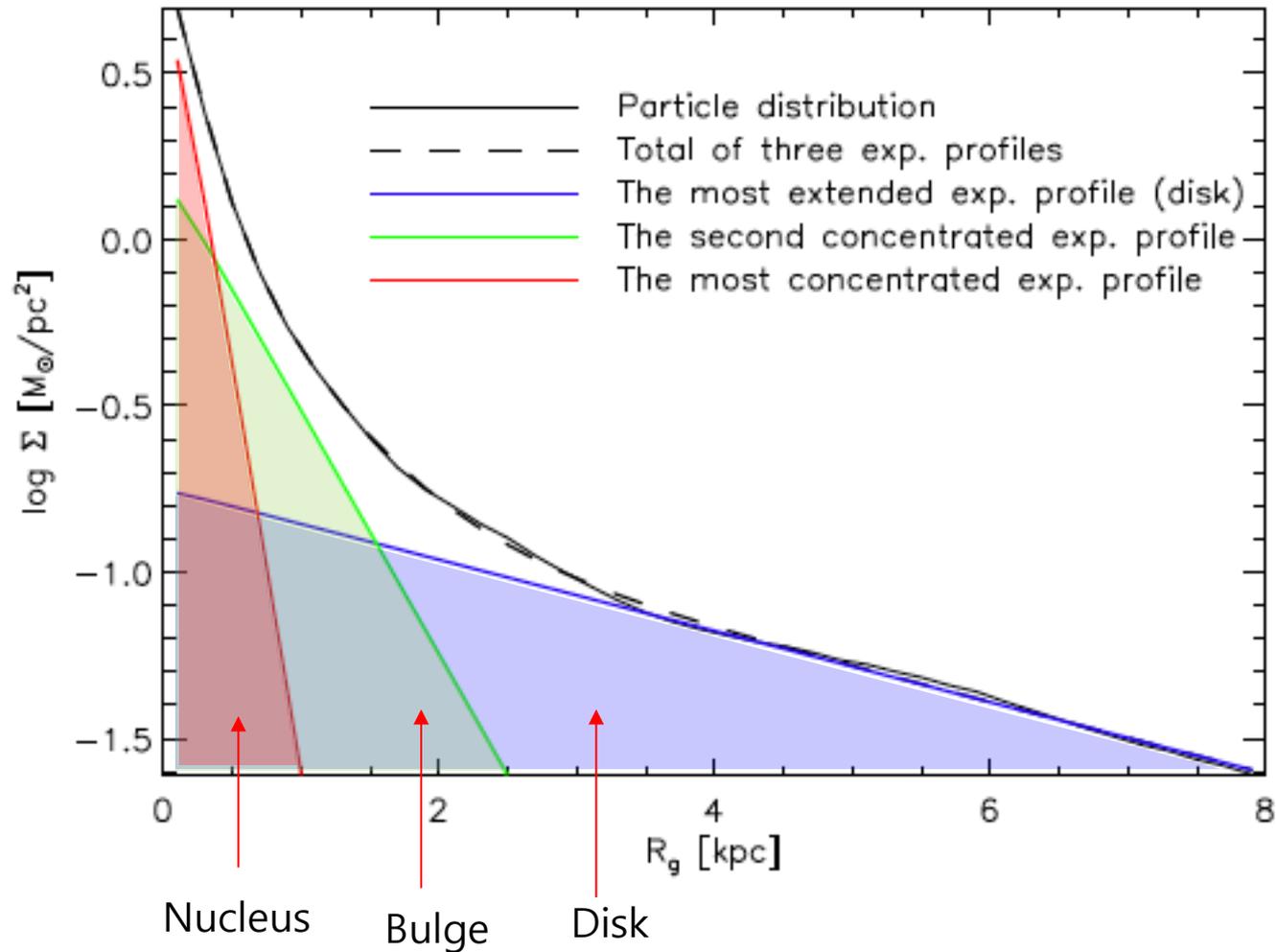
1. 3D-disk model of Earn (1996)



Radial density fluctuation

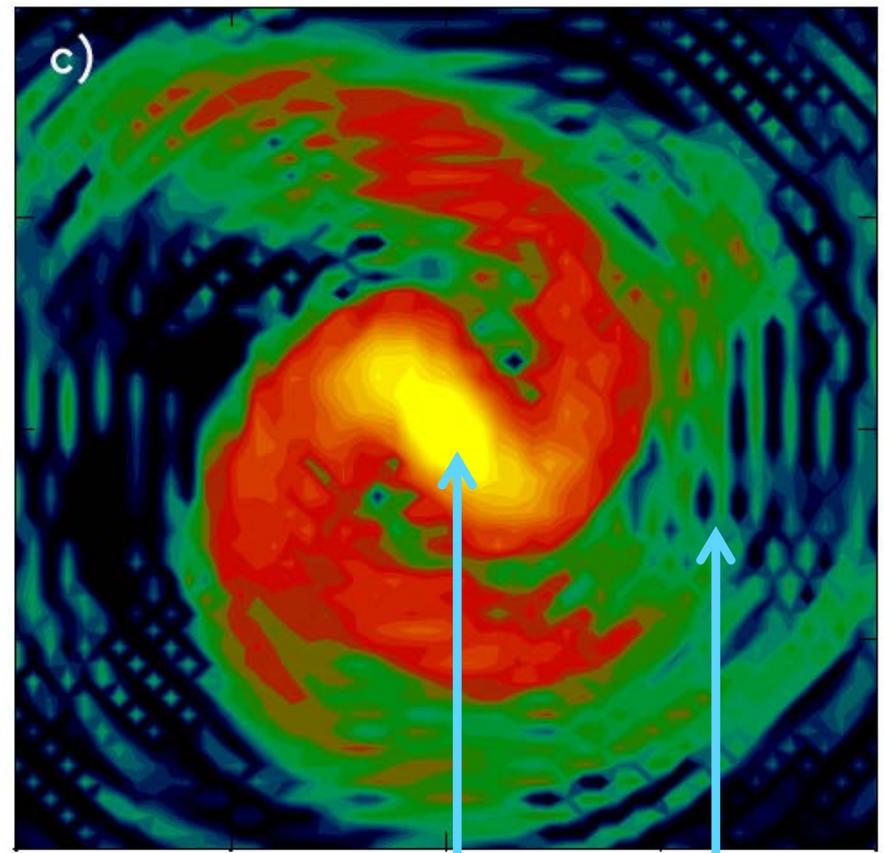
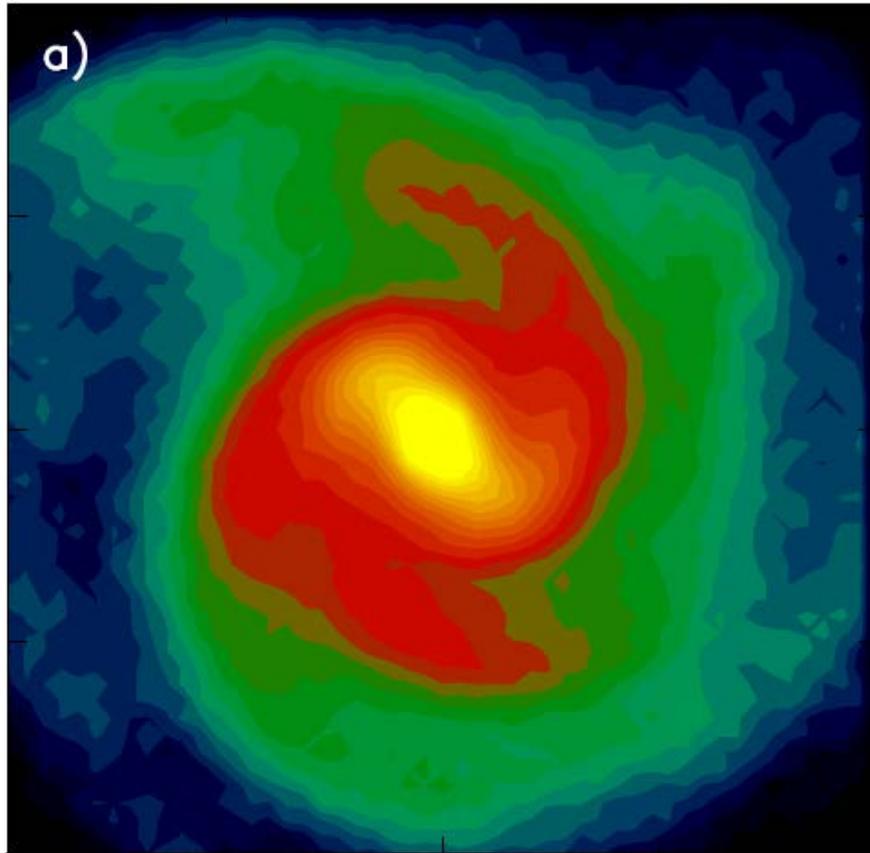
<Projected stellar density distribution>

Decomposition of stellar particle



$$\rho(R_g, z) = \Sigma_0 \exp(-R_g/R_d) [\exp(-|z|/h_z)/(2h_z)]$$

2. Hybrid models (3D-spherical + 3D-disk)



Nucleus+Bulge components : 3D-spherical model of Hernquist & Ostiker (1992)

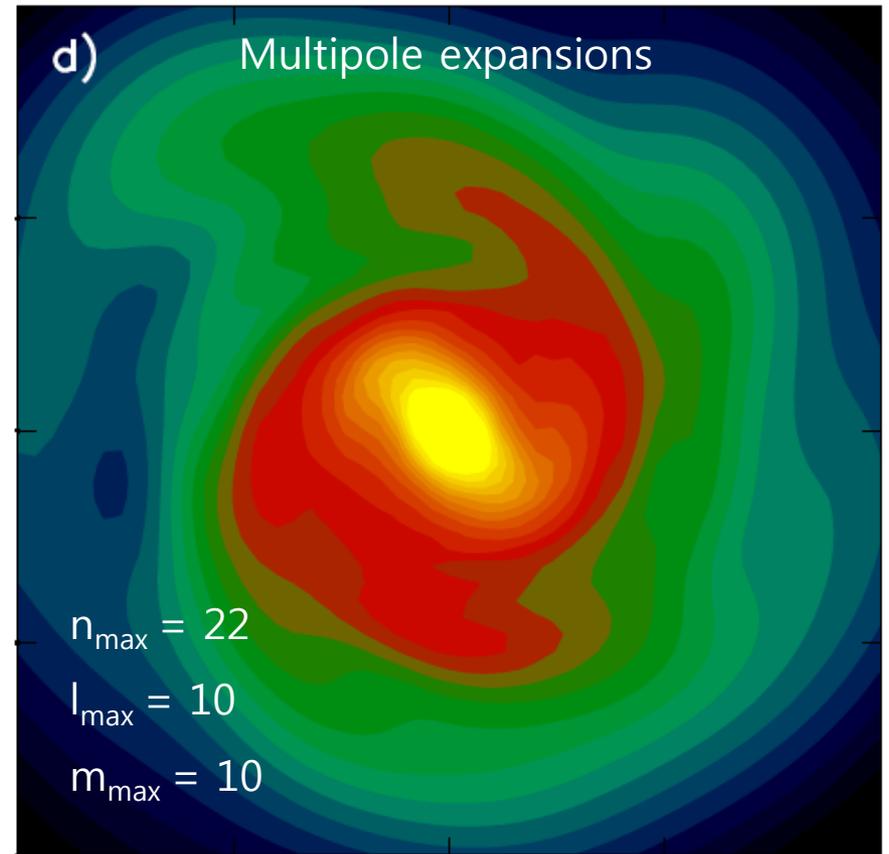
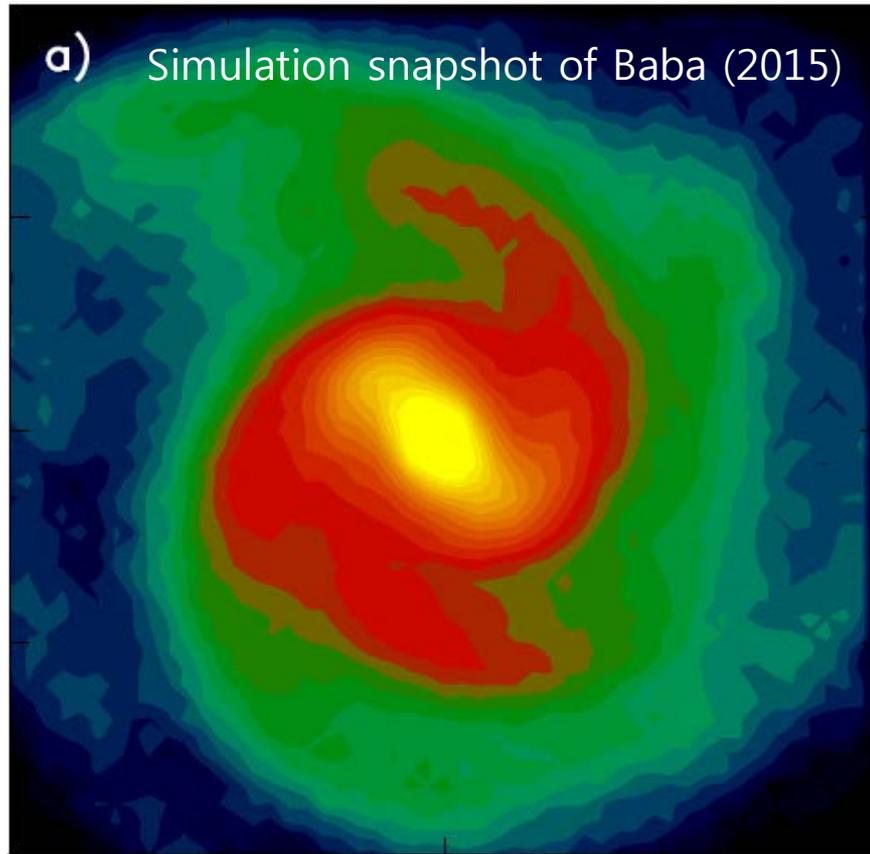
Disk component : 3D-disk model of Earn (1996)

Bulge

Disk

<Projected stellar density distribution>

3. Hybrid models (3D-spherical + 2D-disk)



Nucleus+Bulge components : 3D-spherical model of Hernquist & Ostiker (1992)

Disk component : 2D-disk model of Aoki & Iye (1978)

<Projected stellar density distribution>

Thick disk correction (Binney & Tremaine 2008)

To overcome a limitation of the 2D-disk model...

3D disk

For a thick disk whose density profiles following

$$\rho(R_g, z) = \Sigma_0 \exp(-R_g/R_d) [\exp(-|z|/h_z)/(2h_z)],$$

Potential at a position (R_g, z) can be derived as

$$\Phi(R, z) = -\frac{4G\Sigma_0}{R_d} \int_{-\infty}^{\infty} dz' \frac{\exp(z'/h_z)}{2h_z} \int_0^{\infty} da \sin^{-1} \left(\frac{2a}{\sqrt{+} + \sqrt{-}} \right) a K_0(a/R_d)$$

2D disk

For an instantaneously thin disk whose density profile following

$$\Sigma(R) = \Sigma_0 \exp(-R/R_d),$$

Potential at a position $(R_g, 0)$ can be derived as

$$\Phi(R, 0) = -\pi G \Sigma_0 R [I_0(y)K_1(y) - I_1(y)K_0(y)],$$

Modified Bessel functions



Thick disk correction (Binney & Tremaine 2008)

Difference of a_R between the 2D and 3D exponential disks :

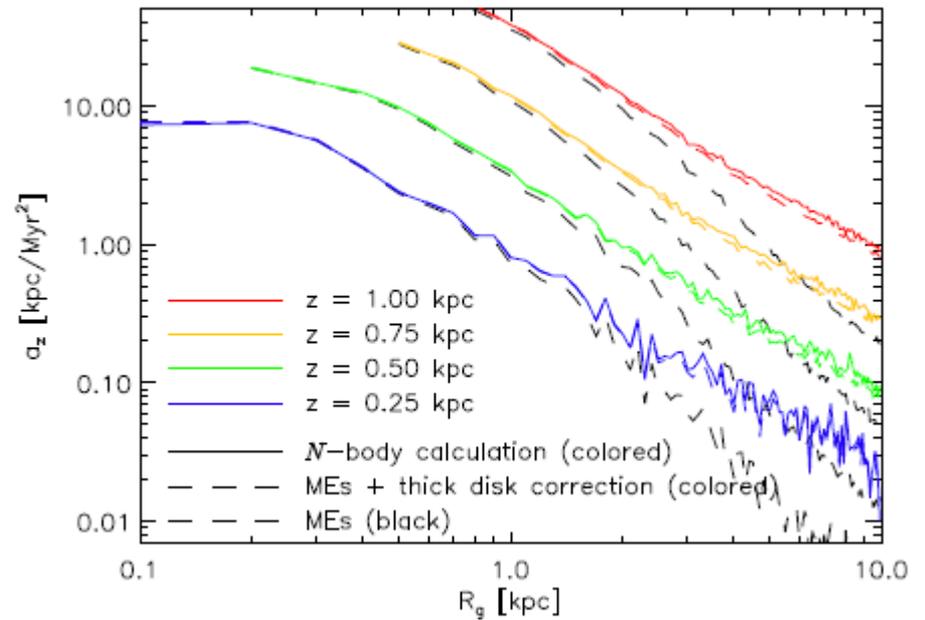
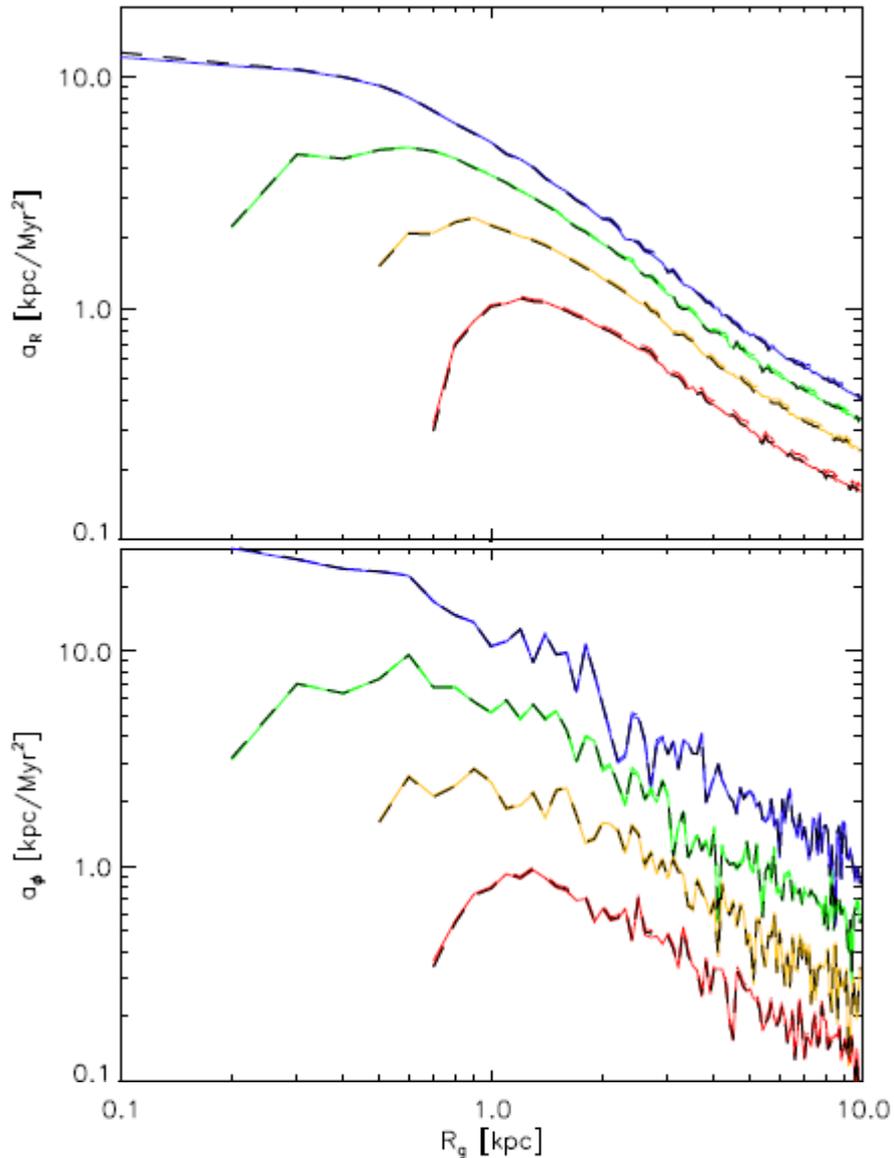
$$\begin{aligned}\Delta a_R(R, z) &= a_R(R, z) - a_{R,0}(R) \\ &= -d\Phi(R, z)/dR + d\Phi_0(R)/dR\end{aligned}$$

a_z due to the non-zero disk thickness :

$$a_z(R, z) = -d\Phi(R, z)/dz$$

Δa_R and a_z : equally for $\phi = 0 - 2\pi$

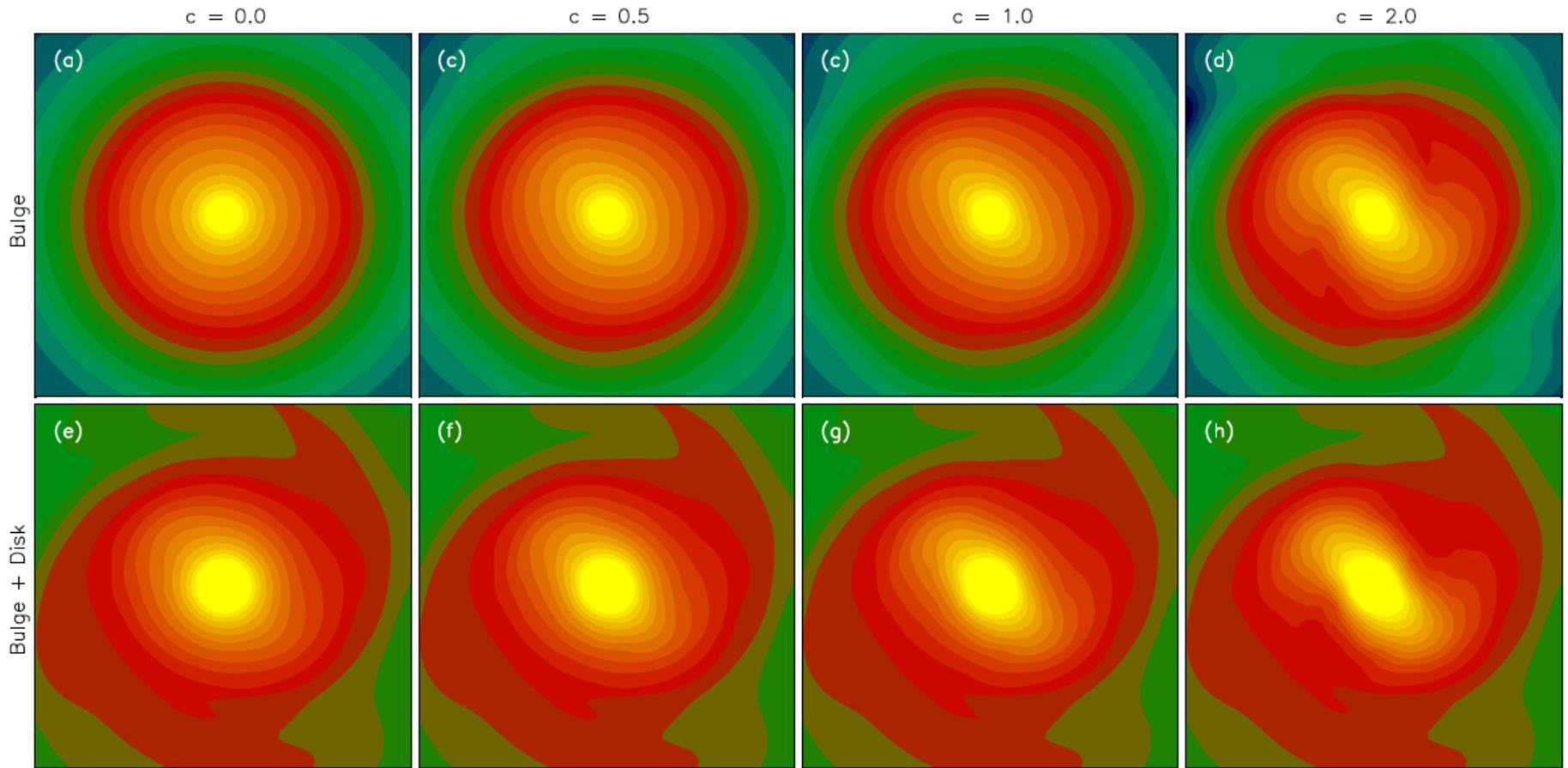
Averaged acceleration profiles



The Galactic structures are well reproduced by **2D ME + 3D ME + thick disk correction.**

Advantages of the ME method

1. We can describe the **realistic mass distribution** of the Milky Way.
2. Performance : **~ 30 times faster** than 10 million live particles
3. We can **enlarge/diminish a bar elongation** by modifying ME coefficients (even modes).



Gradual bar growth

Simulation code

Simulation code : based on a parallel N-body/SPH code, **Gadget-2** (Springel 2005)

- [Radiative Heating & Cooling](#) by Cloudy 90 package (Ferland et al. 1998)
 - : ionizing photon flux = 1, 100 G0 (Habing field)
- [Star formation](#) at high density ($n_{\text{H}} > 100 \text{ cm}^{-3}$) and low temperature ($T < 100\text{K}$)
- [SN II feedback](#) (time-step limiter of Durier & Dallar Vecchia 2012, Saitoh & Makino 2009)
- High performing time-stepping (Saitoh & Makino 2010)
- [Realistic Galactic potential](#) using multipole expansions
 1. tabulating ME coefficients before a simulation
 2. Reading the ME tables when a simulation starts
 3. Calculating external force for a given particle position

- Initial particles : only for gas

Simulation parameters

MODEL PARAMETERS

Model	N_{gas}	$R_{IB} - R_{OB}$ [kpc]	Q	τ_{growth} [Myr]	Ω_{bar} [km/s/kpc]
A	10^5	0 - 6	0.284	0	35
Aq1	10^5	0 - 6	0.252	350	35
Aq2	10^5	0 - 6	0.268	350	35
Aq3	10^5	0 - 6	0.299	350	35
Aq4	10^5	0 - 6	0.315	350	35
At1	10^5	0 - 6	0.284	175	35
At2	10^5	0 - 6	0.284	350	35
At3	10^5	0 - 6	0.284	700	35
Ap1	10^5	0 - 6	0.284	350	30
Ap2	10^5	0 - 6	0.284	350	40
B	10^5	1 - 6	0.284	0	35
Bq1	10^5	1 - 6	0.252	350	35
Bq2	10^5	1 - 6	0.268	350	35
Bq3	10^5	1 - 6	0.299	350	35
Bq4	10^5	1 - 6	0.315	350	35
Bt1	10^5	1 - 6	0.284	175	35
Bt2	10^5	1 - 6	0.284	350	35
Bt3	10^5	1 - 6	0.284	700	35
Bp1	10^5	1 - 6	0.284	350	30
Bp2	10^5	1 - 6	0.284	350	40
Ah	3×10^5	0 - 6	0.284	0	35
At2h	3×10^5	0 - 6	0.252	350	35
Bh	3×10^5	1 - 6	0.284	0	35
Bt2h	3×10^5	1 - 6	0.252	350	35

Initial gas distribution : $R_{IB} - R_{OB}$

- Type 'A' = 0 – 6 kpc,
- Type 'B' = 1 – 6 kpc

Total number of gas particles :

- normal : 10^5
- high resolution : 3×10^5 (Ah, At2h, Bh, Bt2h)

Individual gas mass :

- normal : $7.5-8.4 \times 10^3 M_{sun}$
- high : $2.5-2.8 \times 10^3 M_{sun}$

Bar growth timescale (τ_{growth}) :

- 0, 175, 350, 700 Myr (A, At1, At2, At3)

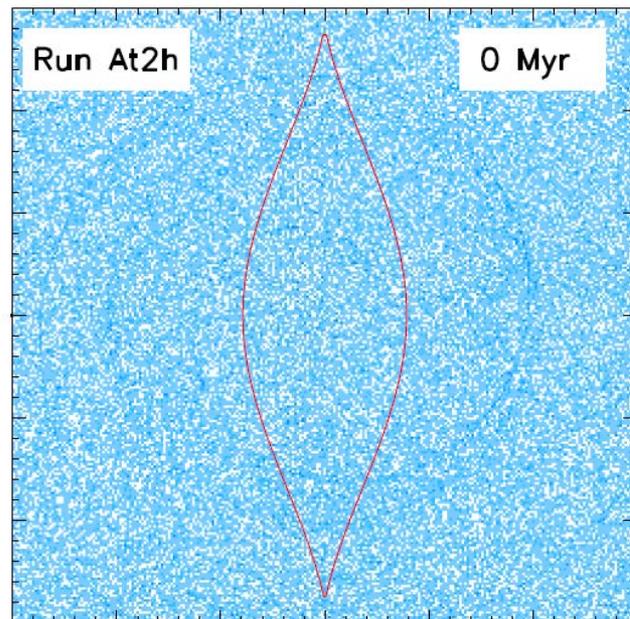
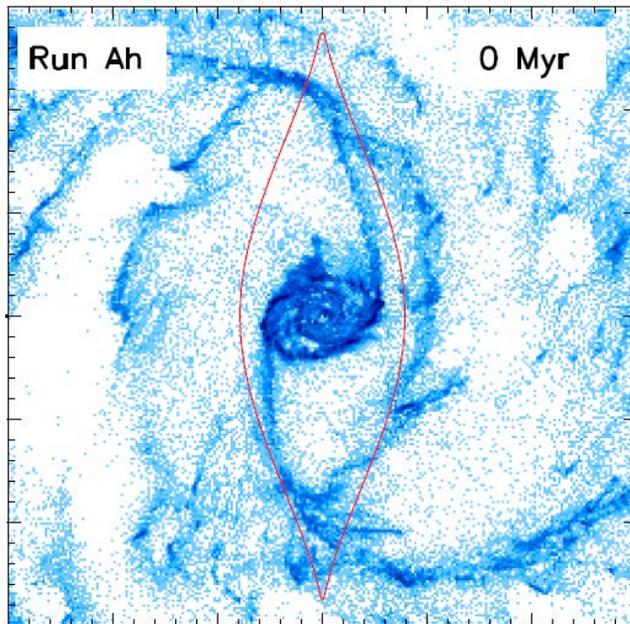
Pattern speed (Ω_{bar}) :

- 30, 35, 40 km/s/kpc (Ap1, A, Ap2)

Bar strength (Q) :

- 0.252, 0.268, 0.284, 0.299, 0.31 (Aq1, Aq2, A, Aq3, Aq4)

Simulation parameters



Initial gas distribution : $R_{IB} - R_{OB}$

- Type 'A' = 0 – 6 kpc,
- Type 'B' = 1 – 6 kpc

Total number of gas particles :

- normal : 10^5
- high resolution : 3×10^5 (Ah, At2h, Bh, Bt2h)

Individual gas mass :

- normal : $7.5-8.4 \times 10^3 M_{\text{sun}}$
- high : $2.5-2.8 \times 10^3 M_{\text{sun}}$

Bar growth timescale (τ_{growth}) :

- 0, 175, 350, 700 Myr (A, At1, At2, At3)

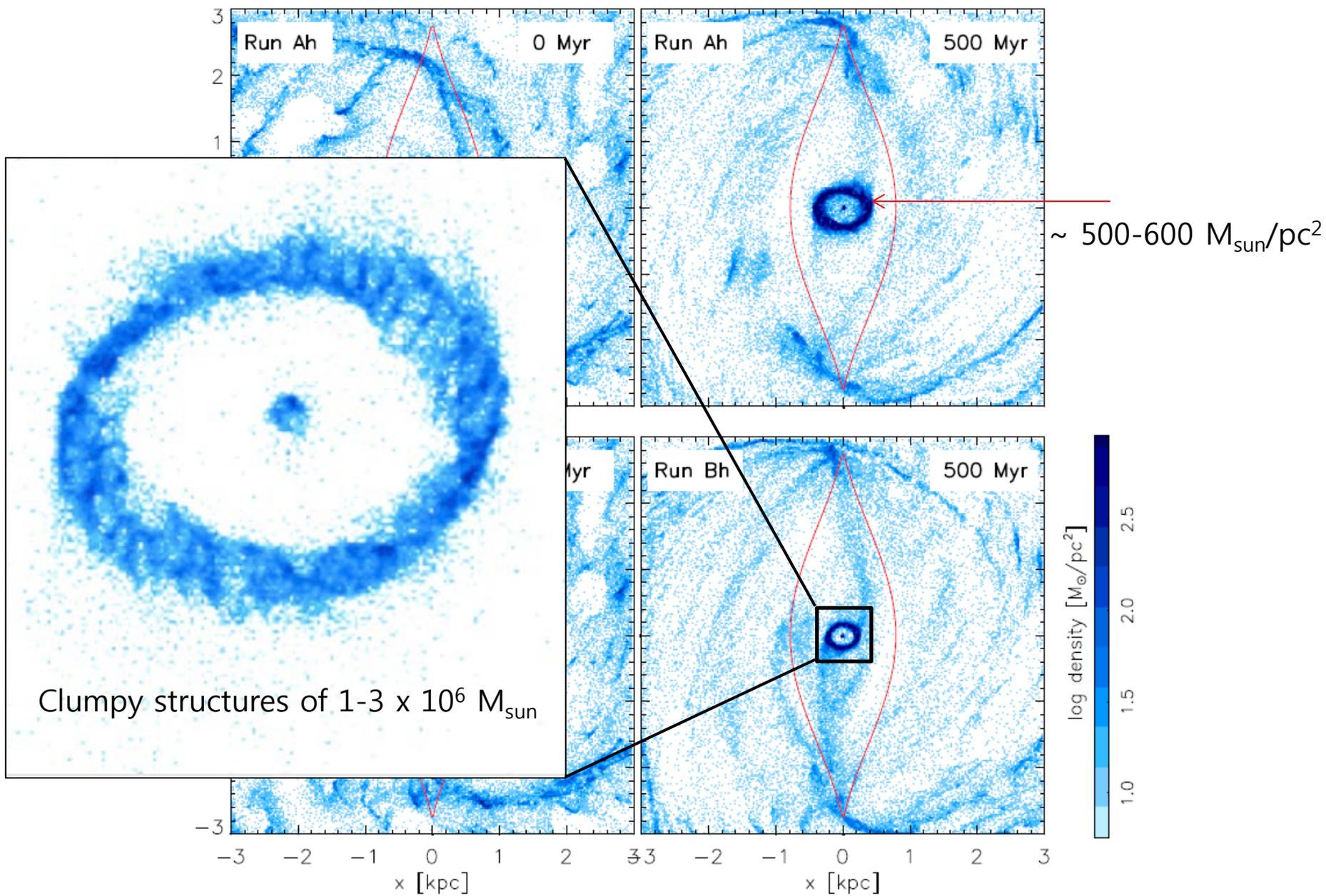
Pattern speed (Ω_{bar}) :

- 30, 35, 40 km/s/kpc (Ap1, A, Ap2)

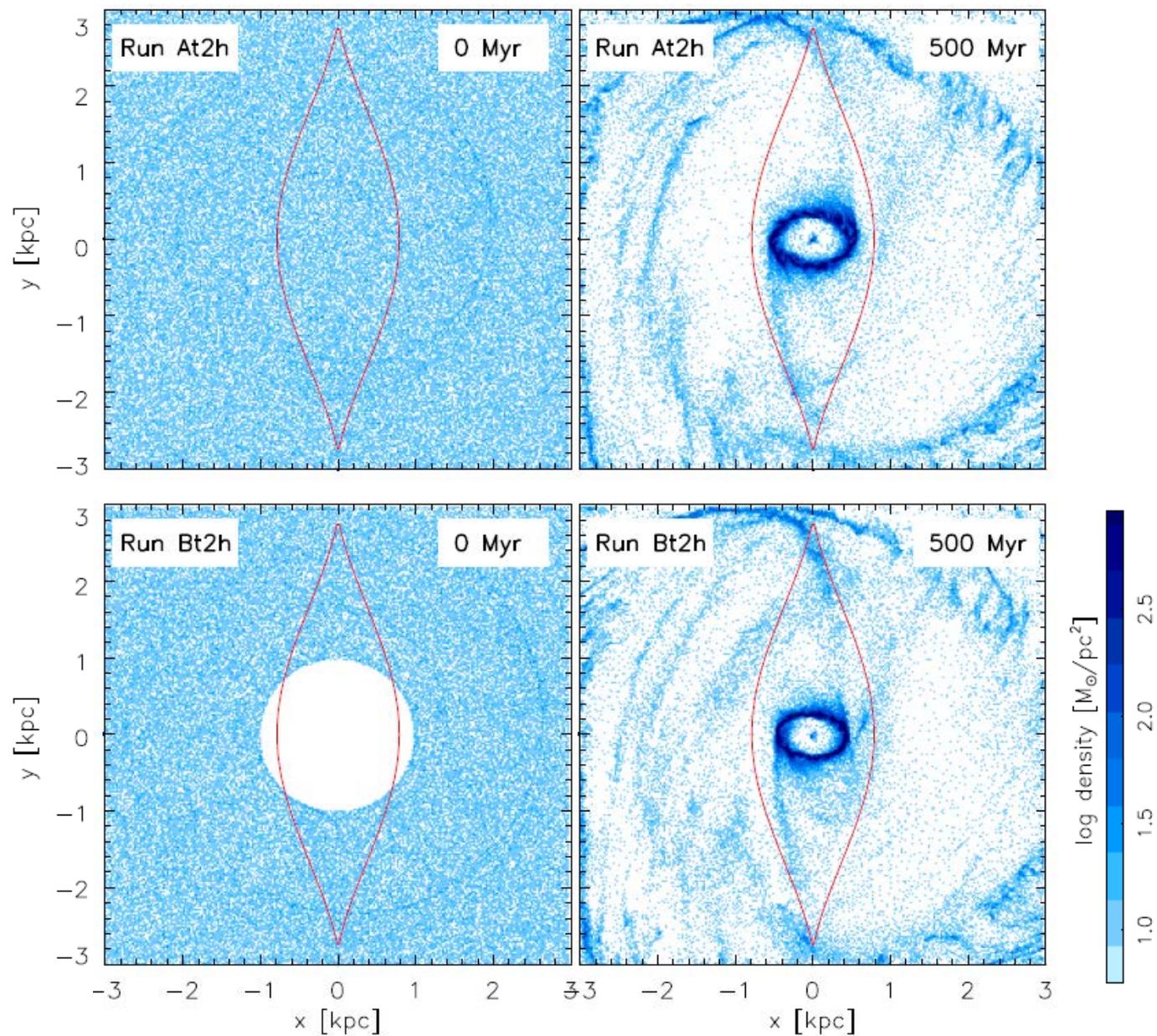
Bar strength (Q) :

- 0.252, 0.268, 0.284, 0.299, 0.31 (Aq1, Aq2, A, Aq3, Aq4)

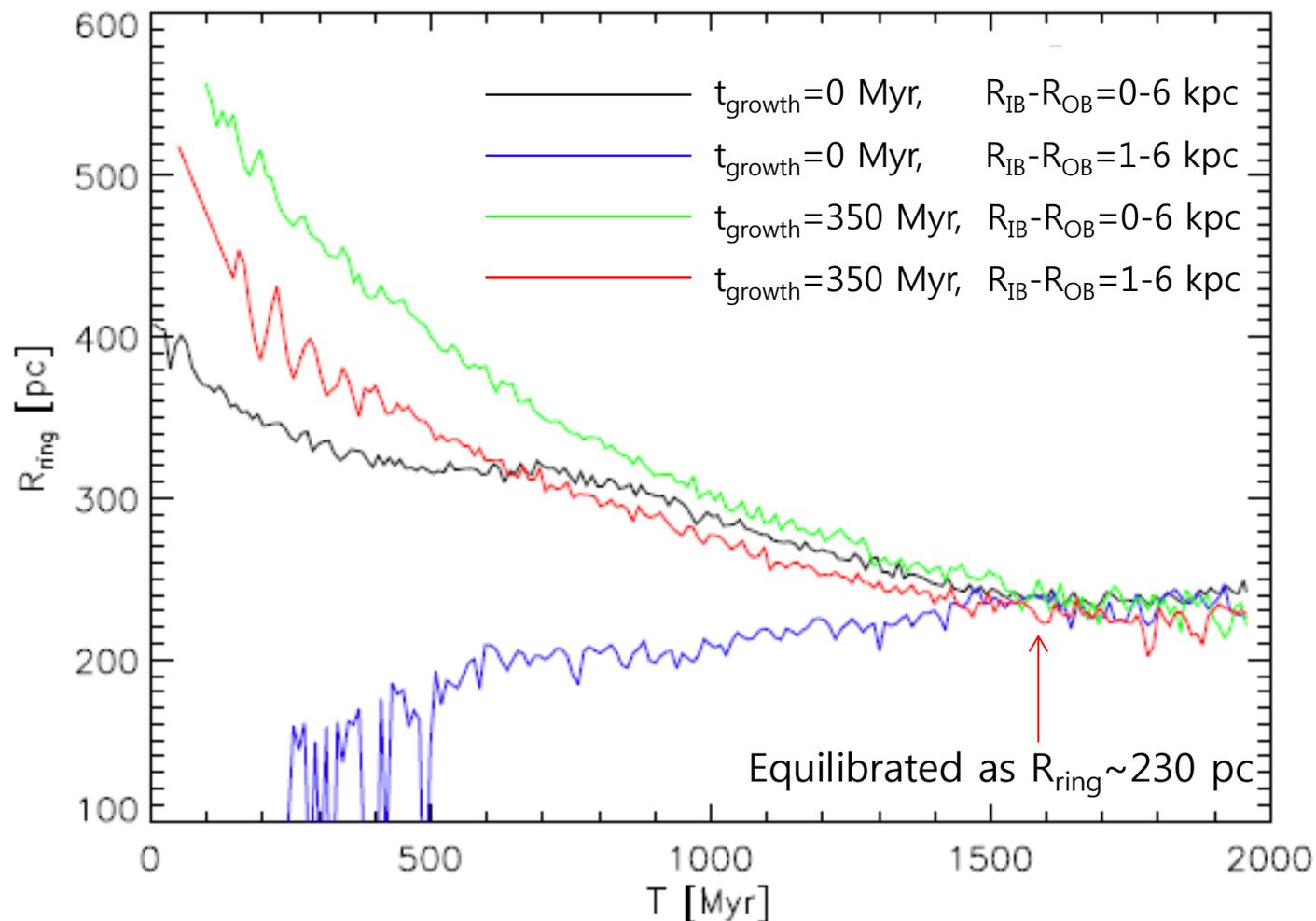
Overall morphology



Overall morphology



Size evolution the CMZ

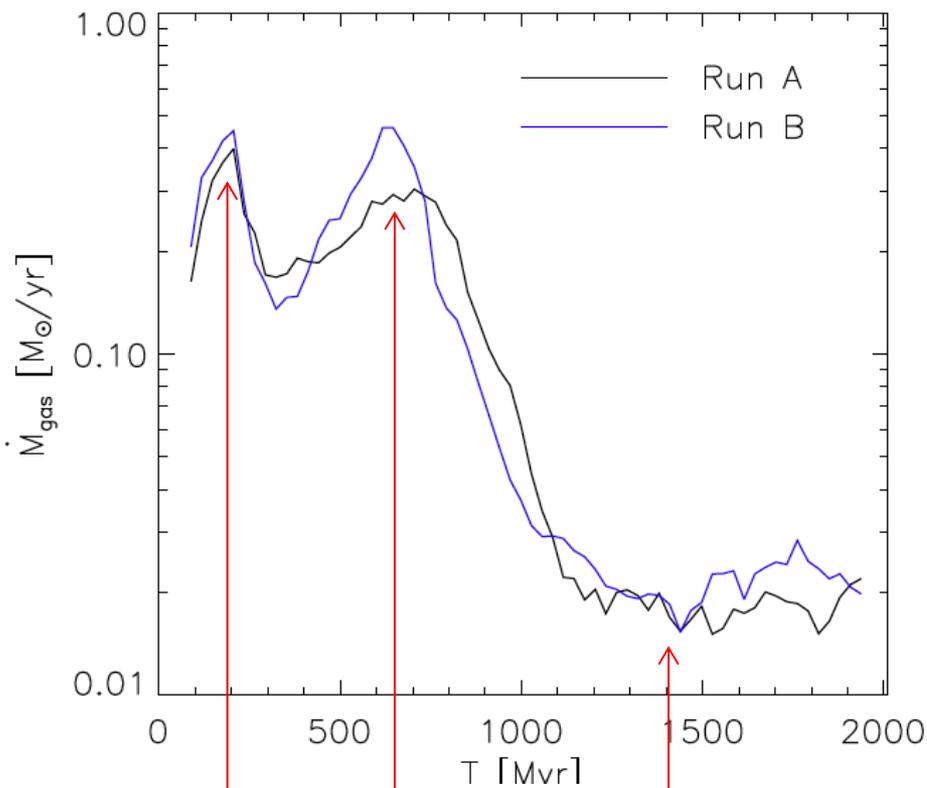


Regardless of the initial value,

all R_{ring} converge to the same value as ~ 230 pc at $T \sim 1500$ Myr.

=> a bit larger than the outer boundary of the CMZ, ~ 200 pc (Morris & Serabyn 1996)

Gas inflows to the CMZ

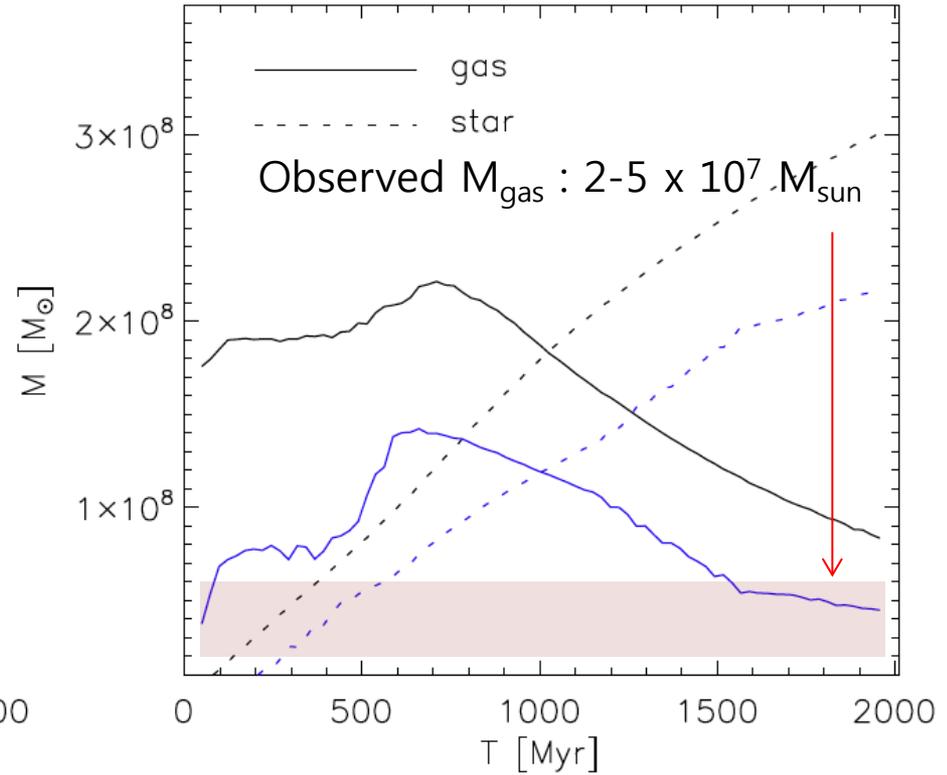
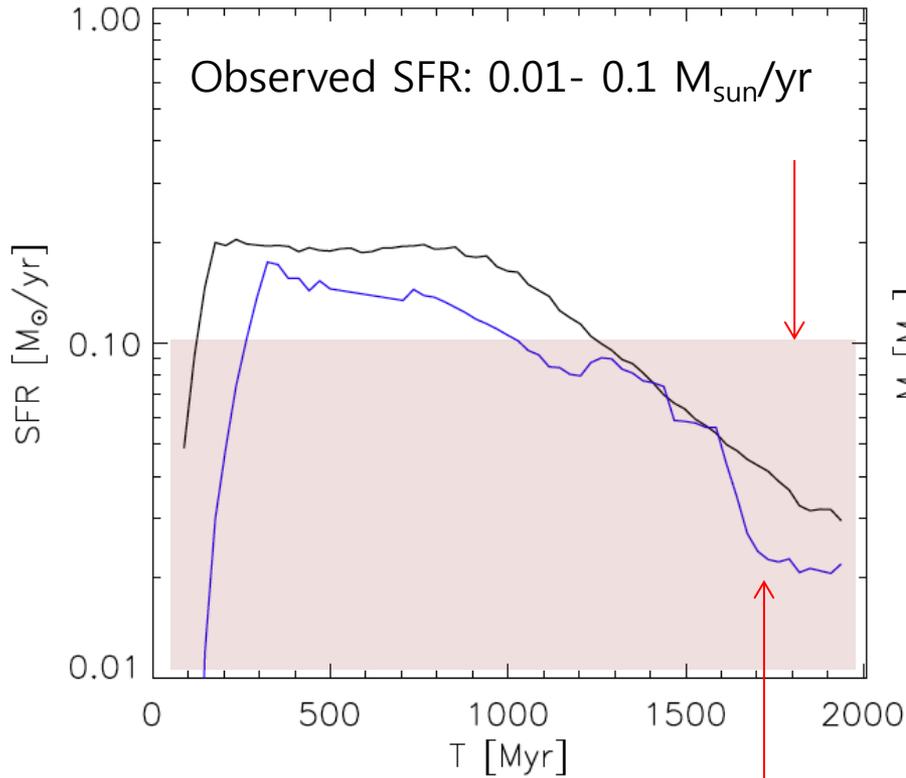


Equilibrated as $\sim 0.02 M_{\text{sun}}/\text{yr}$

Enhanced again by spiral shocks (Kim & Kim 2014, Seo & Kim 2014)

Initial relaxation period (due to unrelaxed gas clouds for the Galactic structure)

Gas inflows to the CMZ

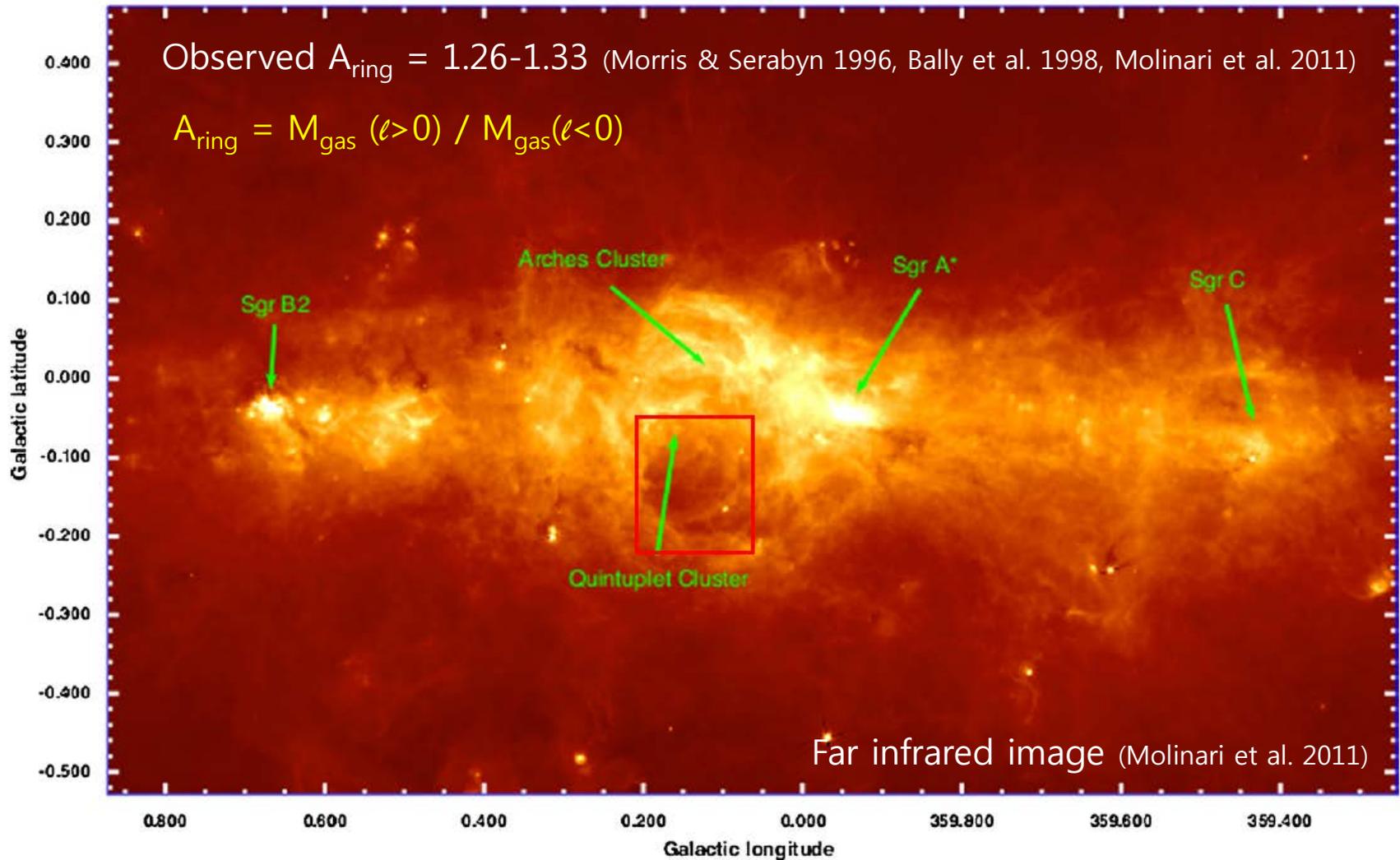


SFR is expected to be equilibrated as $\sim 0.02 M_{\text{sun}}/\text{yr}$

Regardless of t_{growth} and $R_{\text{IB}}-R_{\text{OB}}$,

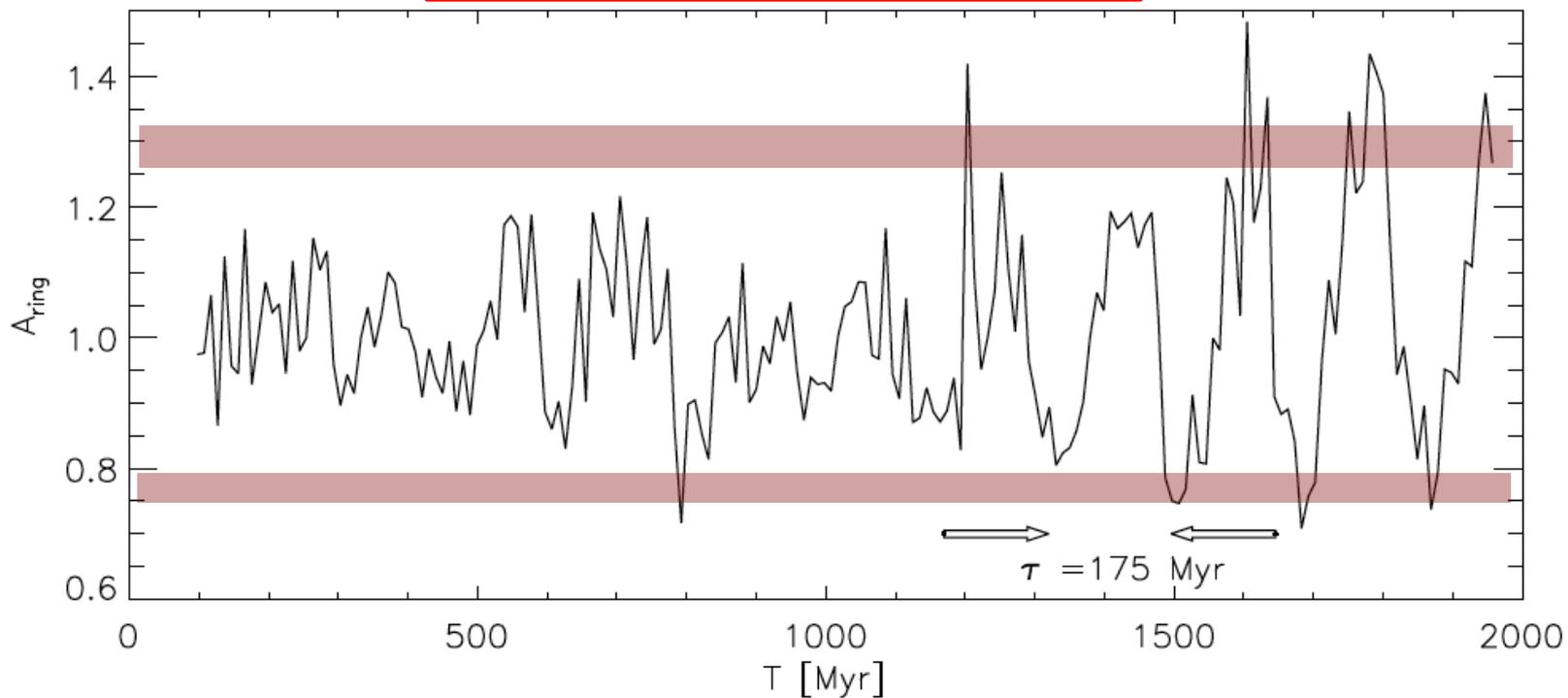
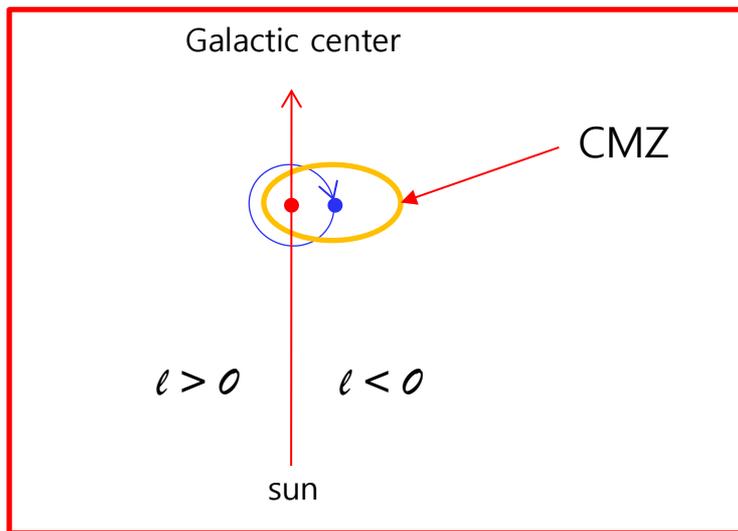
\dot{M}_{gas} , SFR and M_{gas} converge to similar ranges.

Asymmetry of the CMZ

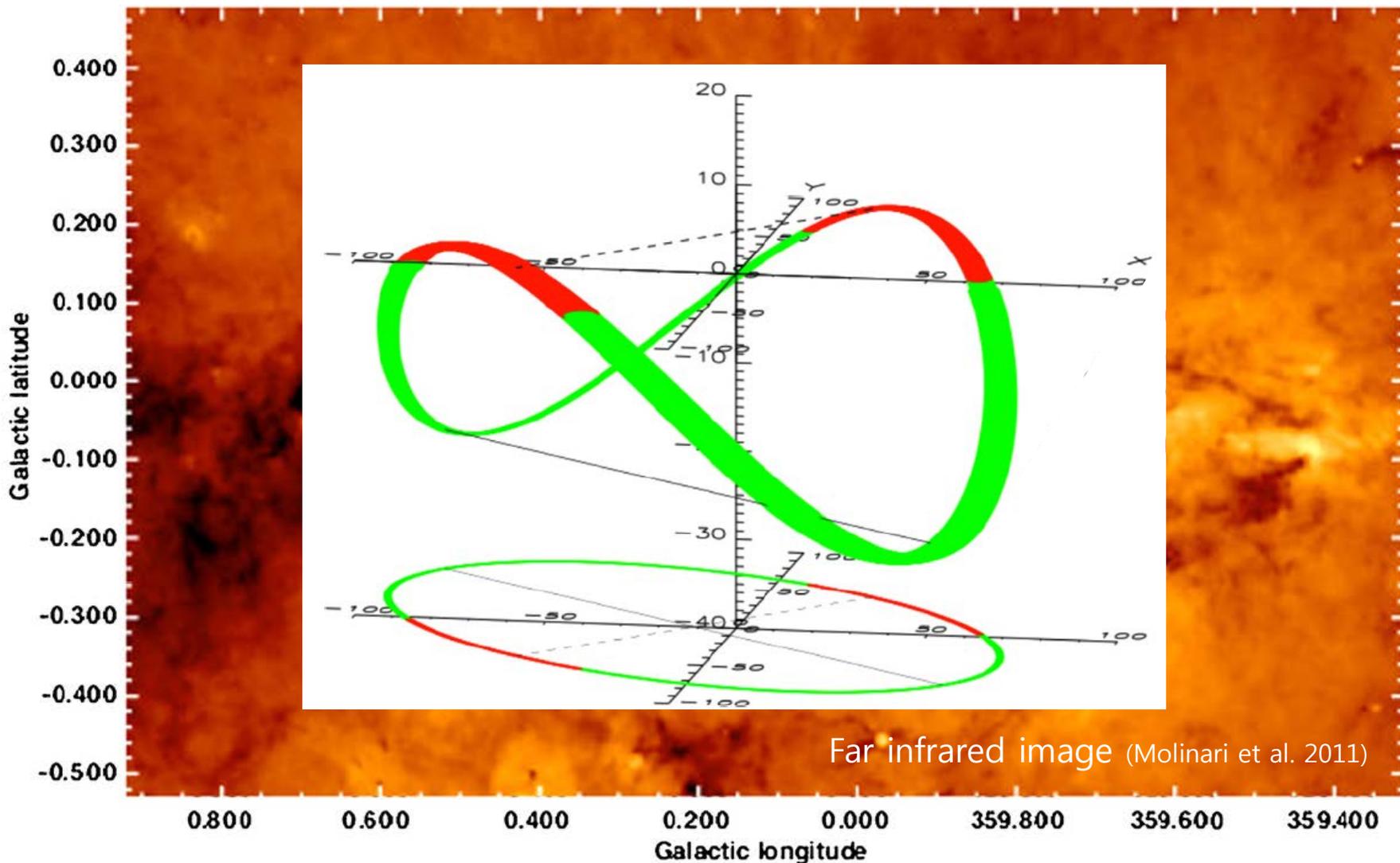


The Galactic structures of Baba (2015) -> a **lopsided central mass distribution**
-> an **off-centered density peak** of central mass distribution to the galactic center

Asymmetry of the CMZ

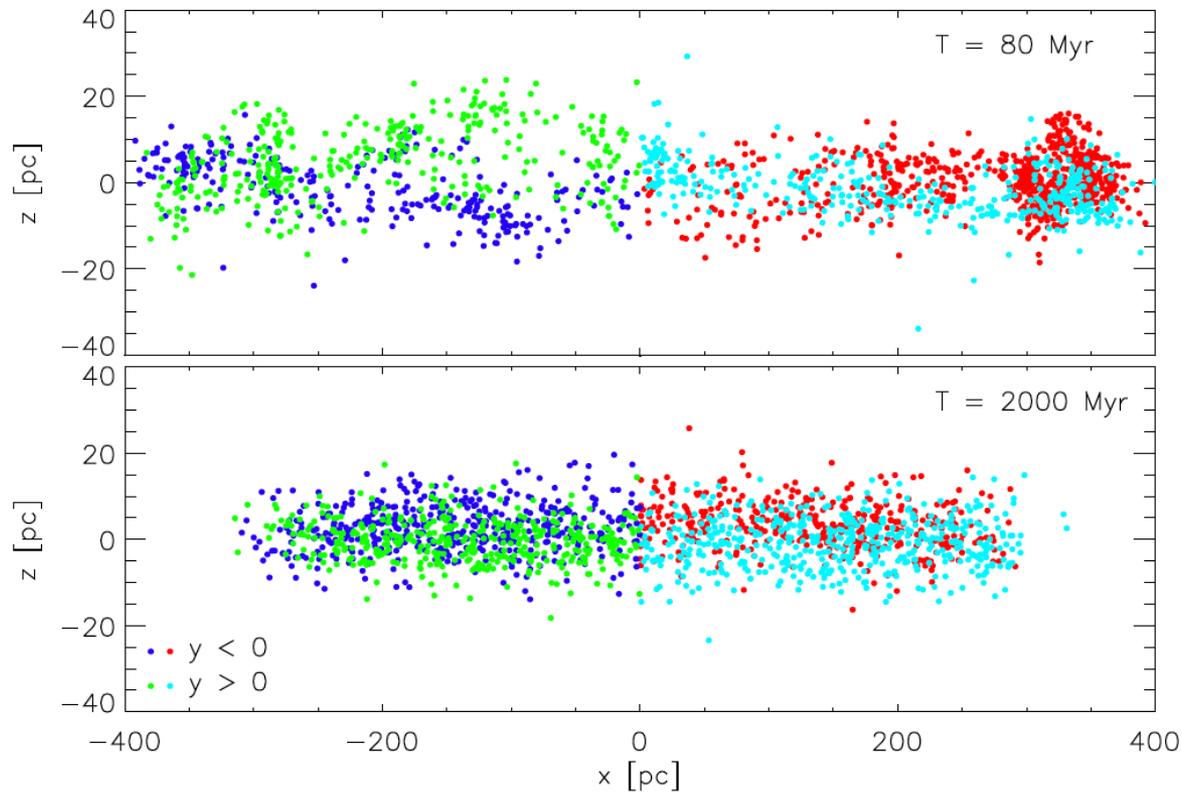


Twisted (∞ -like) shape of the CMZ



A z-directional lopsidedness of the central mass distribution -> z-directional motion

Twisted (infinity-like) shape of the CMZ



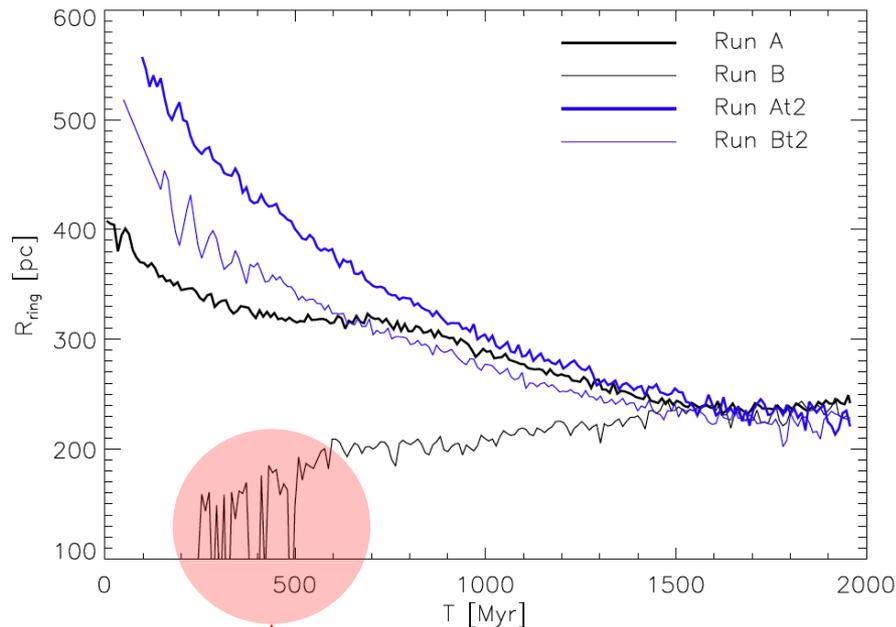
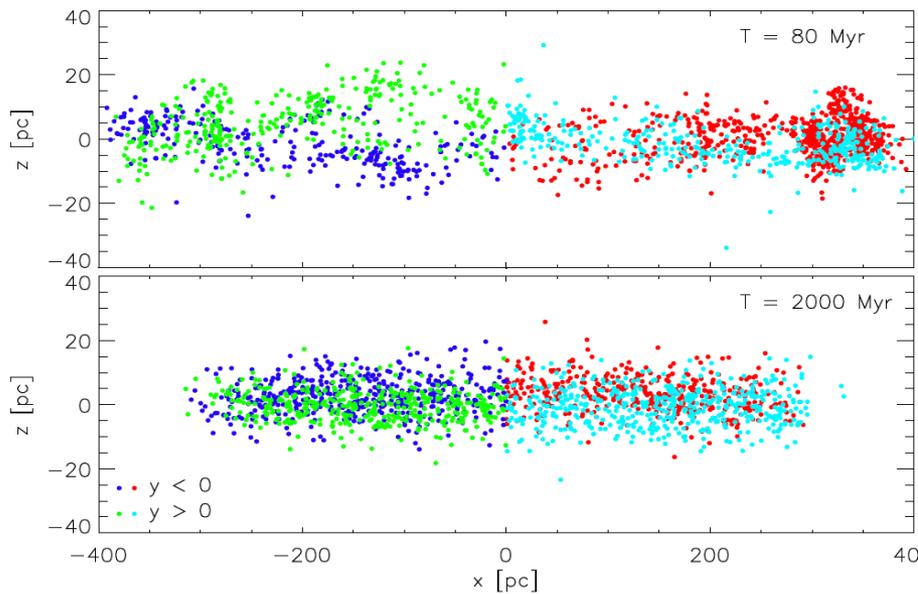
Sinusoidal shape : oscillations are synchronized to have the same phase and frequency to each other.

First 100 Myr : the z-directional frequency = 2 times of an orbital frequency -> Infinity-like shape

After the first 100 Myr : the z-directional frequency = an orbital frequency -> Tilted to the Galactic plane

Thus, the infinity-like shapes of the CMZ is dynamically young feature.

CMZ is dynamically young?



Tilted HI disk

The compact size of the CMZ is only reproduced when

- nuclear gaseous structure is **dispersed** by something,
- but, gas inflow from the disk is **enhanced** by something.

We infer that the "something" might be a minor merging event.

Summary

The [ME method](#) was utilized to efficiently include [the realistic Galactic mass distribution](#).

- : accurate and fast
- : bar strength can be easily enlarged/diminished by modifying ME coefficients.

We trace [the realistic motion of gas clouds](#) from $\sim 3\text{kpc}$ to $\sim 200\text{pc}$.

- : R_{ring} is equilibrated as $\sim 230\text{ pc}$ (most approach to the current CMZ size)
- : observed M_{gas} , SFR and M_{gas} of the CMZ is well reproduced
- : the lopsided central mass distribution results in [the asymmetric mass distribution and also infinity-like feature of the CMZ](#).

[On-going projects](#) related with this work

- : how [a minor merging/fly-by](#) disperses the nuclear gaseous structure, and perturbs the gaseous disk (by Jihye Shin & Jeongsun Hwang)
- : effect of [gas replenishment](#) to the nuclear ring (by Jihye Shin & Kyungwon Chun)
- : Gas migration from [CMZ to CND](#) (by Hannah Morgan, KHU)
- : [tilted HI disk](#) (by Joowon Lee)
- : bar strength of external barred galaxies using [ME method](#) (by Eunbin Kim)